

Simulating Particle Physics with Quantum Computers

Mathieu PELLEN

University of Freiburg

Colloquium - Paul Scherrer Institut, Switzerland

23rd of April 2026



universität freiburg

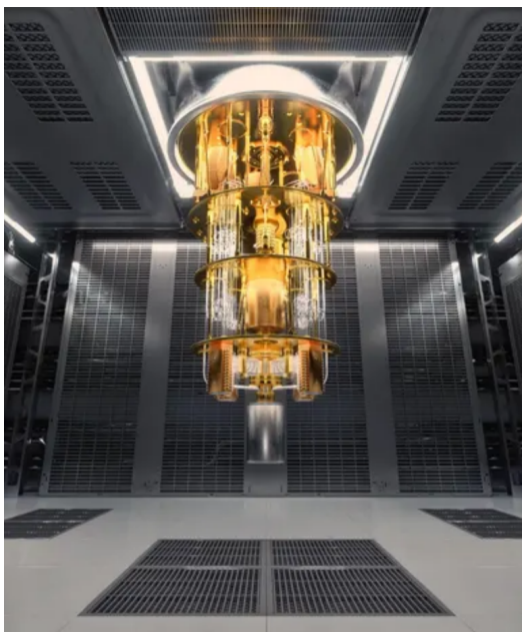
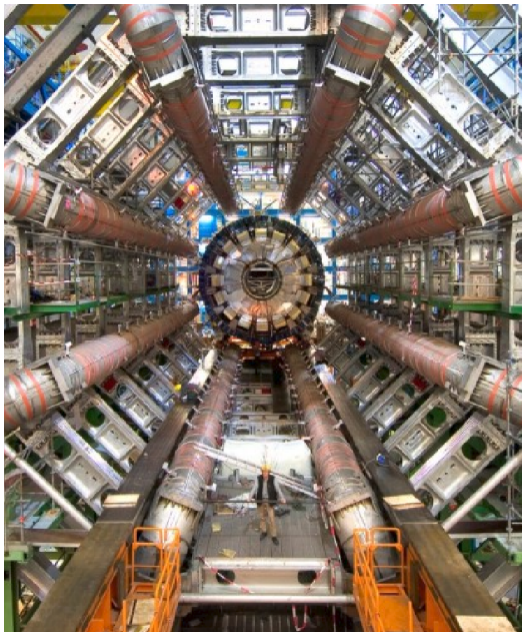
Outline:

Why quantum computing and high-energy physics?

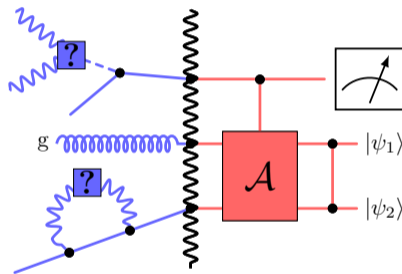
- Basics of quantum computing
 - How to build a quantum circuit
- Context and Basics of high-energy physics

How?

- Applications
 - Cross section → quantum integration
 - Amplitudes in QCD → dedicated quantum algorithms
 - QCD parton shower → dedicated quantum algorithms

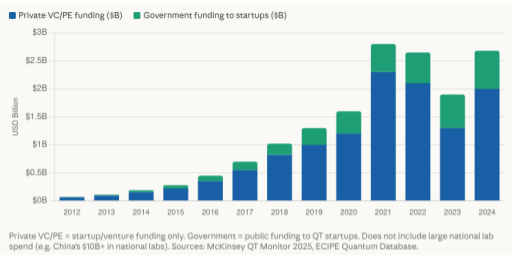


- Quantum computing and high-energy physics are fast fields of research!
 - personal view on quantum computing for high-energy physics
- Beginning of quantum computing (real-world) applications
 - do not expect yet comparisons with classical state of the art now
- Not a review of all applications
 - selected applications for illustration
 - references provided to go beyond



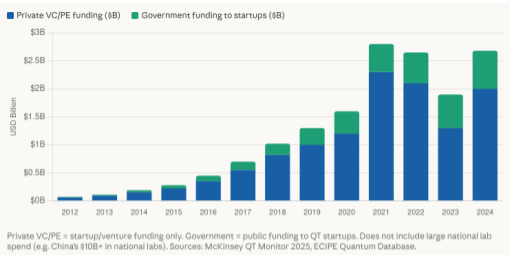
Context in quantum computing

- Great commercial interest



Context in quantum computing

- Great commercial interest
- Various physics approaches (from both big players and start-ups)

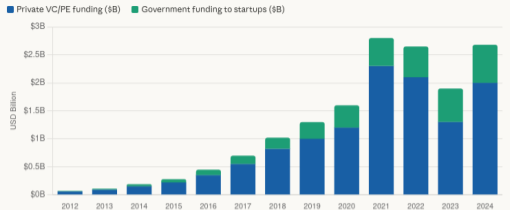


IDTechEx Research Eight Leading Approaches to Building a Commercial Quantum Computer



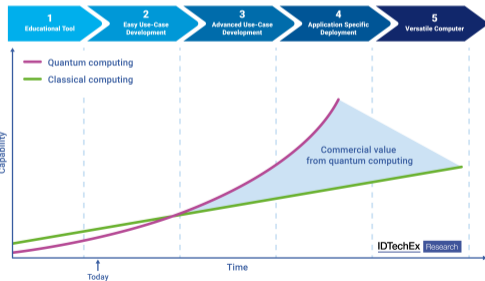
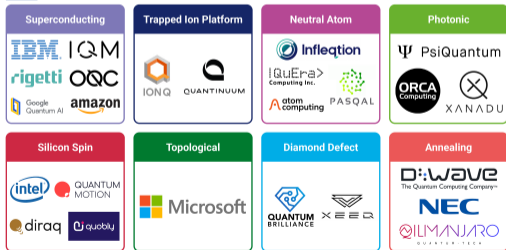
Context in quantum computing

- Great commercial interest
- Various physics approaches (from both big players and start-ups)
- Critical moment



Private VC/PE = startup/venture funding only. Government = public funding to QT startups. Does not include large national lab spend (e.g. China's \$10B+ in national labs). Sources: McKinsey QT Monitor 2025, ECIPE Quantum Database.

IDTechEx Research Eight Leading Approaches to Building a Commercial Quantum Computer



Why quantum computing and high-energy physics?

- Basics of quantum computing
 - How to build a quantum circuit

Literature:

- *Quantum Computation and Quantum Information*, Nielsen and Chuang
- *Programming Quantum Computers*, Johnston, Harrigan, and Gimeno-Segovia

FEBRUARY 17, 2014

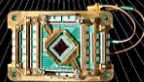
French Advