



### Casimir and Chameleon forces - the universe on the tabletop

René Sedmik

on behalf of the

Cannex team

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# Outline

## Part I: Scientific background

Casimir physics, Chameleon model and dark energy

What are the goals?

## Part II: Setup

Core: Principle of measurement, Control of parallelism,

Vibration insulation: Mechanisms, Preliminary characterization

Status

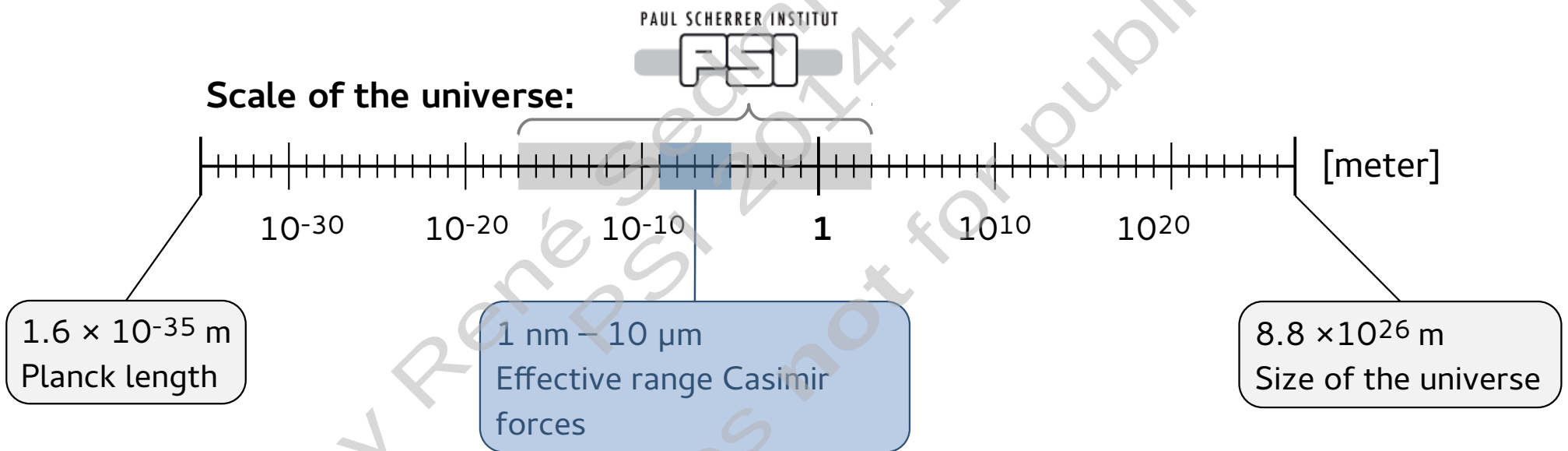
talk by René Seemann, VU Amsterdam  
PSI 2014-11-06  
slides not for public use

# Outline

## Part I: Scientific background

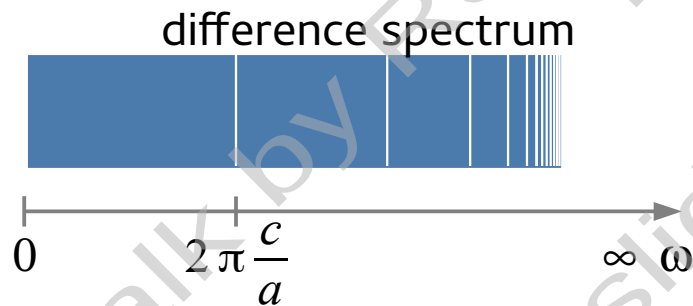
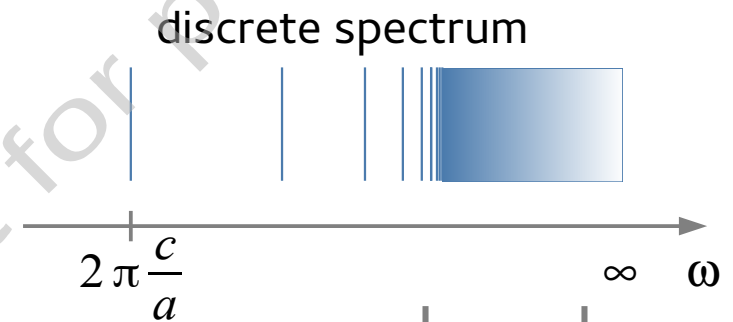
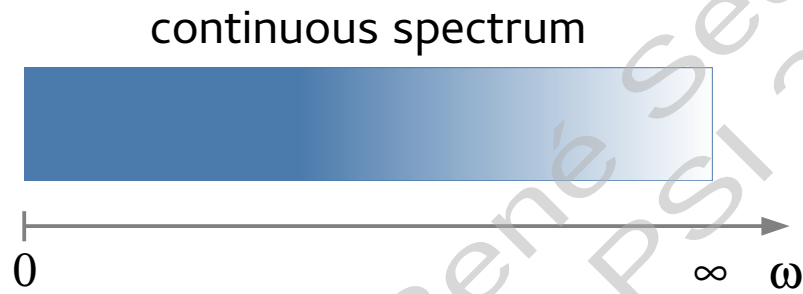
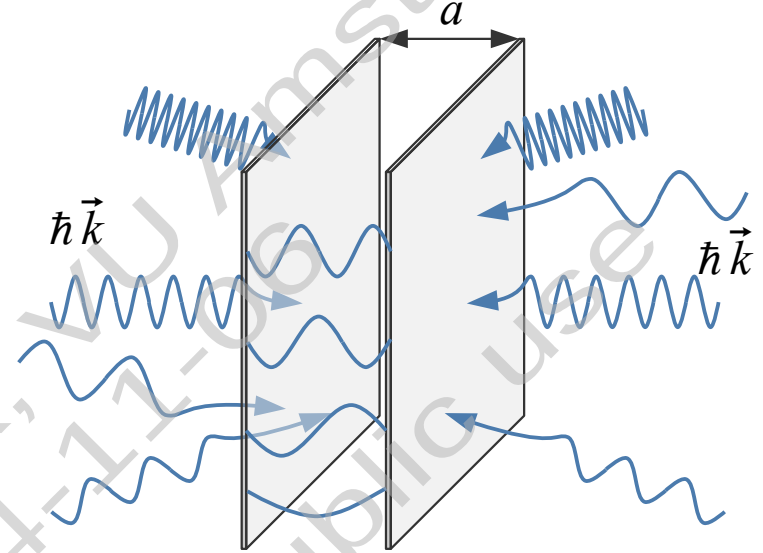
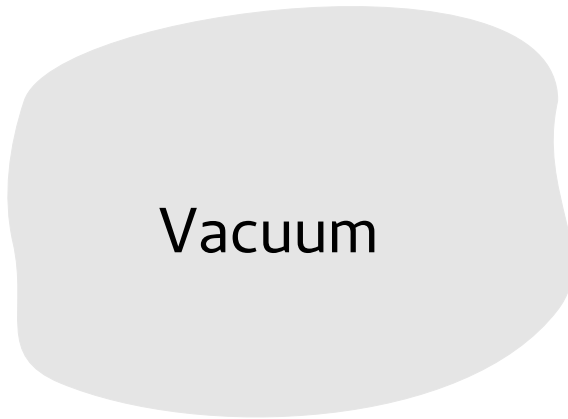
**Casimir physics**, Chameleon model and dark energy

What are the goals?

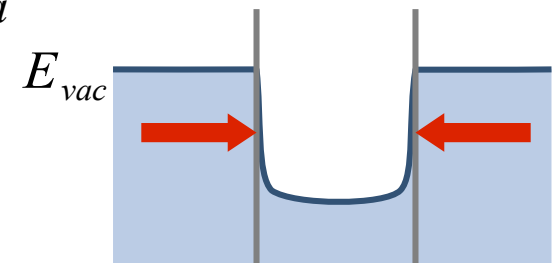


# Casimir force

An intuitive approach ...

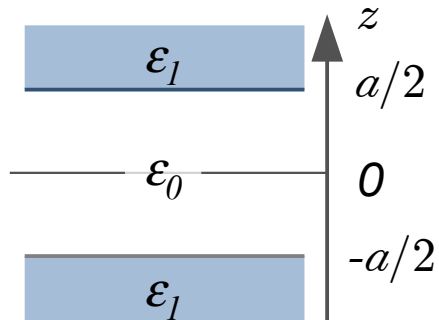


attractive potential:



# Casimir theory (just a little bit)

The realistic approach ...



Electric field:

$$E_\alpha(k, x) = f(k, z)e^{i(k_x x + k_y y - \omega t)}$$

Wave equation:

$$\frac{d^2 f(k, z)}{dz^2} - K^2(k) f(k, z) = 0, \quad K^2(k) = k_x^2 + k_y^2 - \epsilon(\omega) \frac{\omega^2}{c^2}$$

Ansatz:

$$f(k, z) = c_1 e^{Kz} + c_2 e^{-Kz}$$

Boundary conditions at surfaces:

$$\left. \begin{aligned} \epsilon(z) E_\perp(z), E_\parallel(z) &= \text{const.} \\ \frac{\partial E_\perp(z)}{\partial z}, \frac{\partial E_\parallel(z)}{\partial z} &= \text{const.} \end{aligned} \right\} \Rightarrow$$

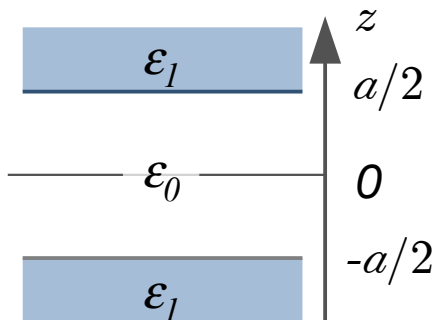
Homogeneous set of eqns.

$$\det C = 0 \equiv \Delta_{\perp, \parallel}(\epsilon, k, \omega, a)$$

(and similar for the magnetic field)

# Casimir theory (a little bit more)

The realistic approach: Lifshitz theory



Sum over modes:

$$\frac{E}{A} = \frac{\hbar}{2} \int_0^\infty \frac{dk_x dk_y}{(2\pi)^2} \underbrace{\sum_{n=0}^\infty \omega_n(a)}$$

Argument  
Theorem:  $\frac{1}{2\pi i} \left[ \int_{-\infty}^\infty \omega d \ln \Delta(\omega) + \int_C \omega d \ln \Delta(\omega) \right]$

infinite

Renormalization: divide  $\Delta(\omega)$  by the expression  $\Delta_\infty(\omega) \equiv \Delta(\omega, a)|_{a \rightarrow \infty}$

$$\frac{E_{\text{ren}}(a)}{A} = \frac{\hbar}{(2\pi)^2 c^2} \int_1^\infty dp \int_0^\infty d\xi p \xi^2 \left[ \ln \frac{\Delta_\perp(i\xi, a)}{\Delta_{\perp, \infty}(i\xi)} + \ln \frac{\Delta_\parallel(i\xi, a)}{\Delta_{\parallel, \infty}(i\xi)} \right], \quad F(a) = -\frac{\partial E_{\text{ren}}(a)}{\partial a}$$

$$\frac{\Delta_\perp(p, i\xi, a)}{\Delta_{\perp, \infty}(p, i\xi)} = 1 - \left( \frac{K_1 \epsilon_0(i\xi) - K_0 \epsilon_1(i\xi)}{K_1 \epsilon_0(i\xi) + K_0 \epsilon_1(i\xi)} \right)^2 e^{-2a \frac{\xi}{c} K_0}, \quad K_j(p, i\xi) = \sqrt{p^2 - 1 + \epsilon_j(i\xi)}$$

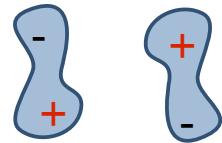
# Casimir vs. van der Waals

Morphology of dipole interactions:

sparse

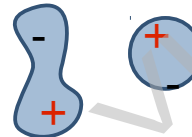
sparse media  $\rightarrow$  single dipole interactions  
described by dipole moments and polarizability  $\alpha$ , **additive**

*van der Waals* interactions



fixed dipoles

*Keesom*



fixed and induced dipoles

*Debye*



induced dipoles

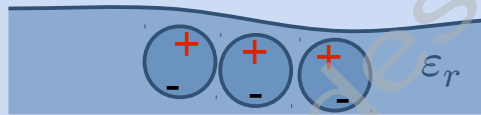
*London*

*Dispersion forces*

Interacting media

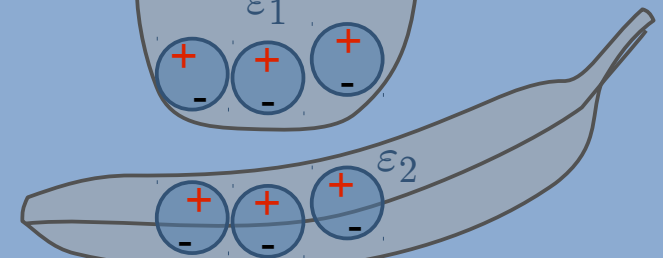
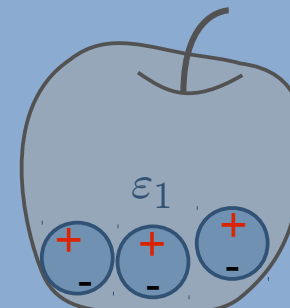
dense

single dipole with dense medium  
*Casimir-Polder* interaction  
 $\rightarrow$  **non-additive** collective response  
depending on the spectral dielectric function  $\epsilon_r$



dense media

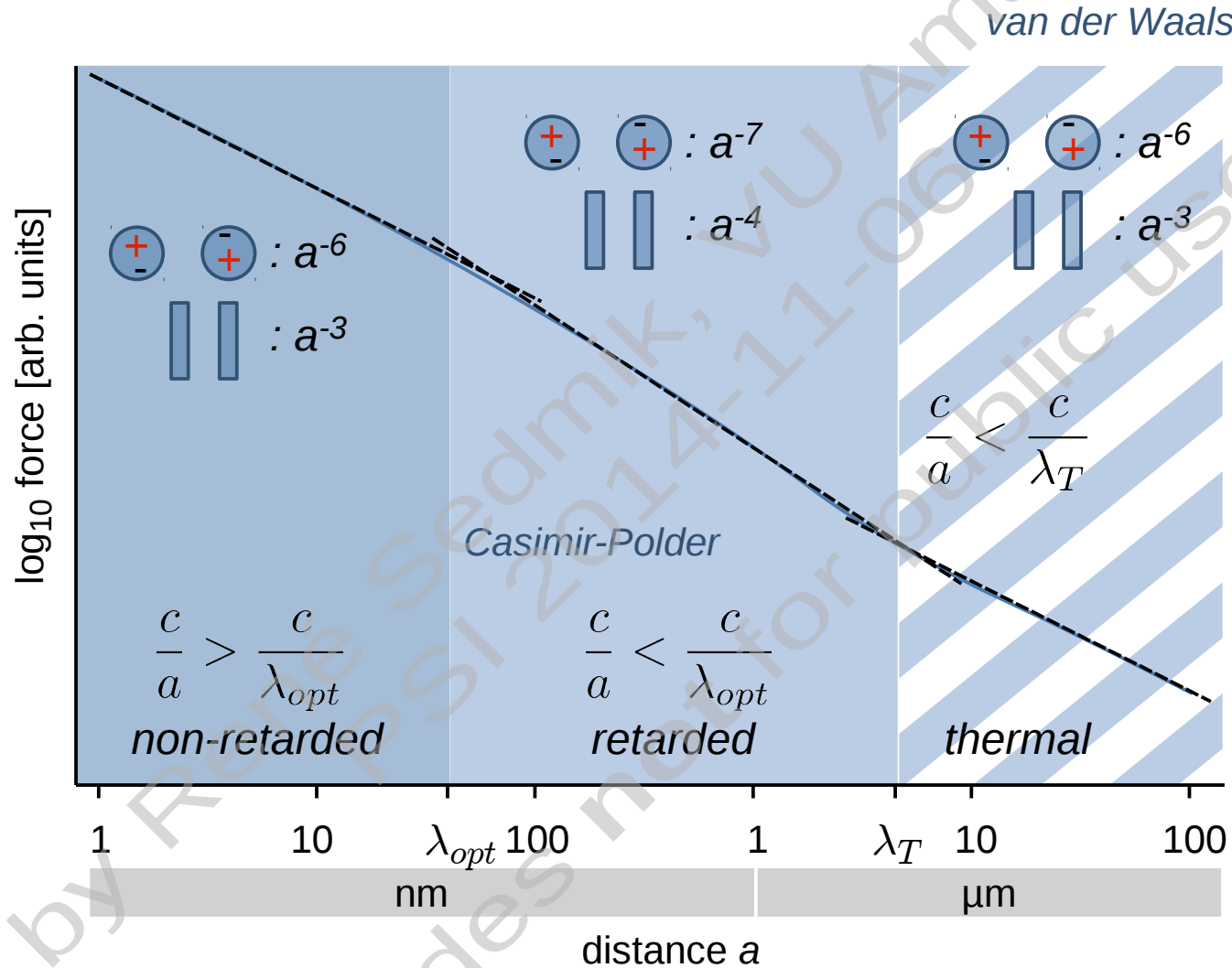
*Casimir* interaction



# Casimir vs. van der Waals

Morphology of dipole interactions:

Interaction distance

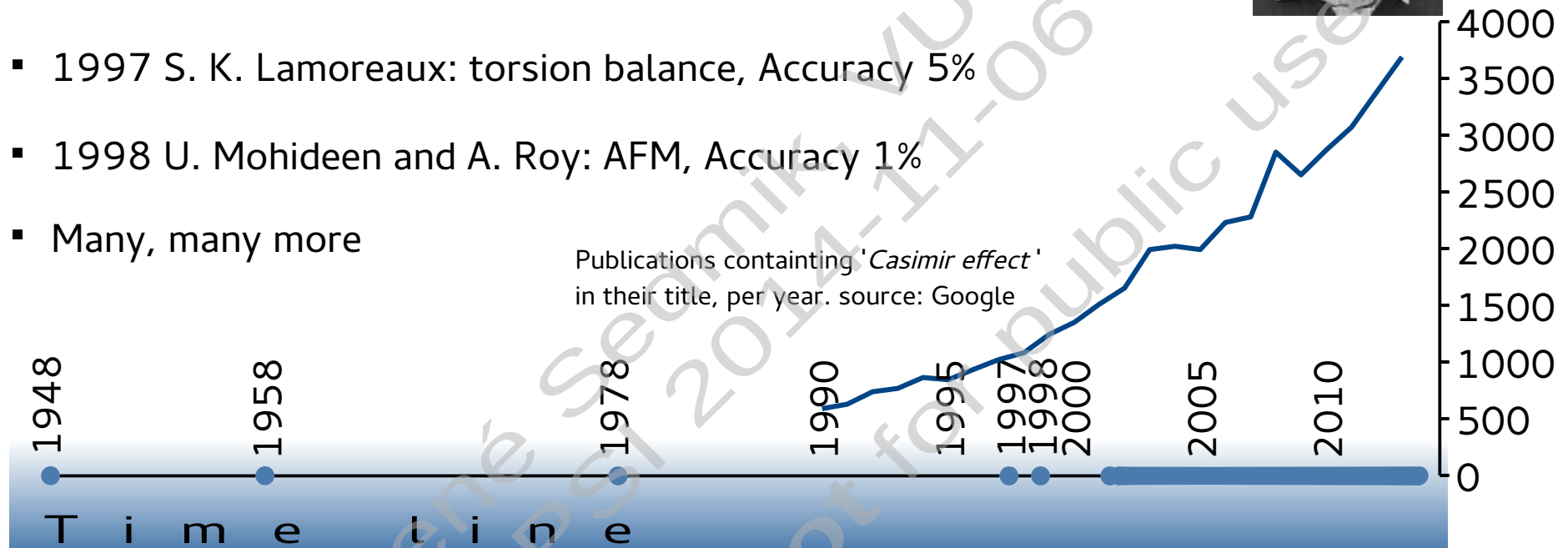


$$\lambda_T = \frac{\hbar c}{k_B T}$$

# Casimir force – a long history

## Landmark contributions

- 1948 H.B.G. Casimir  $\frac{F(a)}{A} = -\frac{\pi^2 \hbar c}{240a^4}$
- 1958 - 1978 Sparnaay, Blokland and Overboek, qualitative
- 1997 S. K. Lamoreaux: torsion balance, Accuracy 5%
- 1998 U. Mohideen and A. Roy: AFM, Accuracy 1%
- Many, many more



H. B. G. Casimir, "On the attraction between two perfectly conducting plates", *Proc. Royal Ned. Acad. Sci.*, Vol. **51**, pp. 793–795 (1948).

M. J. Spaarnay: "Measurement of attractive forces between flat plates", *Physica*, Vol **24**, pp. 751 (1958)

P. H. G. M van Blokland and J. T. G. Overbeek "The measurement of the van der Waals dispersion forces in the range 1.5 to 130 nm", *J. Chem. Soc. Faraday Trans.*, **174**, pp. 2637(1978)

S. K. Lamoreaux, "Demonstration of the Casimir Force in the 0.6 to 6  $\mu\text{m}$  Range", *Phys. Rev. Lett.* **78**, 5–8 (1997)

U. Mohideen and Anushree Roy, "Precision Measurement of the Casimir Force from 0.1 to 0.9  $\mu\text{m}$ ", *Phys. Rev. Lett.* **81**, 004549 (1998)

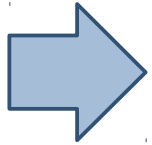
Over 3600 articles in 2013  
Steadily growing interest in the field

# Casimir force – why is it interesting?

**Q:** Why is there such a huge interest in the Casimir effect ?

**A:** Because ...

- it is one of very few **quantum effects with macroscopic action**.
- **it has applications**.
- it becomes the **dominating force** at very small surface separations ( $<100$  nm).
- the Casimir force causes 'stiction' and therefore...

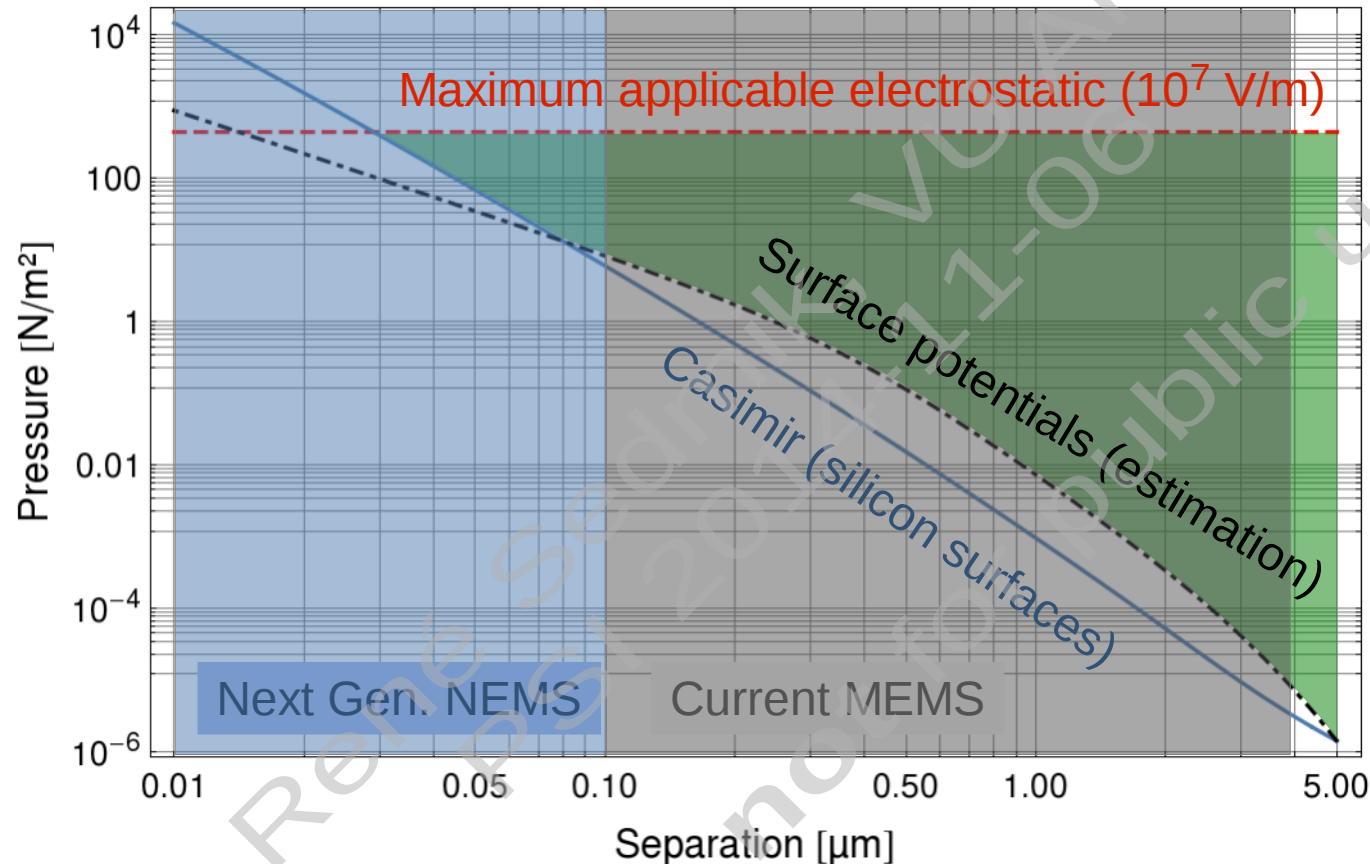


**it is a nuisance for the development of (M)NEMS**

# Casimir force – why is it interesting?

## Micro-electromechanical devices (MEMS)

### Comparison of amplitudes:



### Implications:

- Need to find ways to **eliminate/reduce the Casimir force**
- Find a way to **use the force** to drive actuators

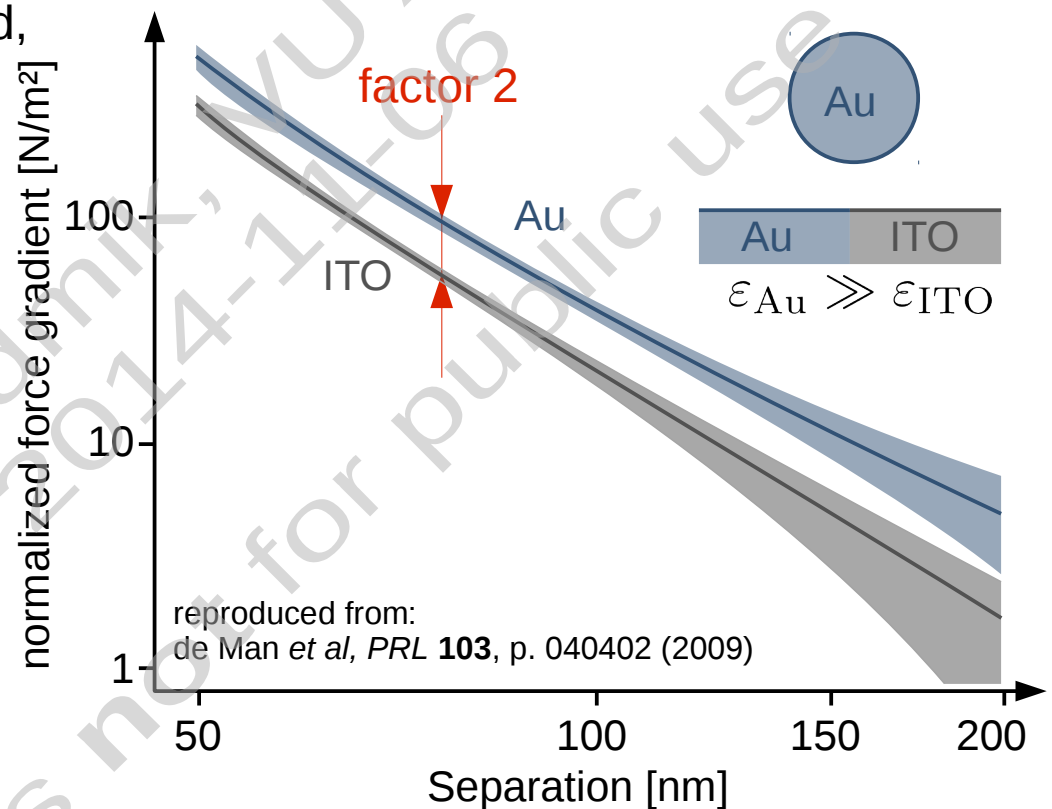
# Casimir force – what has been done already?

## Experiments

... what influences the Casimir force?

Dependence on dielectric functions  $\epsilon_r$

- Range for tunability in vacuum limited, narrow spectral range of differences



$$\frac{E_{\text{reg}}(a)}{A} \propto \int_1^\infty dp \int_0^\infty d\xi p \xi^2 \left[ \ln \left( 1 - \left( \frac{K_1 \epsilon_0(i\xi) - K_0 \epsilon_1(i\xi)}{K_1 \epsilon_0(i\xi) + K_0 \epsilon_1(i\xi)} \right)^2 e^{-2a \frac{\xi}{c} K_0} \right) + \dots \right],$$

# Casimir force – what has been done already?

## Experiments

... what influences the Casimir force?

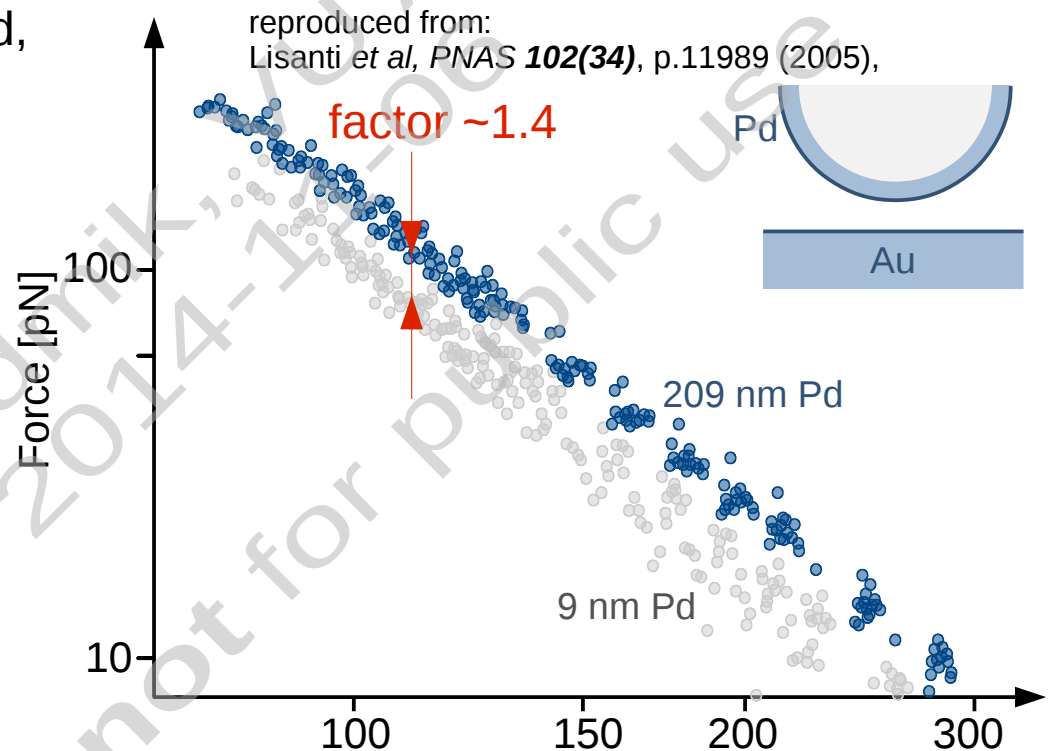
Dependence on dielectric functions  $\epsilon_r$

- ▶ Range for tunability in vacuum limited, **narrow spectral range of differences**

Skin-depth effect

- ▶ materials respond only as bulk if thicker than the **penetration depth**  $\delta_p = \frac{c}{\omega_p} \sim 100 \text{ nm}$

$$\text{with } \omega_p^2 = \frac{4\pi n_e e^2}{m^*}$$



# Casimir force – what has been done already?

## Experiments

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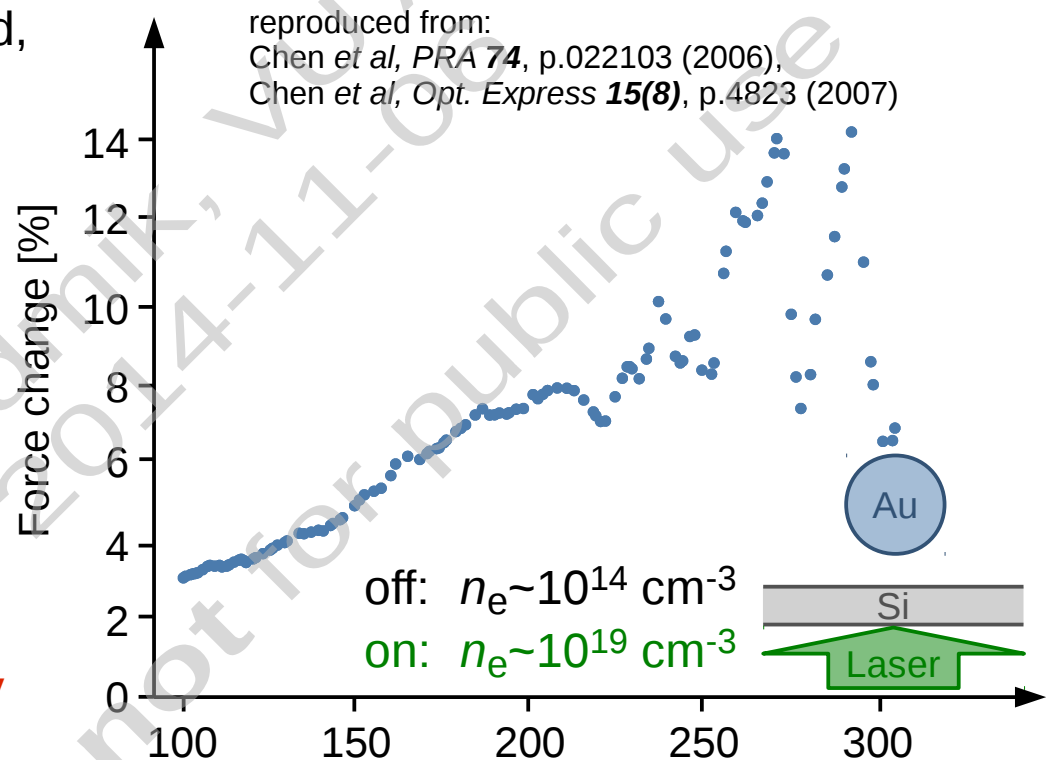
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Modulation of Casimir forces

- ▶ possible but limited again by **narrow spectral range of conductivity**



# Casimir force – what has been done already?

## Experiments

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### Modulation of Casimir forces

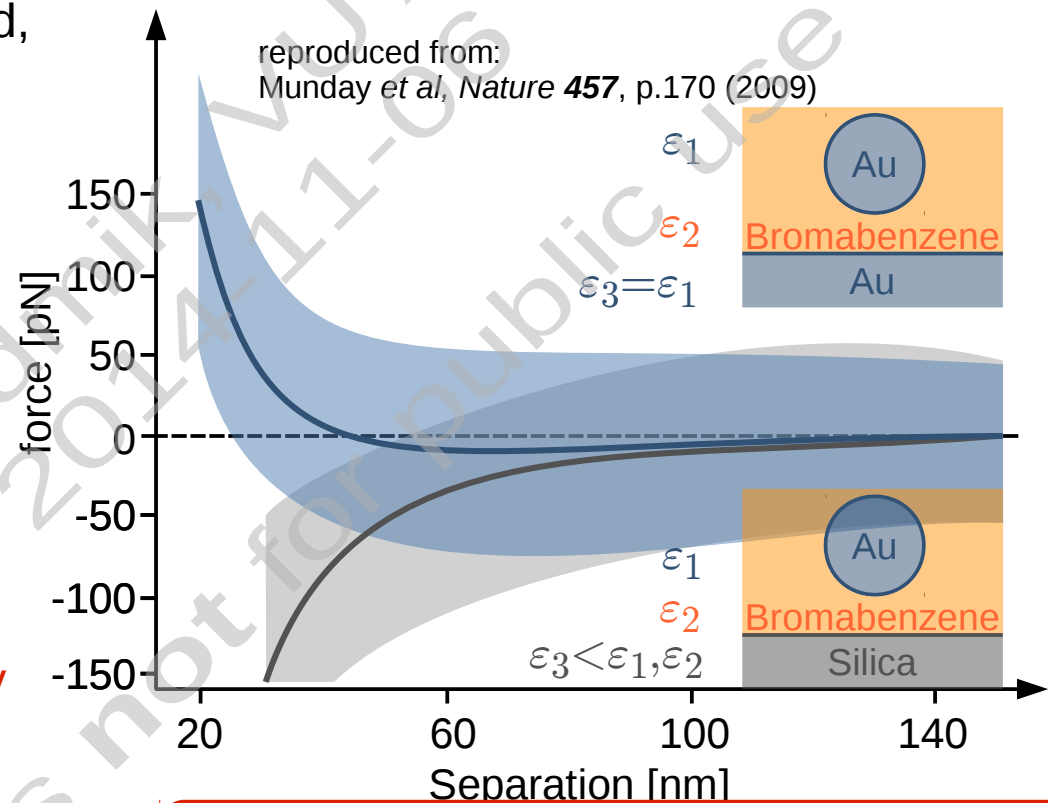
- ▶ possible but limited again by **narrow spectral range of conductivity**

### Repulsive Casimir forces

- ▶ possible using liquids  $\epsilon_1 < \epsilon_2 < \epsilon_3$

$$F(a) \propto -(\epsilon_1 - \epsilon_2)(\epsilon_3 - \epsilon_2)$$

Roughness, temperature, magnetic materials, torques, metamaterials, ...



**No obvious solution  
for NEMS yet.**

# Casimir physics – what problems are we after?

More physical problems in Casimir physics:

## 1. The infamous Drude vs. plasma debate

Review: Klimchitskaya *et al*, Int. J. Mod. Phys. Conf. Ser. **3**(2011), 515

- The Casimir force  $F_C$  depends on dielectric properties via  $\epsilon_r(\omega)$
- There are two main models to extrapolate experimental data on  $\epsilon_r(\omega)$

Drude

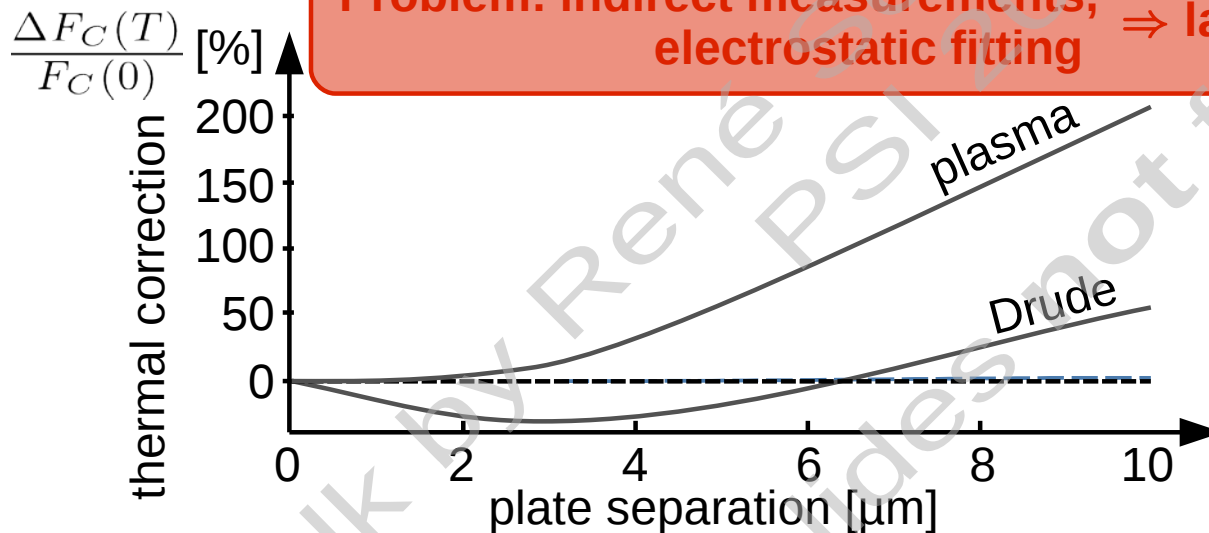
$$\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$$

plasma

$$\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$$

$\omega_p$  plasma frequency  
 $\gamma$  relaxation frequency (dissipation)

**Problem: indirect measurements, electrostatic fitting  $\Rightarrow$  large uncertainties**



An absolute measurement with  $\mathcal{O}(1\%)$  accuracy at  $d > 3 \mu\text{m}$  is required

M.. Bordag *et al* "Advances in the Casimir effect", Oxford Science Publications, (2009).

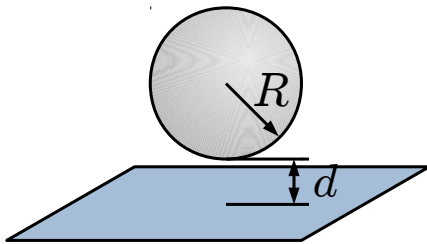
# Casimir physics – what problems are we after?

More physical problems in Casimir physics:

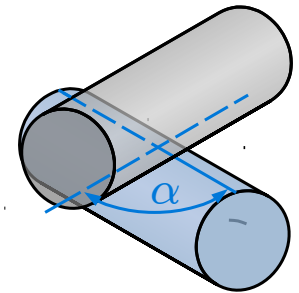
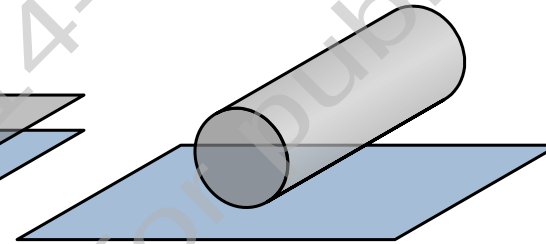
## 2. Geometries without symmetry

- Several numerical approaches for precise computation of forces need experimental verification
- Data on different geometries, some involving parallelism, are required

most experiments:



required data:



Need a precision experiment allowing for **control of relative orientation**

# Casimir physics – what do we need?

More specifically...

- Sensitivity: Drude vs. Plasma problem  
measurement at  $d > 3 \mu\text{m}$  with **accuracy**  $< 1 \%$  (1 pN/cm<sup>2</sup>)  
Sushkov *et al*, Nature Physics **7** (2011), 230, Klimchitskaya *et al*, arXiv:1108.5696 (2011)
- Geometry: Parallelism (plates, cylinders, ...)  
measurement at  $d > 3 \mu\text{m}$  with **max. deviation**  $0.1 \mu\text{rad}$   
Rodriguez-Lopez Emig Phys. Rev. A **85** (2012), 032510, Bimonte *et al*, EPL **97** (2012), 50001

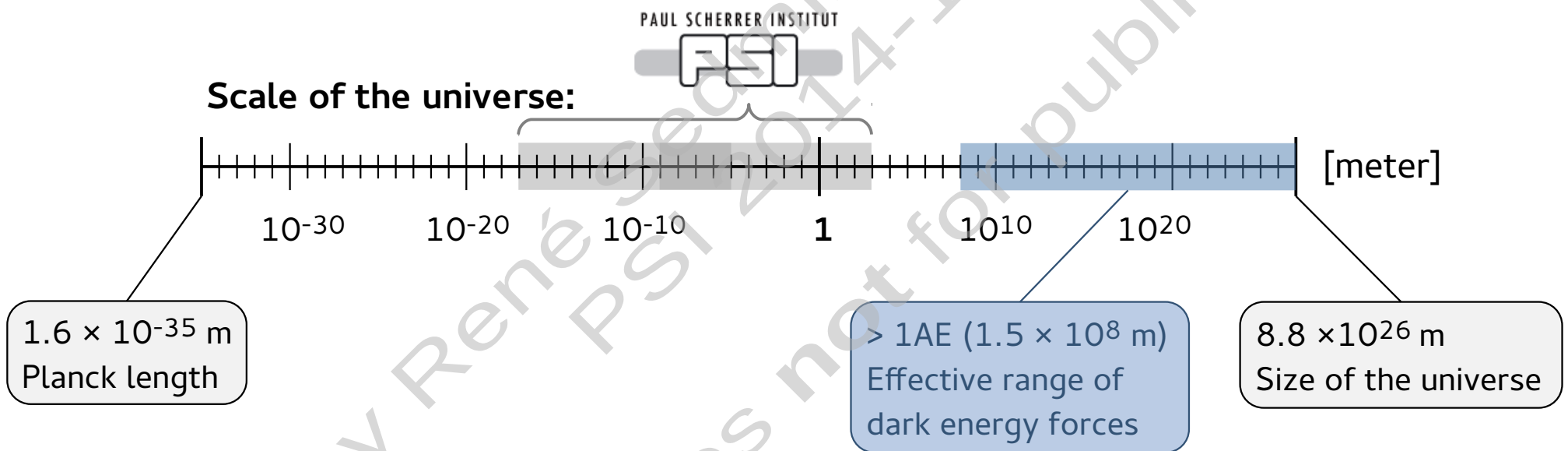
**Need a high-precision, high-accuracy experiment  
with control of orientation (parallelism)**

# Outline

## Part I: Scientific background

Casimir physics, **Chameleon model** and dark energy

What are the goals?



$$\begin{aligned}
\mathcal{L}_{\text{SM}} = & -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{8}\text{Tr}(\mathbf{W}_{\mu\nu}\mathbf{W}^{\mu\nu}) - \frac{1}{2}\text{Tr}(\mathbf{G}_{\mu\nu}\mathbf{G}^{\mu\nu}) \\
& + (\bar{\nu}_L, \bar{e}_L)\tilde{\sigma}^\mu iD_\mu \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} + \bar{e}_R\sigma^\mu iD_\mu e_R + \bar{\nu}_R\sigma^\mu iD_\mu \nu_R + \text{h.c.} \\
& - \frac{\sqrt{2}}{v} \left[ (\bar{\nu}_L, \bar{e}_L)\phi M^e e_R + \bar{e}_R\bar{M}^e \bar{\phi} \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} + (-\bar{e}_L, \bar{\nu}_L)\phi^* M^e \nu_R + \bar{\nu}_R\bar{M}^e \phi^\top \begin{pmatrix} -e_L \\ \nu_L \end{pmatrix} \right] \\
& + (\bar{u}_L, \bar{d}_L)\tilde{\sigma}^\mu iD_\mu \begin{pmatrix} u_L \\ d_L \end{pmatrix} + \bar{u}_R\sigma^\mu iD_\mu u_R + \bar{d}_R\sigma^\mu iD_\mu d_R + \text{h.c.} \\
& - \frac{\sqrt{2}}{v} \left[ (\bar{u}_L, \bar{d}_L)\phi M^d d_R + \bar{d}_R\bar{M}^d \bar{\phi} \begin{pmatrix} u_L \\ d_L \end{pmatrix} + (-\bar{d}_L, \bar{u}_L)\phi^* M^d u_R + \bar{u}_R\bar{M}^d \phi^\top \begin{pmatrix} -d_L \\ u_L \end{pmatrix} \right] \\
& + \overline{(D_\mu\phi)}D^\mu\phi - \frac{m_h^2}{2v^2} \left[ \bar{\phi}\phi - \frac{v^2}{2} \right]^2
\end{aligned}$$

## dark matter

indication: grav. pull  
theory: unknown  
contents: unknown  
part: 27%



## visible universe

indication: em radiation, gravity  
theory:  $\mathcal{L}_{SM}$   
contents: particles of the SM  
part: 5%

## dark energy

indication: acceleration  
theory: unknown  
contents: unknown  
part: 68%

## proposition:

add a new scalar  
Yukawa interaction:  $\mathcal{L}_C$

J. Khoury and A. Weltman,  
*Phys. Rev. Lett.* **93** (2004)  
171104

talk by Renee Sedmik, VU Amsterdam  
PSI 2014-11-06  
slides not for public use

$$\mathcal{L}_{SM} \rightarrow \mathcal{L}'_{SM} + \mathcal{L}_C$$

Einstein-Hilbert action with added dynamics:

$$\frac{e^{\frac{\beta\gamma}{M_{Pl}}\phi}}{4} F^2 + \sum_i \mathcal{L}^{(i)} \left( e^{\frac{\beta_i}{M_{Pl}}\phi}, \psi^{(i)} \right) \rightarrow \sqrt{-g} \left[ \mathcal{R} \frac{M_{Pl}^2}{2} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$

equations of motion for  $\phi$ :

$$\partial^2 \phi = \frac{\partial V(\phi)}{\partial \phi} + \frac{\beta}{M_{Pl}} \rho e^{\frac{\beta}{M_{Pl}}\phi}$$

effective potential

$$V_{eff} = V(\phi) + \rho e^{\frac{\beta}{M_{Pl}}\phi}$$

depend on the local mass density  $\rho$

$\mathcal{L}_{SM}$

$\mathcal{L}_C$

talk by René Sedmik, VU Amsterdam  
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 2014-11-06

$$m_\phi(\rho) = \sqrt{\frac{\partial^2 V_{eff}(\phi)}{\partial \phi^2}} \Big|_{\phi=\phi_{min}(\rho)}$$

$$\propto \frac{\beta}{M_{Pl}} \rho$$

Chameleon potential:

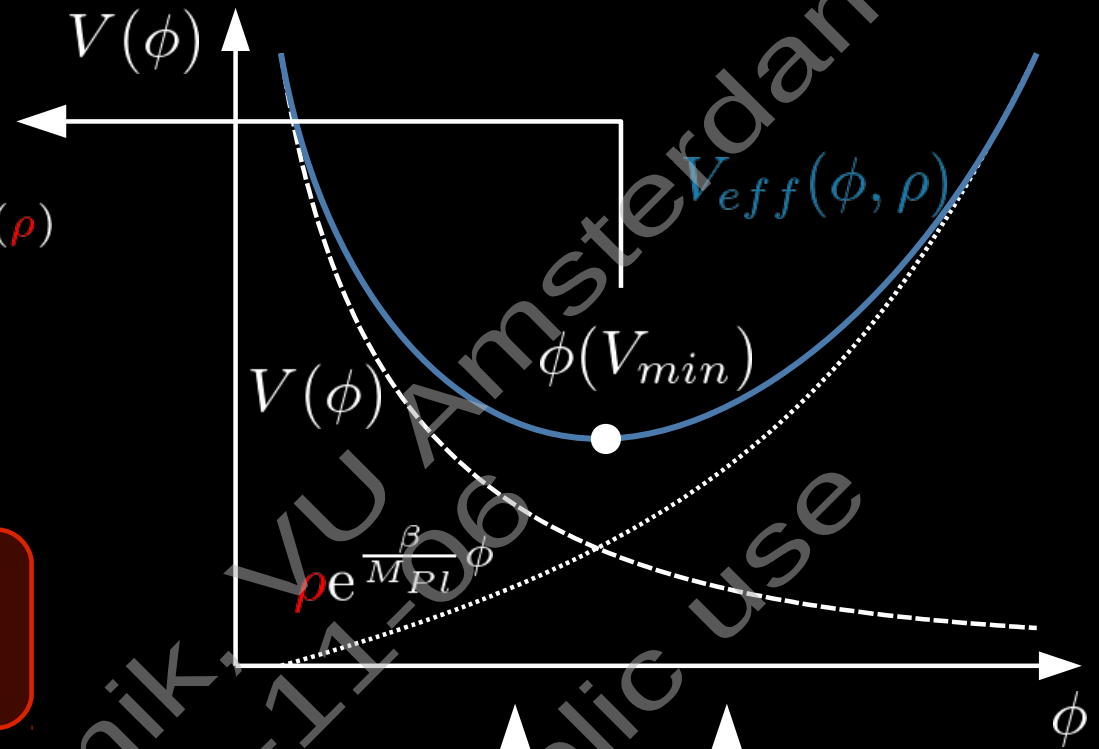
$$V(r) \propto \beta \frac{M}{M_{Pl}} \frac{e^{-m_\phi(\rho)r}}{r}$$

$$\partial^2 \phi = \frac{\partial V(\phi)}{\partial \phi} + \frac{\beta}{M_{Pl}} \rho e^{\frac{\beta}{M_{Pl}} \phi}$$

equations of motion for  $\phi$ :

$$V_{eff} = V(\phi) + \rho e^{\frac{\beta}{M_{Pl}} \phi}$$

effective potential



'quintessence'  
runaway potential

Yukawa  
coupling

**Q:** What does that mean?

free space:

$$\rho \ll 1 \Rightarrow m_\phi \ll 1$$

strong interaction

$$V(r) \propto \beta \frac{M}{M_{Pl}} \frac{e^{-m_\phi(\rho)r}}{r}$$

large massive object:

$$\rho \gg 1 \Rightarrow m_\phi \gg 1$$

weak interaction



For this adaptivity, the model is dubbed **Chameleon**

**A:** The interaction strength is inversely dependent on the local mass (energy) density.

**Q:** What does that mean?

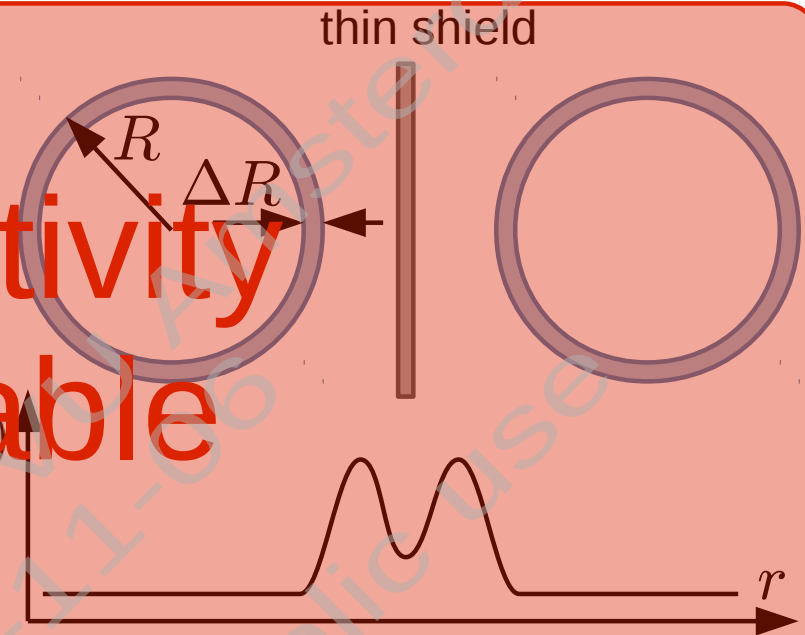
# How can we detect the Chameleon on Earth?

## Idea #1:

Precision Cavendish-type 5<sup>th</sup> force experiment  
Measure force as a function of distance

$$V(r) \propto \beta \frac{M}{M_{Pl}} \frac{\Delta R}{R} \frac{e^{-m_\phi(\rho)r}}{r}$$

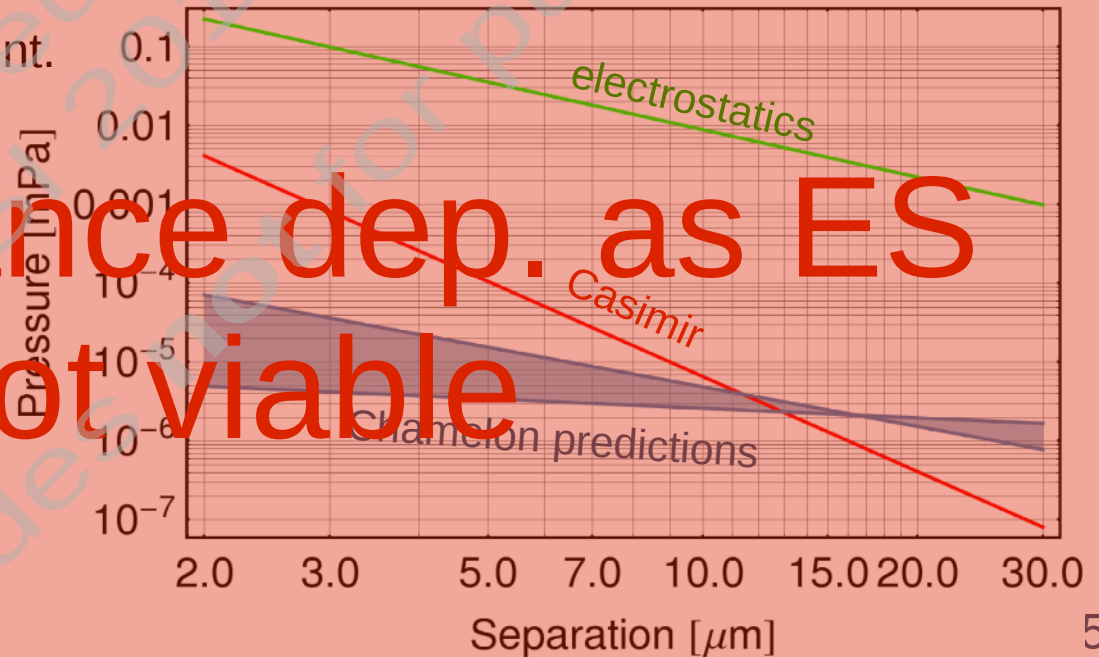
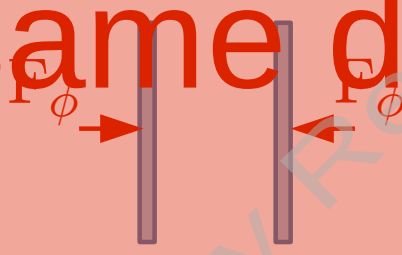
thin shell effect



low sensitivity  
not viable

## Idea #2:

Precision Casimir-type force experiment.  
Measure force as a function of distance without shield



same distance dep. as ES  
not viable

de Man, Heeck, Wijngarden, Iannuzzi,  
PRL. **103** (2009), 040402

# How can we detect the Chameleon on Earth?

## Idea #3:

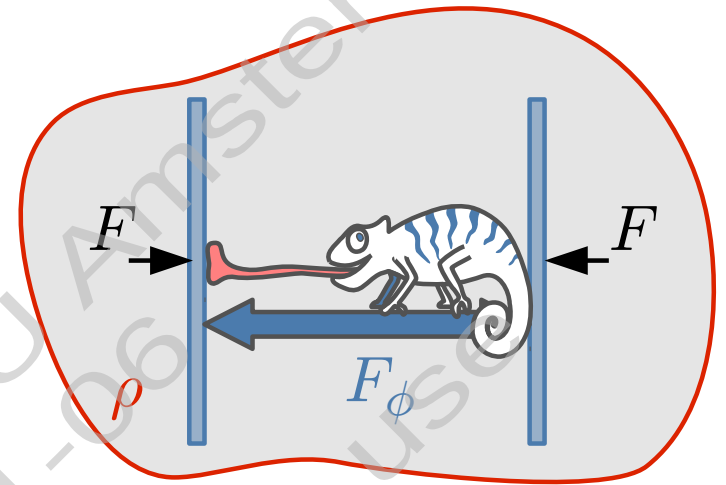
Measure at **constant plate separation**

the change in the force for **different gas density  $\rho$**

Brax, van de Bruck, Davis, Shaw, Iannuzzi,  
*Phys. Rev. Lett.* **104** (2010) 241101

$$F(\rho) = F_{ES} + F_C + F_G + F_\phi$$

$$\Delta F(\rho) = F(\rho) - F(0)$$



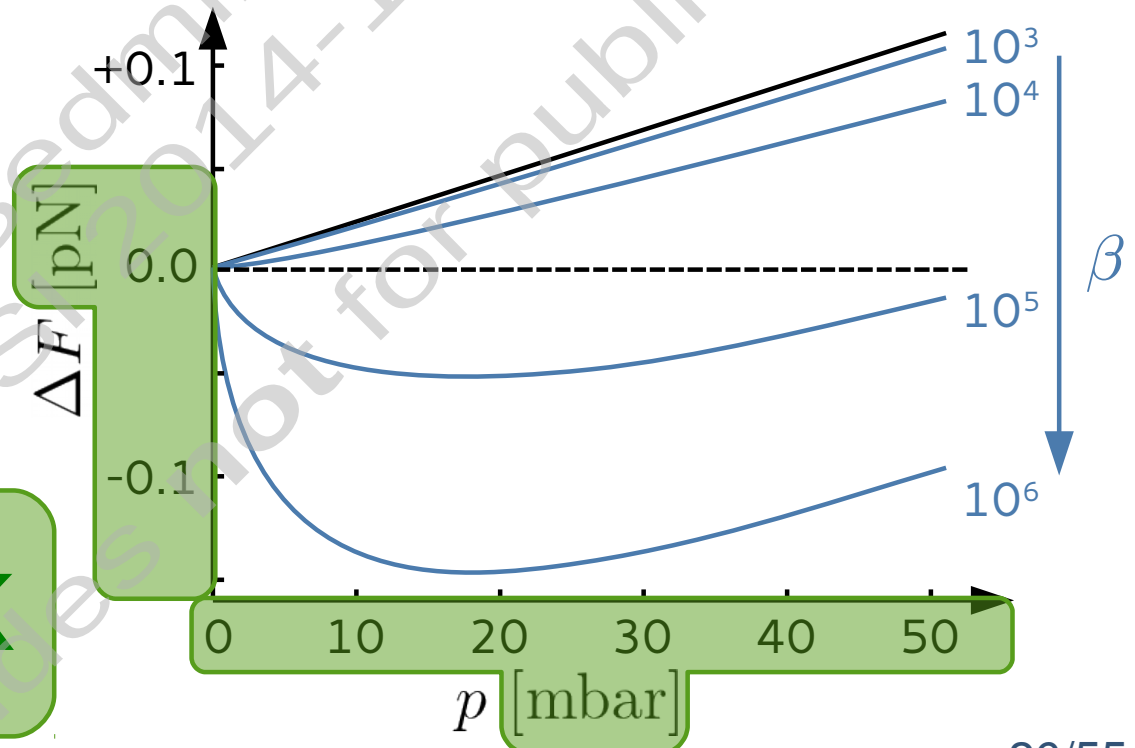
## Assumptions:

- Xe gas
- plate area 1 cm<sup>2</sup>

$$V(\phi) = \Lambda^4 \left( 1 + \frac{\Lambda^n}{\phi^n} \right)$$

- $\Lambda = 2.4$  meV

✓ can work

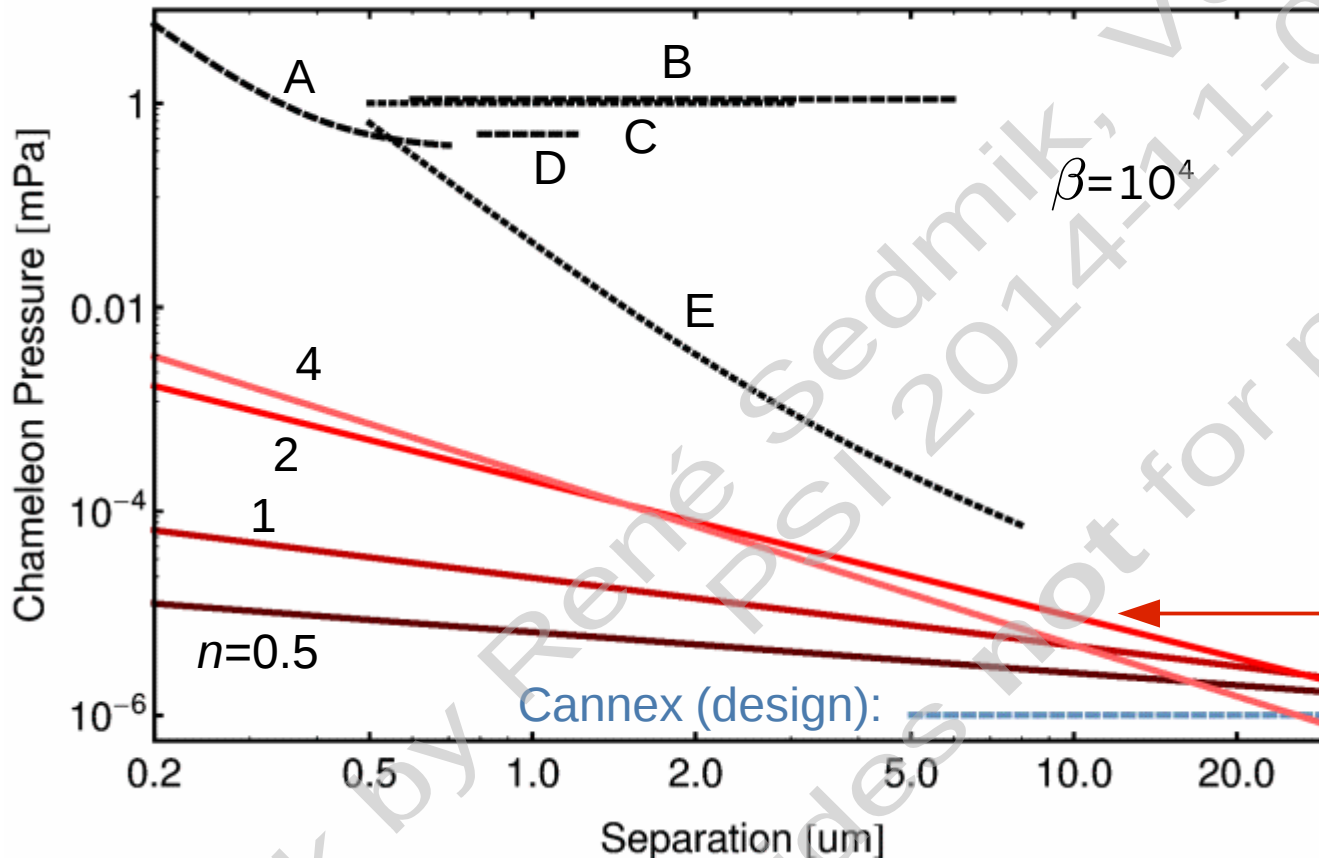


# What can we reach?

## New limits...

- to the interaction parameter  $\beta$
- to the form of potentials  $V(\phi)$ , exclusion of specific  $n$

$$V(\phi) = \Lambda^4 \left( 1 + \frac{\Lambda^n}{\phi^n} \right)$$



- A:** Decca *et al*, Ann. Phys. (N.Y.) **318** (2005), 37
- B:** Lamoreaux PRL. **78**(1997), 5
- C:** Bressi *et al*, PRL. **88** (2002), 041804
- D:** Decca *et al*, PRL. **91** (2003), 050402
- E:** Sushkov *et al*, Nature Phys. **7** (2011), 230

Parameter range not excluded by any experiment

# Chameleon physics – what do we need?

More specifically...

- Sensitivity: New limits can be seen for a measurement at  $d > 10 \mu\text{m}$  with accuracy  $< 0.1 \text{ pN/cm}^2$ .  
Brax *et al*, Phys. Rev. D **76** (2007), 124034
- Geometry: Parallel plates give the maximum force sensitivity. Measurement at  $d > 10 \mu\text{m}$  with max. deviation **0.1  $\mu\text{rad}$** .  
Brax *et al*, Phys. Rev. Lett. **104** (2011), 241101

compare:  
Casimir

$d > 3 \mu\text{m}$ ,  
 $1 \text{ pN/cm}^2$

0.1  $\mu\text{rad}$

Again: Need a high-precision, high-accuracy experiment with control of orientation (parallelism)

Both forces could be measured with the same setup

## CANNEX

Casimir And  
Non-Newtonian force  
EXperiment

**Aim:** Measure Casimir as well as Chamelon forces  
with the same setup

talk by René S. Schmik, VU Amsterdam  
PSI 2014-11-06  
slides not for public use

# Outline

## Part I: Scientific background

Casimir physics, Chameleon model and dark energy

What are the goals?

## Part II: Setup

Core: Principle of measurement, Control of parallelism,

Vibration insulation: Mechanisms, Preliminary characterization

Status

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# What do we need to implement Cannex?

Cooking recipe:  
*sub-pN Force measurement in  
parallel plates geometry*

Sensitivity: **0.1 pN**  
Geometry: flat, parallel plates  
Area: 1 cm<sup>2</sup>  
Separation: 3-30 μm  
Pressure: vacuum – ambient

difficulty level: **challenging**  
preparation time: **3 years**

## Shopping List:

Force transducer for <1 N/m

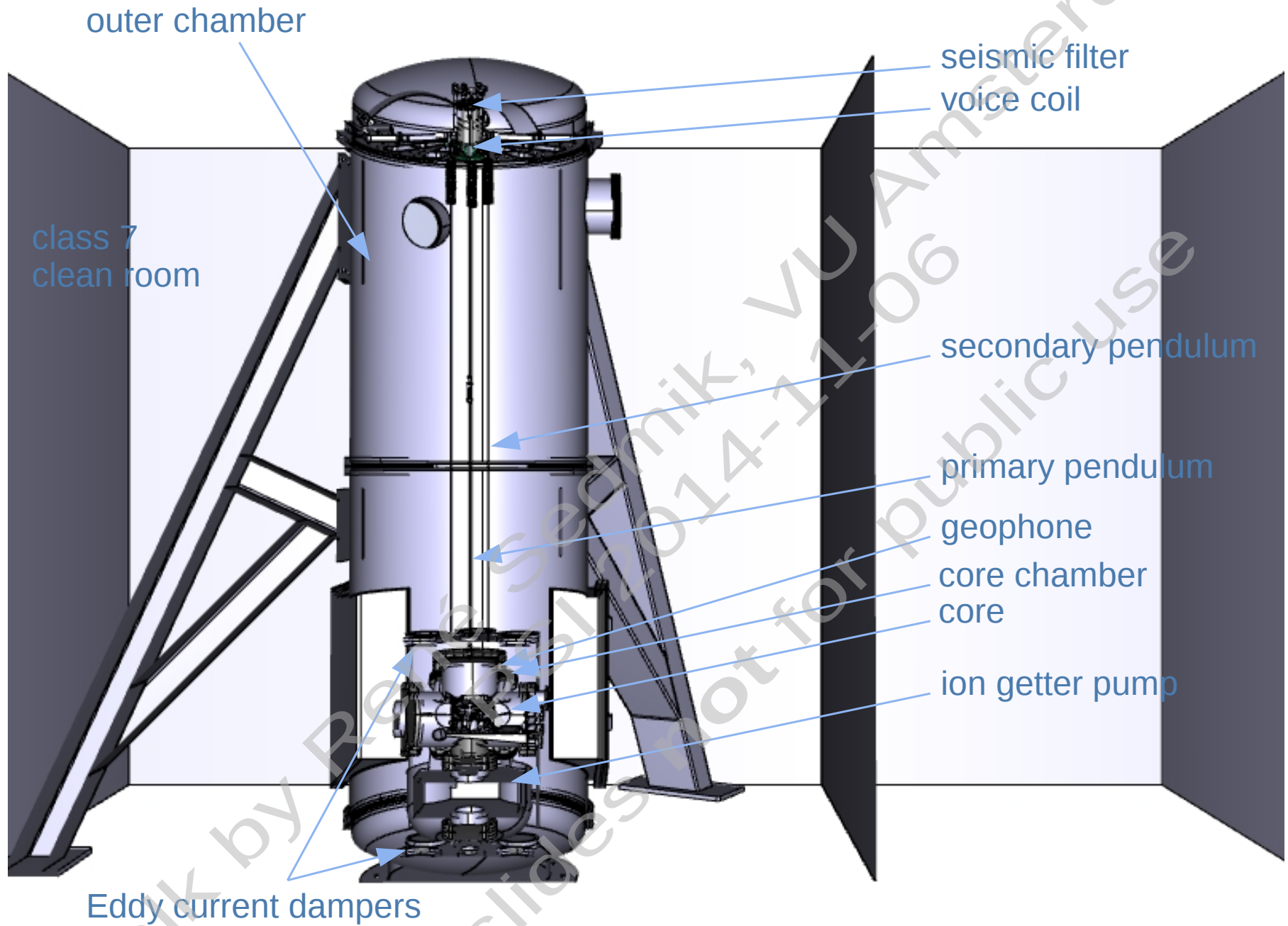
Capacitive detection system:  
(C-bridge 1 aF at nom. 100 pF)

Parallelism control ~0.1 μrad

Ultra-flat surfaces (waviness < 20 nm)

Vibration isolation (-60 dB at 1 Hz)

# Setup



# Outline

## Part I: Scientific background

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## Part II: Setup

**Core: Principle of measurement,** Control of parallelism,  
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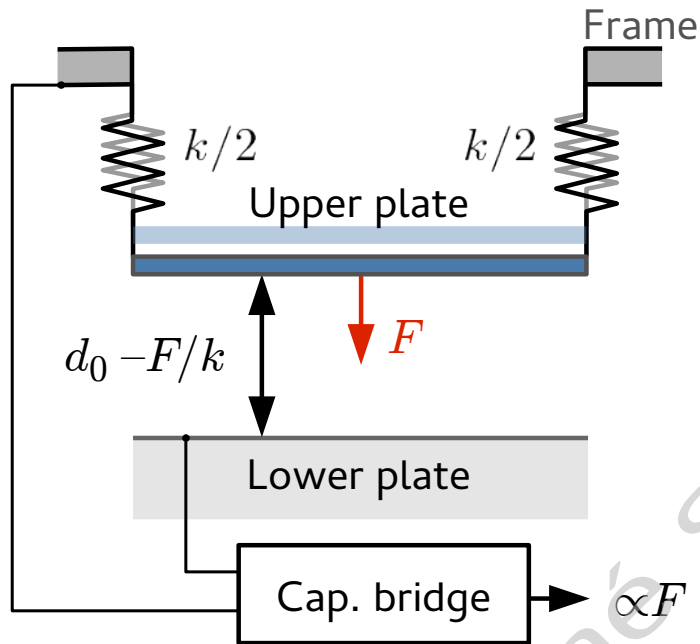
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# Force detection

## Principle:

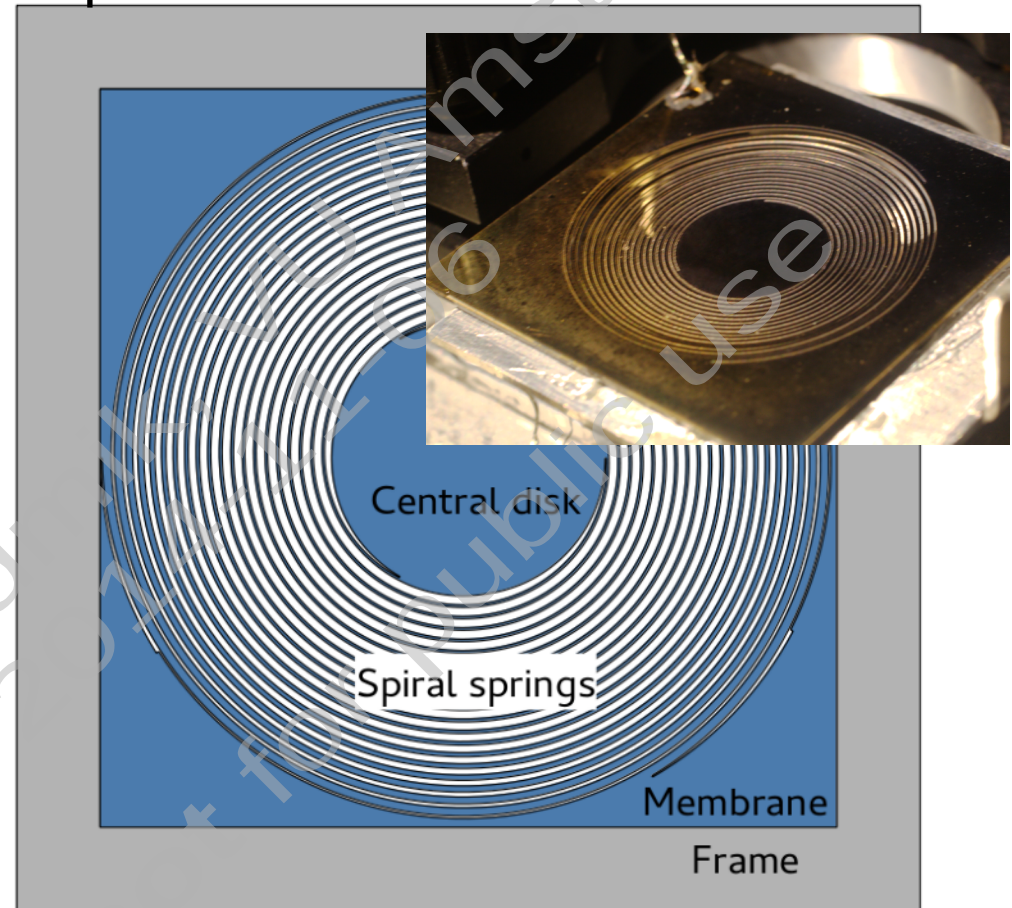
Measure capacitively the displacement of a spring.



Now: Andeen-Hagerling 2700A  
(0.4 ppm resolution  $\sim 1$  pN)

Target: Custom differential bridge  
(0.01 ppm resolution  $< 0.1$  pN)

## Implementation: (patents pending)



Custom-fabricated Silicon membrane

Force constant	0.25 N/m
Eigenfrequency	11.4 Hz
Disk area	1 cm <sup>2</sup>
Waviness(disk)	$< 15$ nm (whole area)

# Force detection

## Principle:

Measure capacitively the displacement of a spring.

## Implementation:

(patents pending)



### Shopping List:

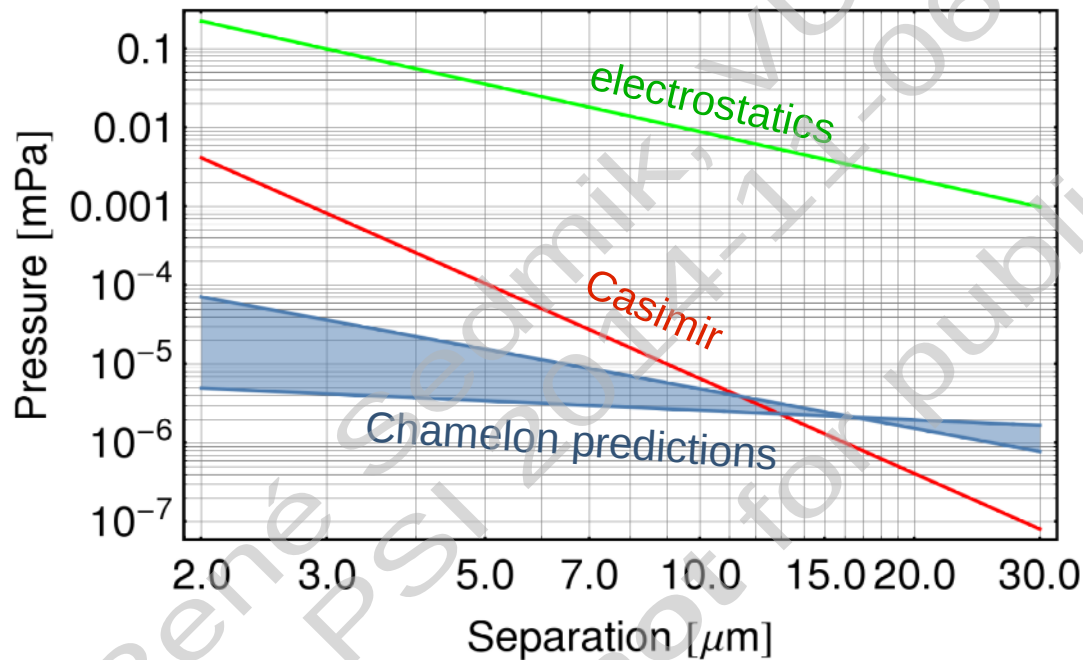
- ✓ Force transducer for  $< 1 \text{ N/m}$   
 $0.25 \text{ N/m}$
- Capacitive detection system:  
(C-bridge  $1 \text{ aF}$  at nom.  $100 \text{ pF}$ )
- Parallelism control  $\sim 0.1 \mu\text{rad}$
- ✓ Ultra-flat surfaces (waviness  $< 20 \text{ nm}$ )
- Vibration isolation ( $-60 \text{ dB}$  at  $1 \text{ Hz}$ )

Custom-fabricated Silicon membrane	
Spring constant	$0.25 \text{ N/m}$
Resonance frequency	$11.4 \text{ Hz}$
Area	$1 \text{ cm}^2$
Waviness (disk)	$< 15 \text{ nm}$ (whole area)

# Force detection: Elimination of parasitic forces

## Problem diagnosis:

- Omnipresent varying Fermi potentials  $V_0$  give rise to **parasitic electrostatic forces**.
- **Electrostatic and hydrodynamic effects** are normally far stronger than **Casimir or Chameleon interactions**.

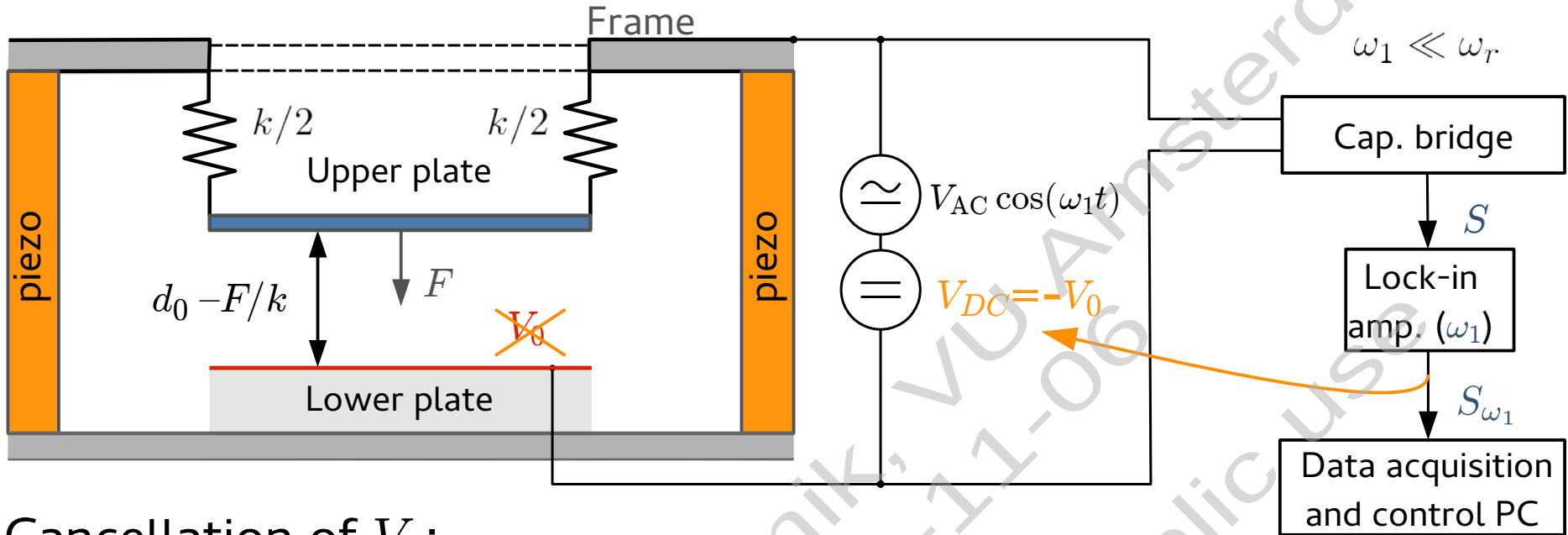


## Solution:

- **Active feedback circuit** to cancel electrostatics
- Use **UHV** ( $10^{-9}$  mbar) for Casimir measurements to cancel hydrodynamics
- Chameleon measurements at **quasi-DC** (<100 mHz) frequencies

de Man, Heeck, Iannuzzi,  
*Phys. Rev. A* **79**, (2009), 024102

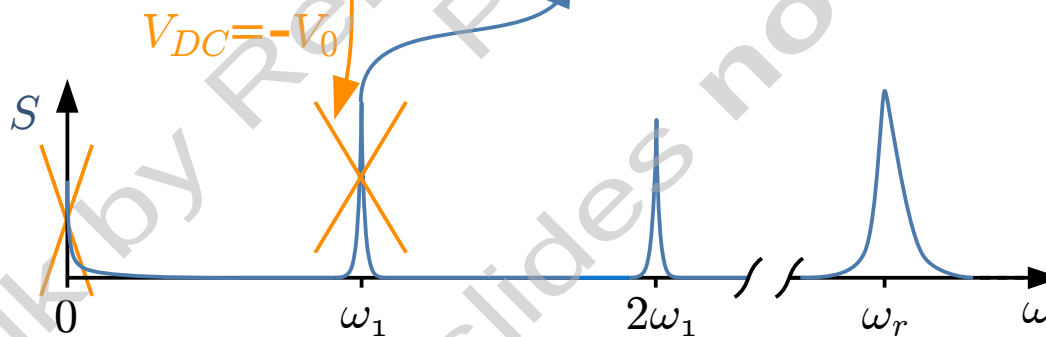
# Force detection – AC measurements



Cancellation of  $V_0$ :

**Signal:**  $S \equiv C(F) \underset{F/k \ll d}{\approx} \frac{\epsilon_0 A}{d - F/k}$       **Force:**  $F = \frac{\epsilon_0 A}{d^2} [\cancel{V_0} + V_{DC} + V_{AC} \cos(\omega_1 t)]^2$

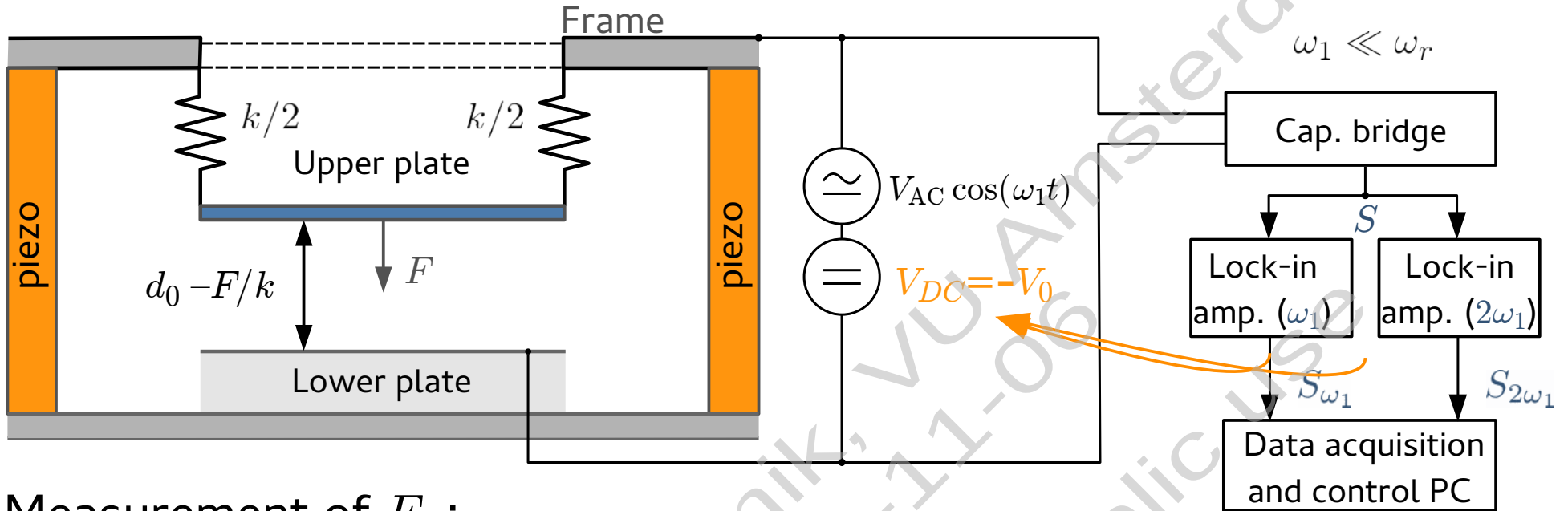
$$= \frac{\epsilon_0^2 A^2}{8d^4 k^2} [2(\cancel{V_0} + V_{DC}) \cos(\omega_1 t) \left( 4k + \frac{\epsilon_0 A (3V_{AC}^2 + 4[V_0 + V_{DC}]^2)}{d^3} \right) + \dots]$$



**Effect:**

- Replace unknown force by small known one.
- Residual potential < 50  $\mu\text{V}$  (compared to initially >10 mV)
- Experience shows that the residual force  $\ll$  Casimir force

# Force detection – AC measurements



Measurement of  $F_C$ :

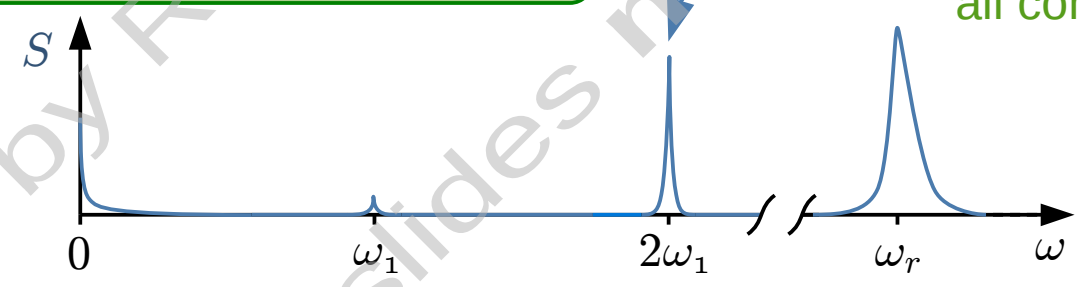
**Signal:**  $S \equiv C(F) \underset{F/k \ll d}{\approx} \frac{\epsilon_0 A}{d - F/k}$

**Force:**  $F = \frac{\epsilon_0 A}{d^2} [V_{AC} \cos(\omega_1 t)]^2 + F_C(d)$

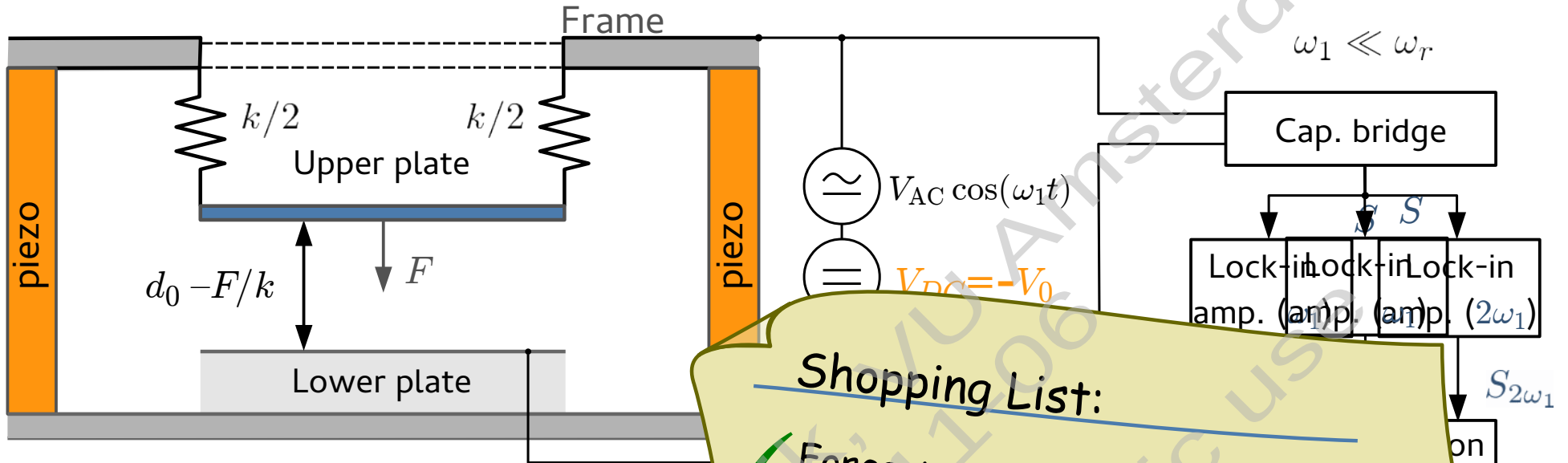
Measurement of  $F_C$  from the  $S_{2\omega_1}$  signal

$$F_C = -\frac{dk}{2} - \frac{\epsilon_0 A}{4d^2} V_{AC}^2 + \frac{2d^5 k^2}{\epsilon_0^2 A^2 V_{AC}^2} S_{2\omega_1}$$

all components known!



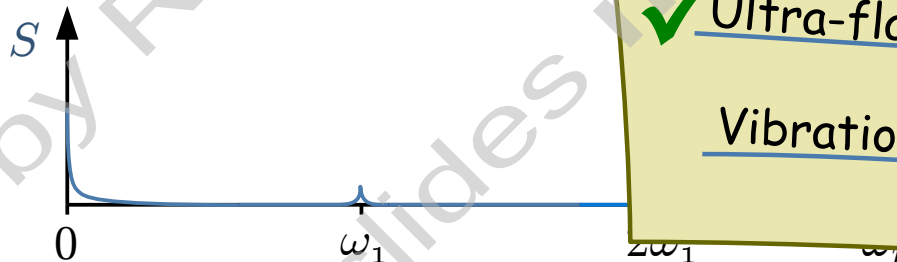
# Force detection – AC measurements



Measurement of  $F_C$ :

Signal:  $S \equiv C(F) \underset{F/k \ll d}{\approx} \frac{\epsilon_0 A}{d - F/k}$

Measurement of  $F_C$   
from the  $S_{2\omega_1}$  signal



Shopping List:

- ✓ Force transducer for  $< 1 \text{ N/m}$   
**0.25 N/m**
- Capacitive detection system:  
(C-bridge ~~1 aF~~ at nom. 100 pF)  
**currently 40 aF, need to improve**
- Parallelism control  $\sim 0.1 \mu\text{rad}$
- ✓ Ultra-flat surfaces (waviness  $< 20 \text{ nm}$ )
- Vibration isolation ( $-60 \text{ dB}$  at 1 Hz)

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**Core:** Principle of measurement, **Control of parallelism,**

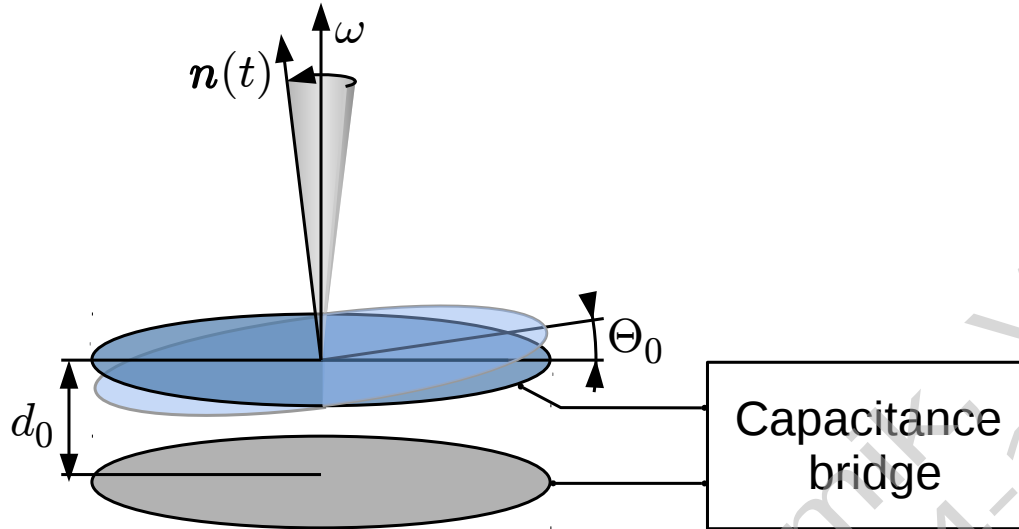
Vibration insulation: Mechanisms, Preliminary characterization

Status

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PSI 2014-11-06  
slides not for public use

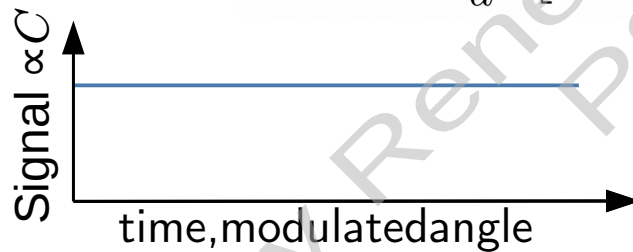
# Parallelism control: Principle

assume: **parallel plates**

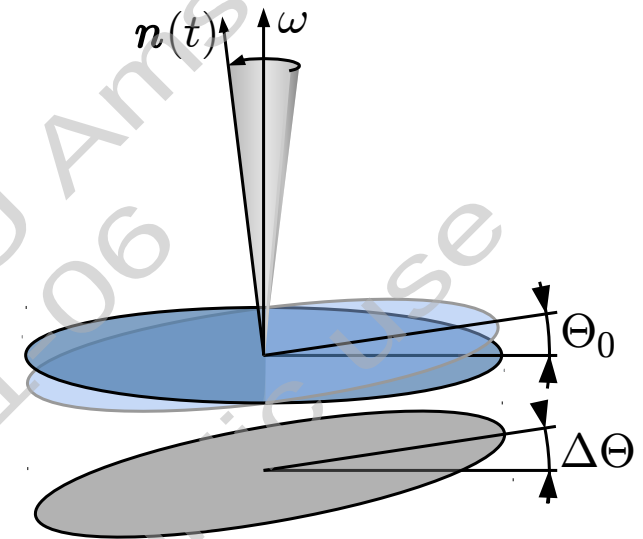


$$\Theta(t) = \Theta_0 = \text{const.}$$

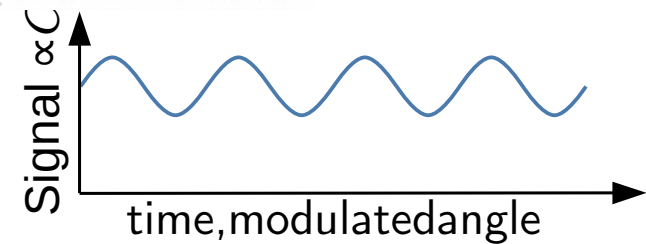
$$C(t) \approx \varepsilon_0 \frac{R^2 \pi}{d} \left[ 1 + \left( \frac{R}{2d} \right)^2 \left( \Theta_0^2 + \Delta\Theta^2 + 2\Theta_0 \Delta\Theta \cos \omega t \right) \right]$$



**plates with relative tilt  $\Delta\Theta$**

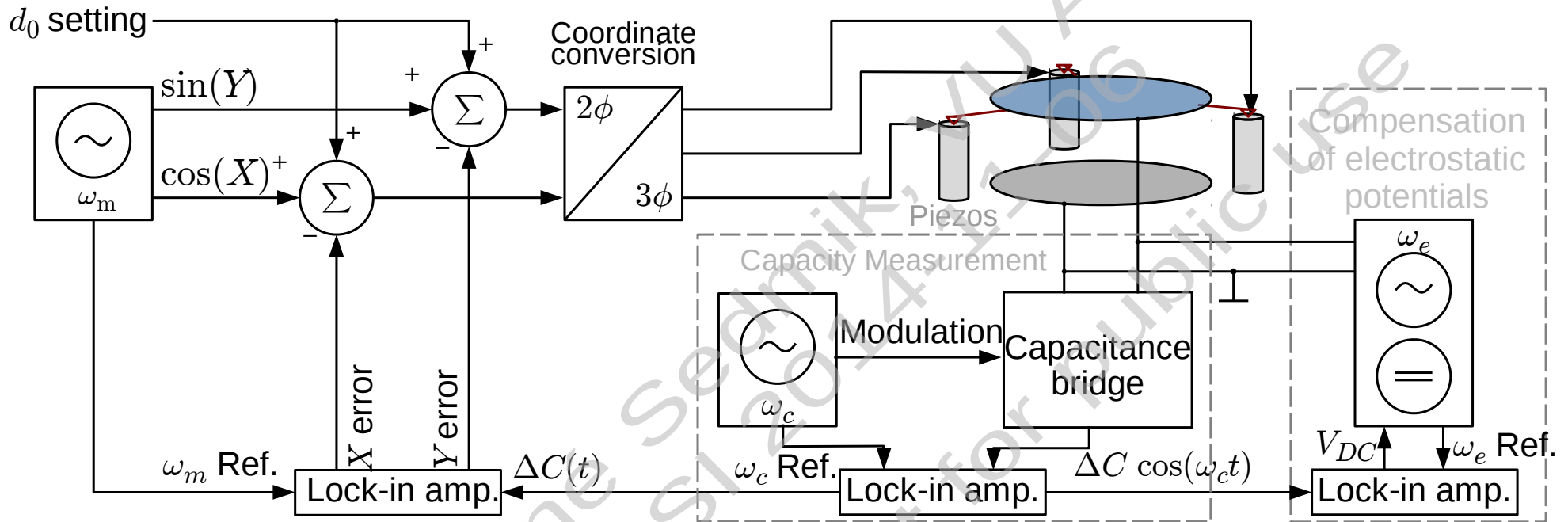


$$\Theta(t) = \Theta_0 + \Delta\Theta \cos \omega t$$



**Idea:** Use feedback circuit to compensate  $\Delta\Theta$

## A feedback circuit



# Parallelism control: Proof of principle

## Proof of principle

### Step response

under very bad conditions

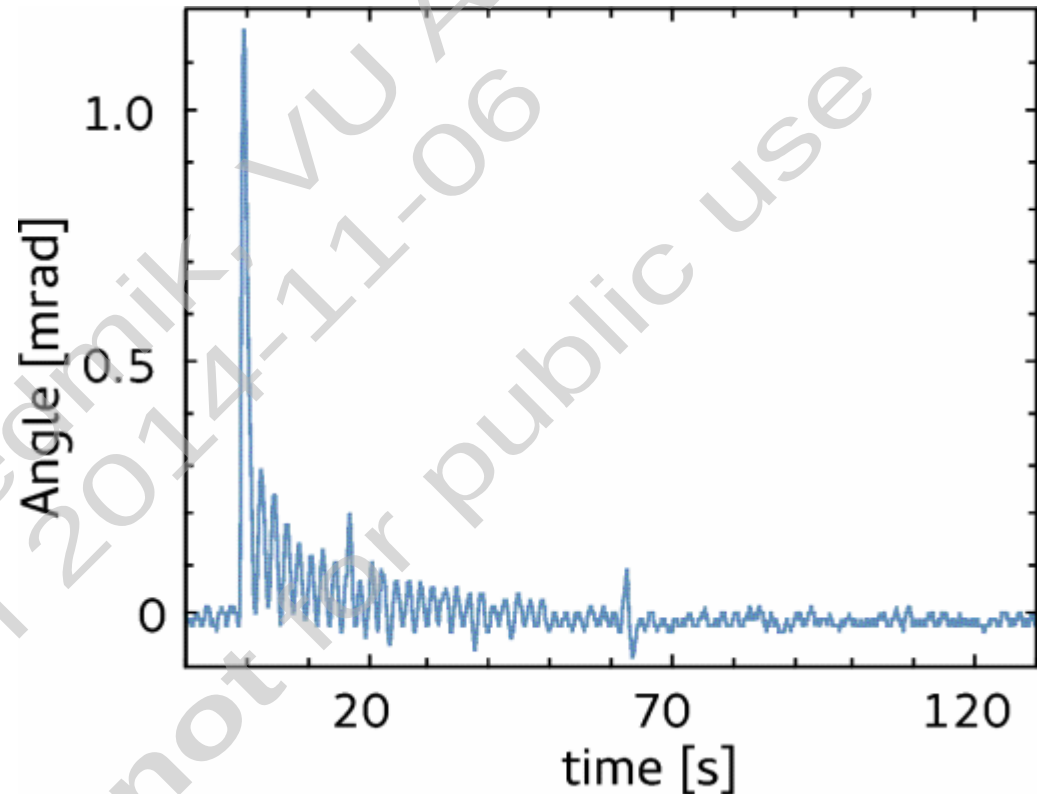
- First test
  - in air
  - without anti-vibration
  - with rough surfaces
  - insufficient shielding
- 6  $\mu\text{m}$  single-sided step
- nominal distance 90  $\mu\text{m}$
- plate area 1  $\text{cm}^2$

### Long-term stability

- Same conditions
- **3  $\mu\text{rad}$ (RMS)**

### Target

- Assumptions: good surface quality, vacuum, anti-vibration
- **0.1  $\mu\text{rad}$  (~1 nm total tilt)**



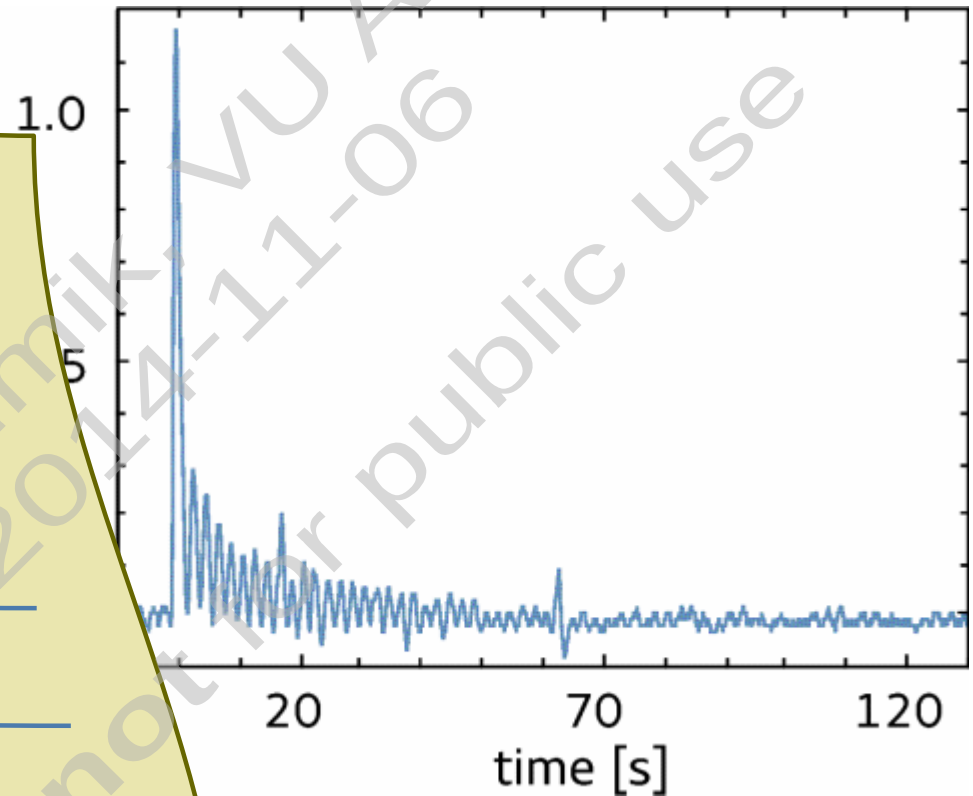
# Parallelism control: Proof of principle

## Proof of principle

### Step response

#### Shopping List:

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**0.25 N/m**
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(C-bridge ~~1 aF~~ at nom. 100 pF)  
**currently 40 aF, need to improve**
- Parallelism control  ~~$\sim 0.1$   $\mu$ rad~~  
**rough test:  $\sim 3$   $\mu$ rad**
- ✓ Ultra-flat surfaces (waviness  $< 20$  nm)
- Vibration isolation ( $-60$  dB at 1 Hz)



anti-vibration

# Outline

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**Vibration insulation:** Mechanisms, Preliminary characterization

Status

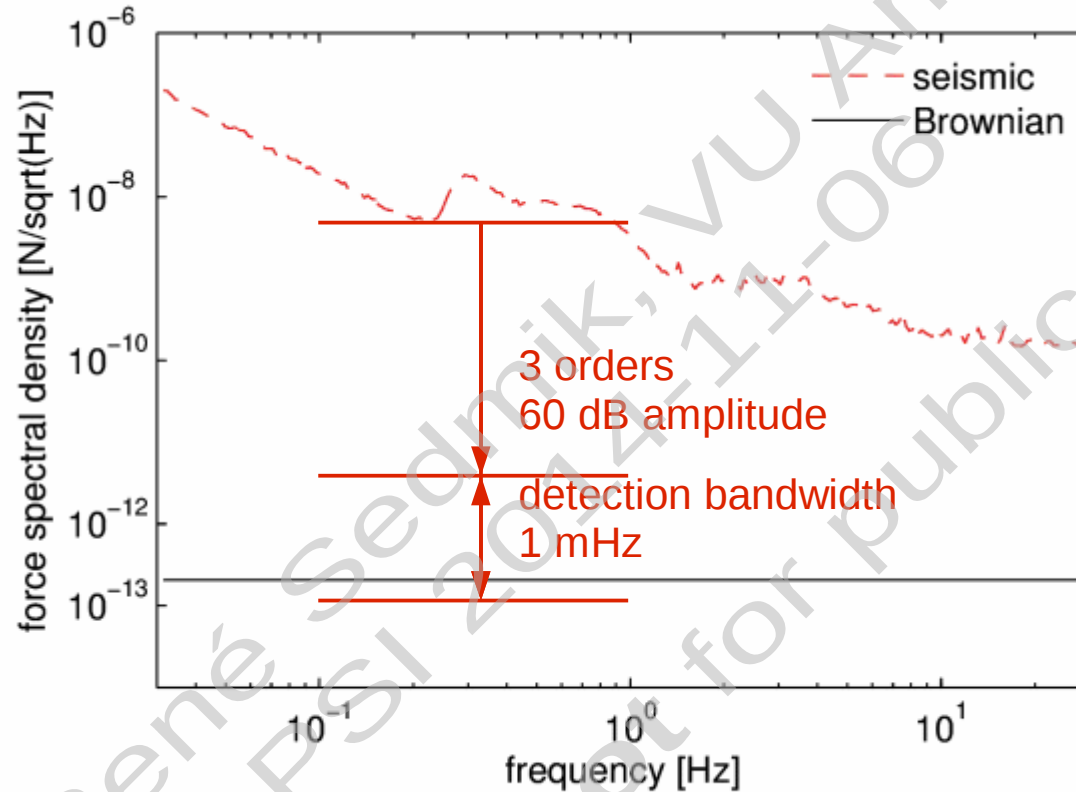
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# Vibration insulation: Required performance

**Q:** What performance do we need?

Aim: Want to measure 0.1 pN at < 1 Hz, vertical direction

Disturbances:



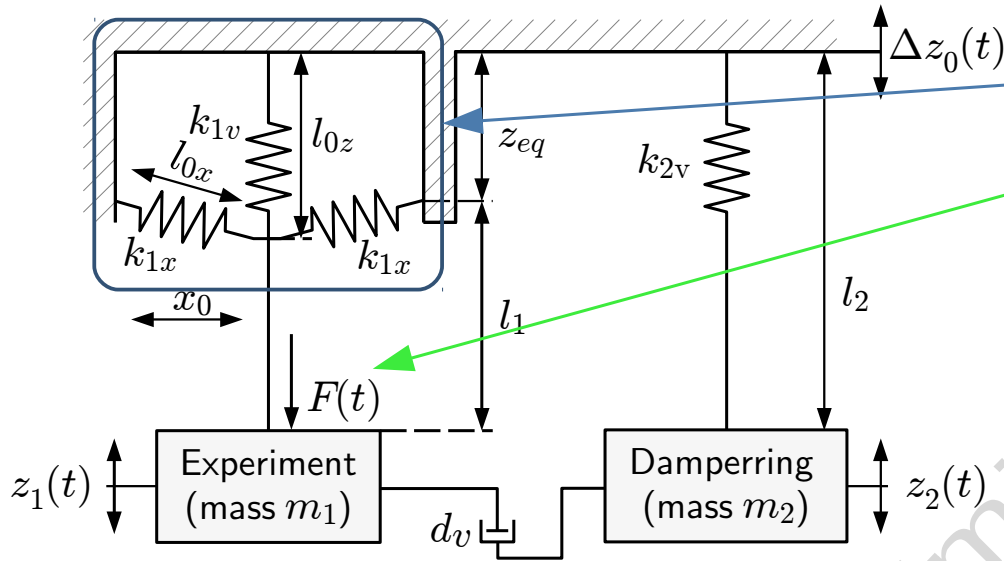
geophone  
measurement

$$F_B = \sqrt{4k_B T \frac{\gamma}{k}}$$

$\gamma$  ... structural damping  
 $k$  ... spring constant

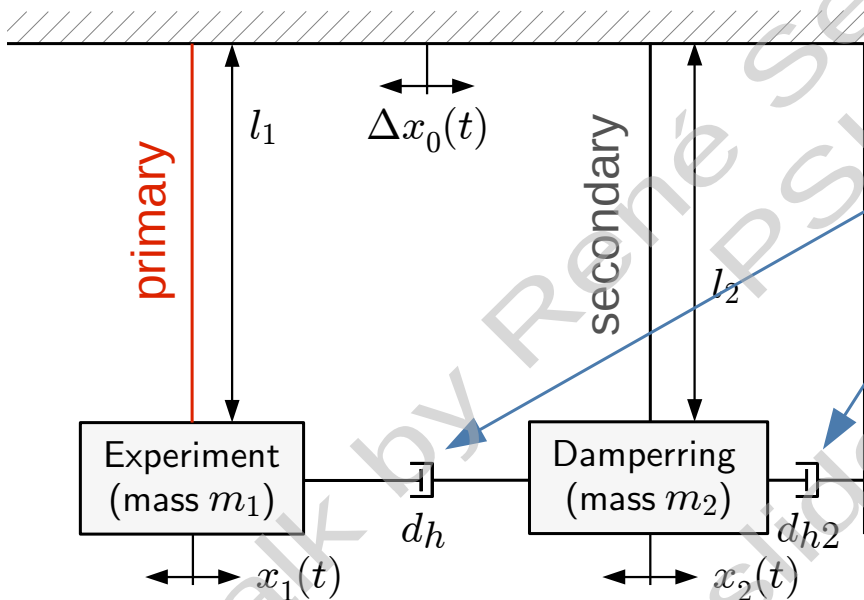
**A:** Improvement by 3 orders in vertical amplitude, 0.1 Hz – 1 Hz.

# Vibration insulation: Principle



## Vertical seismic:

- GAS\* filter
- active  $\mathcal{H}_\infty$  feedback with voice coil driver

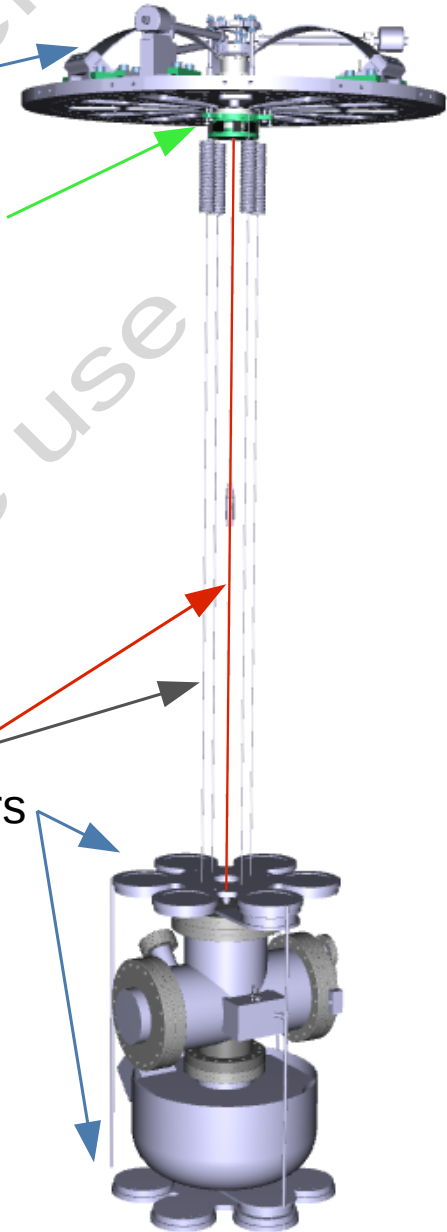


## Horizontal seismic:

- Double-pendulum
- Eddy current dampers

## Acoustic:

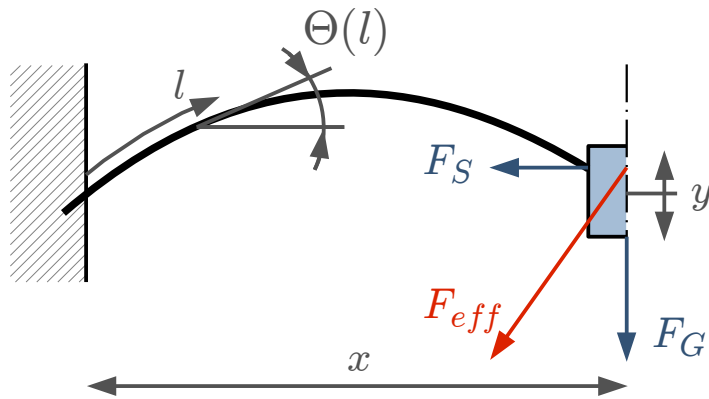
- rigid all-enclosing vacuum chamber  $10^{-5}$  mbar (not shown)



\* Gravimetric anti-spring

# Vibration insulation: Vertical: GAS filter

## Principle: Euler spring

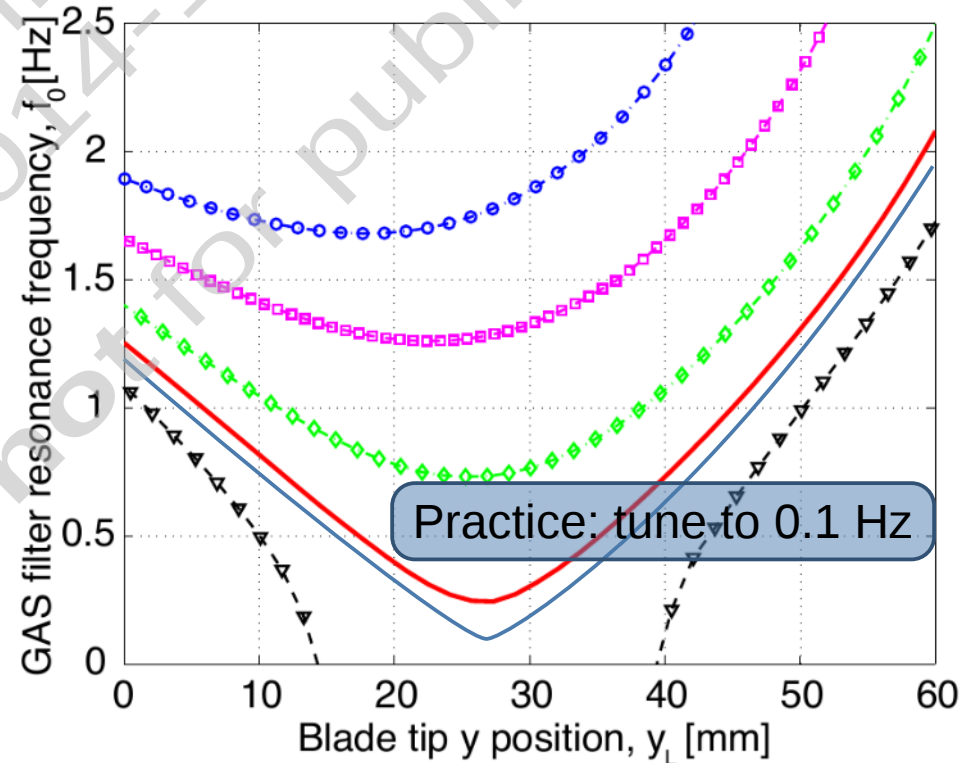
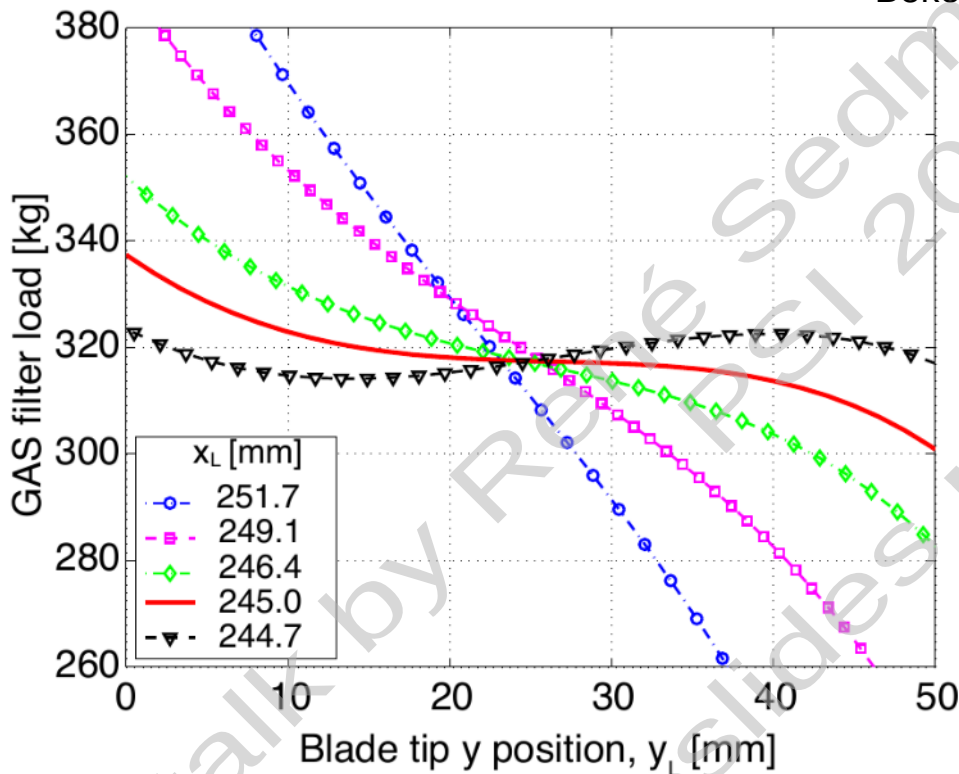


### Model:

$$U = \int_0^L dl \frac{EI(l)}{2} \left( \frac{\partial \Theta(l)}{\partial l} \right)^2 - F_S \cos \Theta(l) - F_G \sin \Theta(l)$$

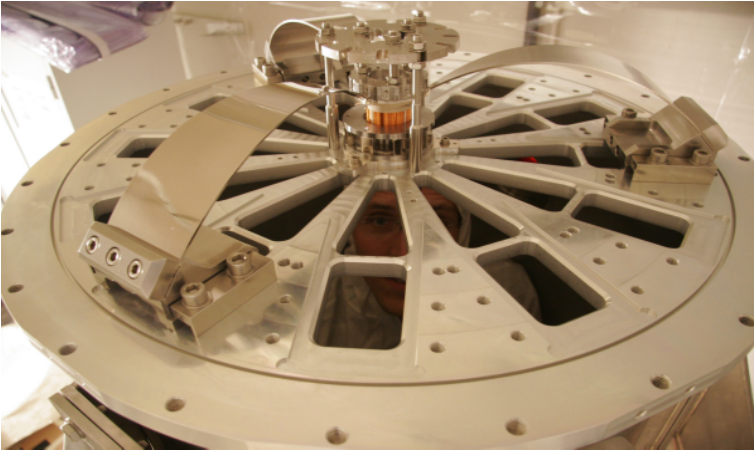
Solution for dynamics from:  $\frac{\delta U(\Theta)}{\delta \Theta} = 0$

Bertolini *et al*, *Nucl. Inst. Meth. Phys. Res A*, **435** (1999), 475  
 Beker, PhD Thesis, VU Amsterdam (2013)



# Vibration insulation: Vertical: GAS filter

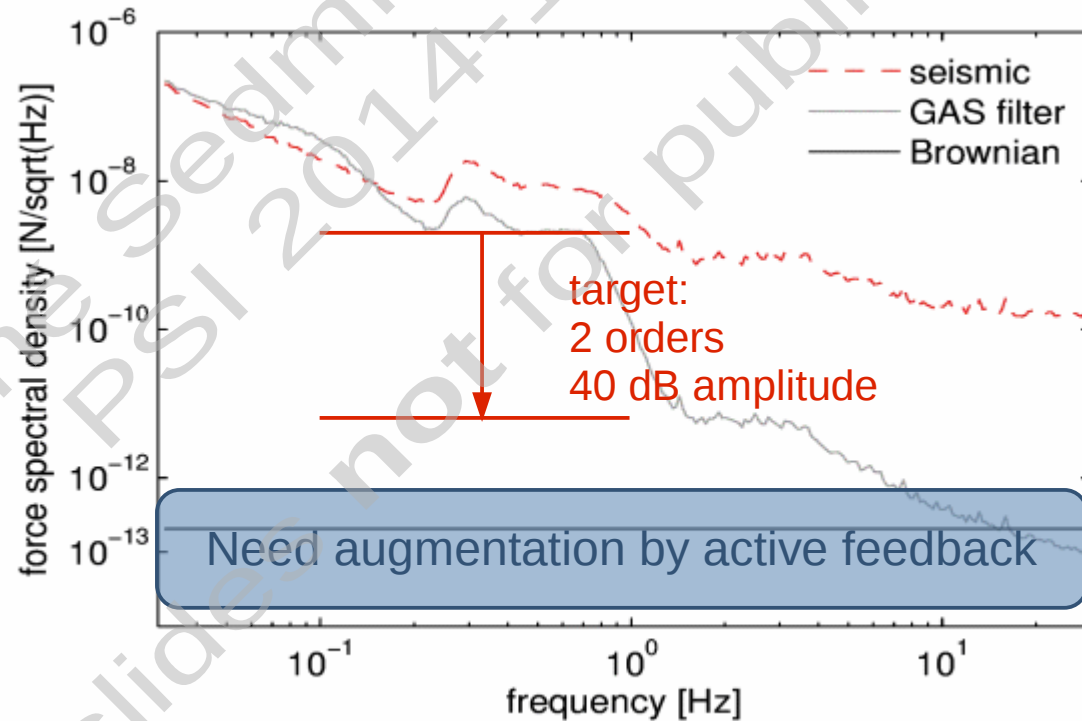
## Performance



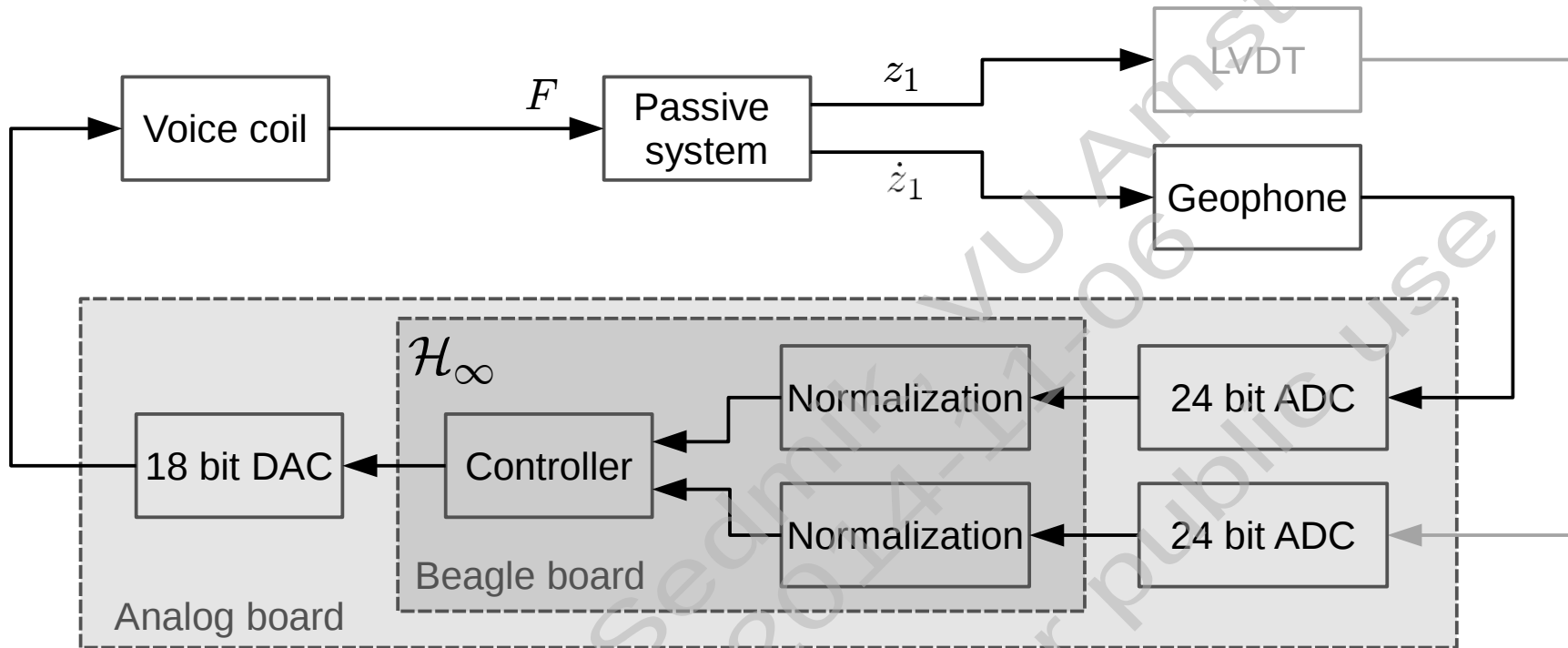
## Implementation:

- 3 springs
- suspended mass 85 kg
- tuning in air: 150 mHz

## Numerical results:



## Schematic

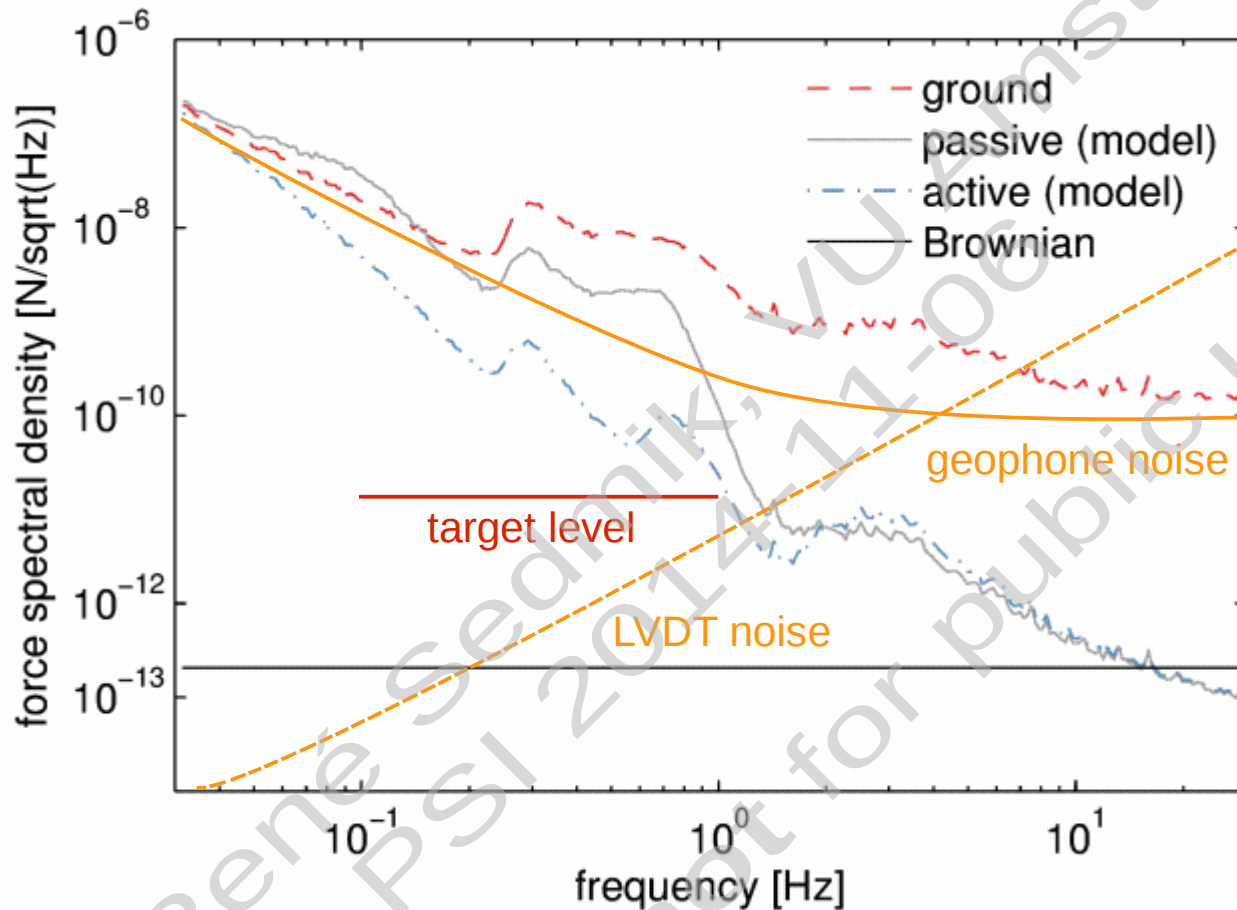


**Q: Why is a simple PID-controller not enough?**

- A:**
1. Sensors and system give too much phase delay.
  2. SNR is  $\sim 1$  at low frequencies

Need a predictive system for the passive dynamics for stability

Performance:



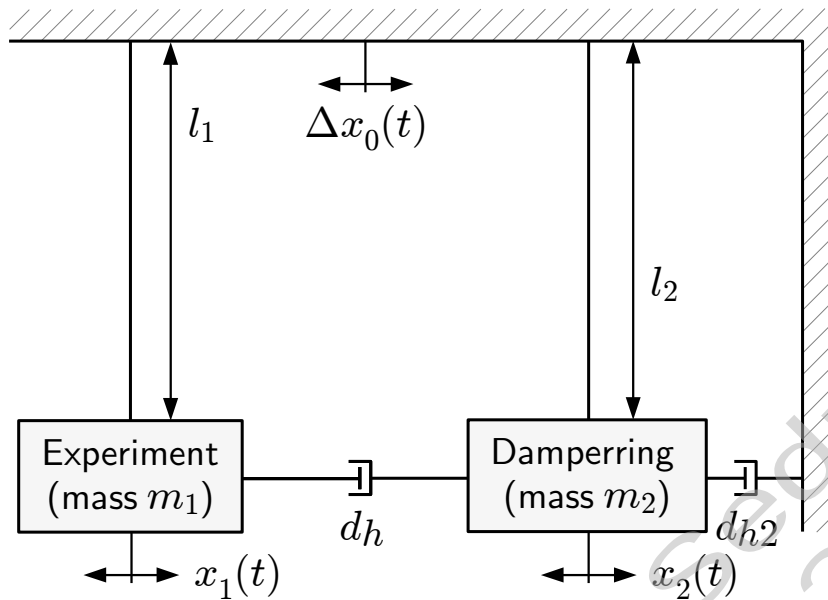
Still ~ one order of magnitude missing  
Problem: sensor noise of the geophone

Approach to improve:

- enhance preamplifiers (factor 10)
- Add LVDT for better signal at low frequencies

# Vibration insulation: Horizontal: Double-pendulum

Principle: Low-pass  
and energy dissipation

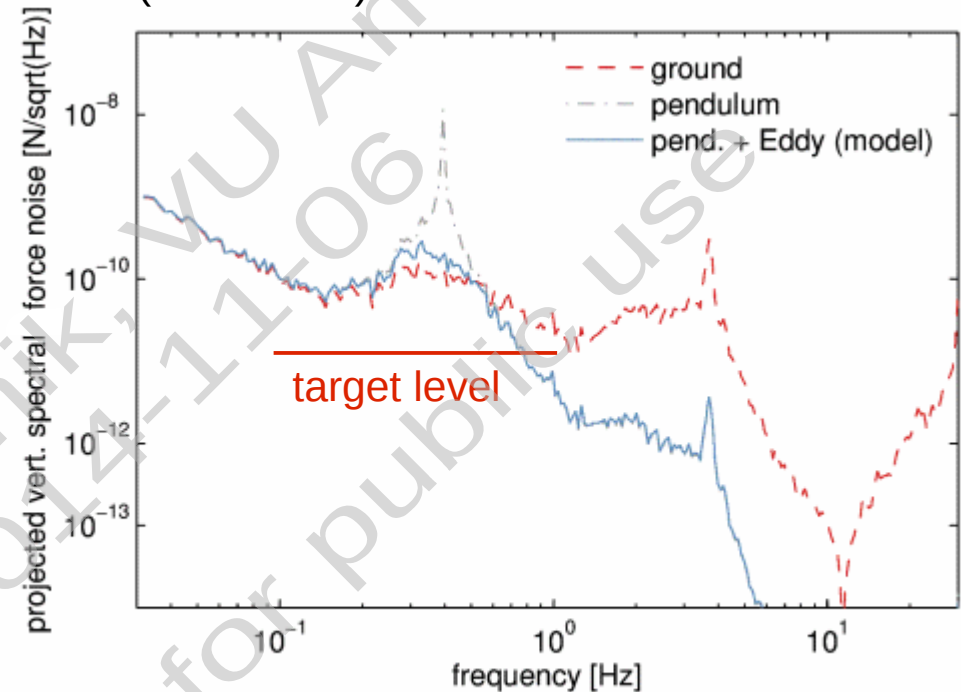


## Requirements:

$10^{-11}$  N/ $\sqrt{\text{Hz}}$  at 0.1 Hz – 1 Hz

Performance:

projected to vertical, 10% coupling  
(worst case)

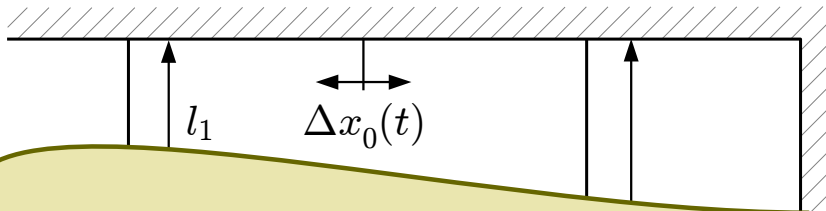


## Approach for improvement:

- increase damping
- assure that sensor has  $< 3\%$  coupling  $H \rightarrow V$

# Vibration insulation: Horizontal: Double-pendulum

Principle: Low-pass  
and energy dissipation

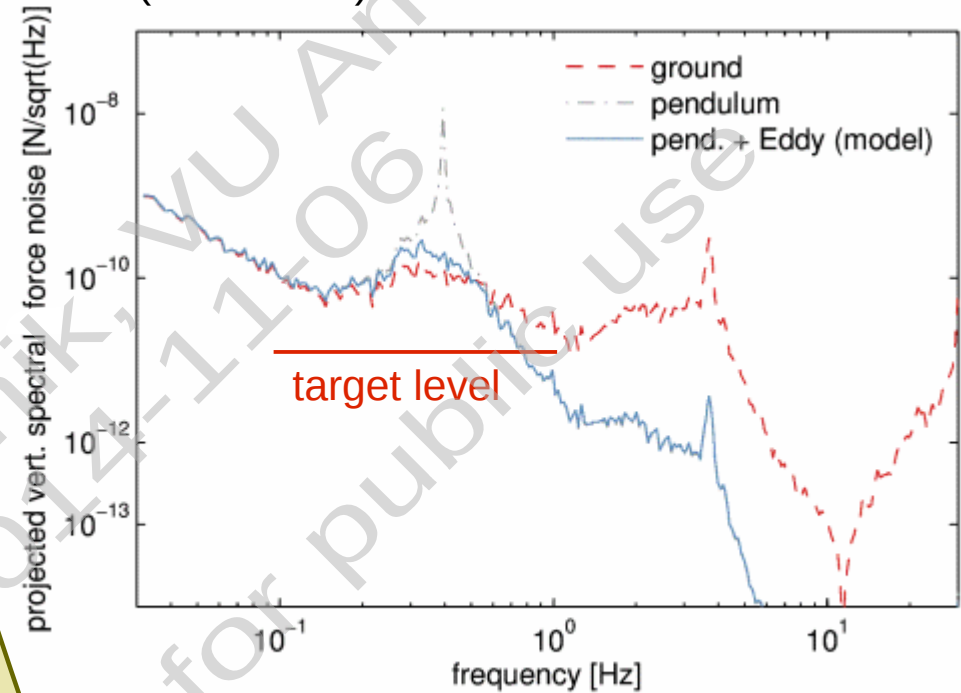


## Shopping List:

- ✓ Force transducer for  $< 1 \text{ N/m}$   
**0.25 N/m**
- Capacitive detection system:  
(C-bridge ~~1 aF~~ at nom. 100 pF)  
**currently 40 aF, need to improve**
- Parallelism control  ~~$\sim 0.1 \mu\text{rad}$~~   
**rough test:  $\sim 3 \mu\text{rad}$**
- ✓ Ultra-flat surfaces (waviness  $< 20 \text{ nm}$ )
- Vibration isolation  ~~$\sim 60 \text{ dB}$  at 1 Hz~~  
 **$\sim 40 \text{ dB}$   
need. improvement**

Performance:

projected to vertical, 10% coupling  
(worst case)

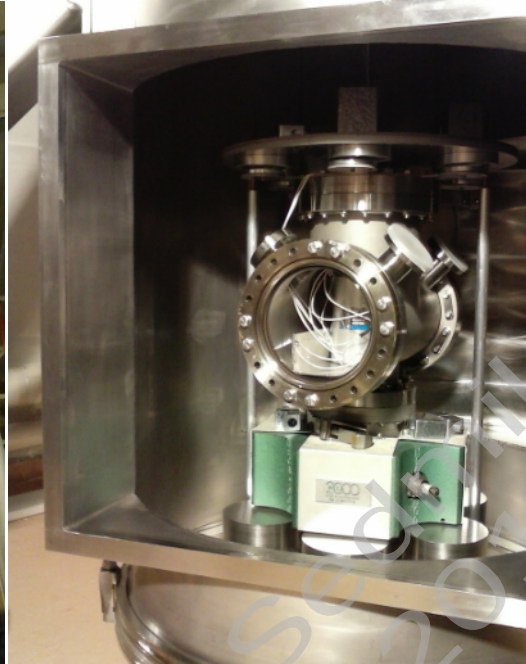


Approach for improvement:

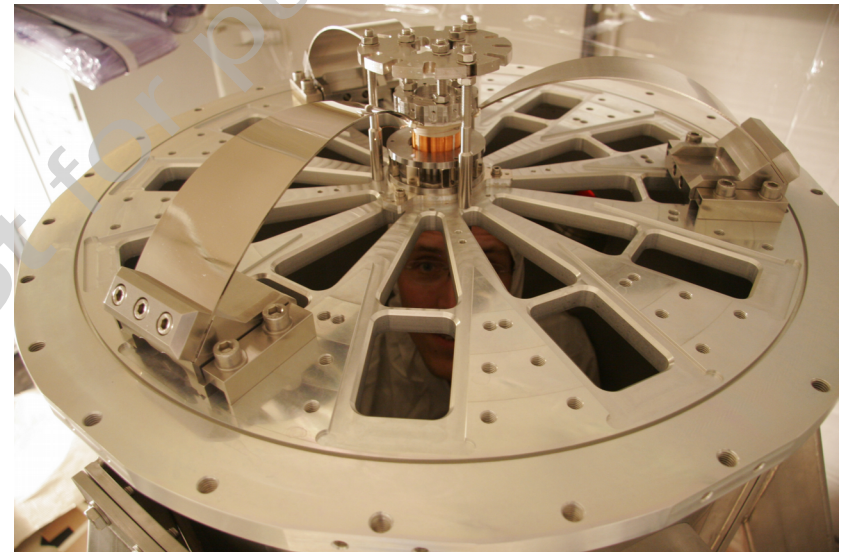
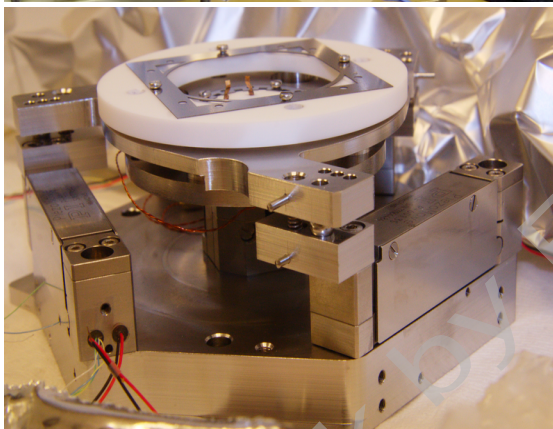
- increase damping
- assure that sensor has  $< 3\%$  coupling  $H \rightarrow V$

# Outlook & Conclusion

Status:



- The experiment is still under construction
- Subsystems are being tested
- **First force measurements planned still this yer: Proof of principle**
- For final sensitivity, several improvements required (long term)



# Outlook & Conclusion

## Conclusion:

- Casimir forces potentially play a role in the development of MEMS but are interesting in many ways.
- Chameleon forces could explain the accelerated expansion of the universe.
- Both forces could be measured in a high-accuracy parallel plate experiment
- Status: Experiment has been designed and built, close to turning the switch
- First preliminary results (hopefully) later this year.

# Acknowledgments

## Thank you for your kind attention!

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Han Voet  
Rob J. Limburg  
Jan Rector  
Davide Iannuzzi

Nikhef:

Alessandro Bertolini  
Eric Hennes  
Arnold Rietmeijer  
Martin Duets  
Johannes v. d. Brand

Stay tuned on our website  
[cannex.vu.nl](http://cannex.vu.nl)



partially crowd-funded via **Flintwave**  
Make Research & Innovation Visible

and **ASML**

# Finite Temperature

## Gibbs free energy

Remember Lifshitz at  $T=0$

$$\frac{E_0}{A} = \frac{\hbar}{2} \int_0^\infty \frac{dk_x dk_y}{(2\pi)^2} \sum_{n=0}^\infty \omega_n(a)$$

Lifshitz at  $T \neq 0$

$$\frac{\mathcal{F}(T)}{A} = \frac{\hbar}{2} \int_0^\infty \frac{dk_x dk_y}{(2\pi)^2} \sum_{n=0}^\infty \left[ \omega_n(a) + \frac{k_B T}{\hbar} \left( 1 + e^{\frac{\hbar \omega_n(a)}{k_B T}} \right) \right]$$

$$\mathcal{F}(T) = E_0 + \Delta\mathcal{F}(T)$$

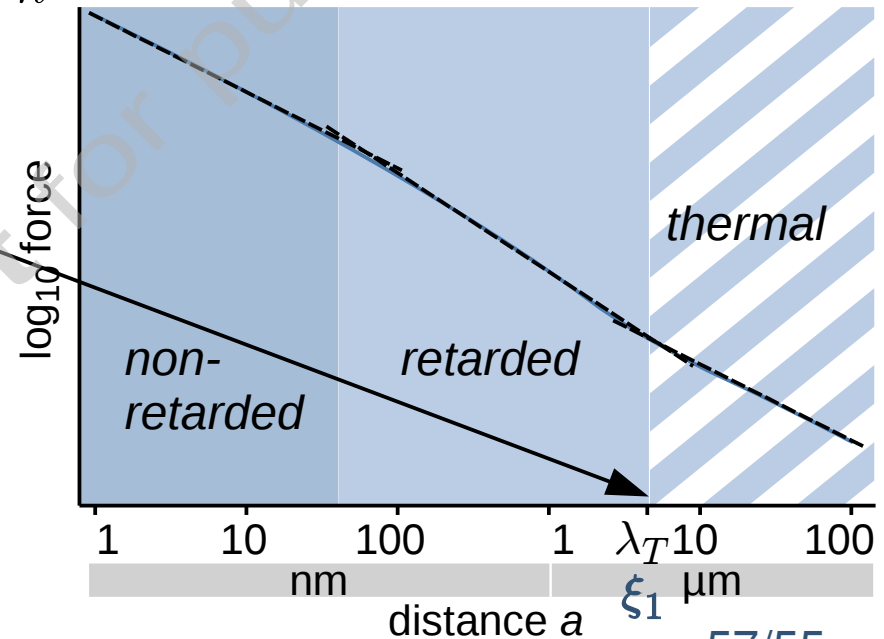
Matsubara formalism: Formal substitution

$$\frac{\hbar}{2\pi} \int_0^\infty d\xi \rightarrow k_B T \sum'_n \quad \text{and} \quad \xi \rightarrow \xi_n := 2\pi \frac{k_B T}{\hbar} n \quad \dots \text{ Matsubara frequencies}$$

term at  $n=0$  has an additional factor 1/2

$$\xi_0 = 0$$

$$\xi_1|_{T=300\text{K}} \approx 2.5 \times 10^{14} \text{ rad/s} \Rightarrow 7.6 \mu\text{m}$$



$$\mathcal{L}_{\text{SM}} \rightarrow \mathcal{L}'_{\text{SM}} + \mathcal{L}_{\text{C}}$$

$$\frac{e^{\frac{\beta_\gamma}{M_{Pl}} \phi}}{4} F^2 + \sum_i \mathcal{L}^{(i)} \left( e^{\frac{\beta_i}{M_{Pl}} \phi}, \psi^{(i)} \right) \quad \sqrt{-g} \left[ \mathcal{R} \frac{M_{Pl}^2}{2} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$

modified metric

$$\tilde{g}_{\mu\nu} = e^{\frac{2\beta^{(i)} \phi}{M_{Pl}}} g_{\mu\nu} \quad g = \det g_{\mu\nu}$$

energy momentum tensor

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}}{\delta g^{\mu\nu}}$$

Energy (natural units)

$$\tilde{\rho} \equiv \tilde{g}^{\mu\nu} T_{\mu\nu} \quad \rho_i \equiv \tilde{\rho} e^{\frac{3\beta_i}{M_{Pl}} \phi}$$

effective potential

$$V_{\text{eff}} = V(\phi) + \sum_i \rho_i e^{\frac{\beta_i}{M_{Pl}} \phi}$$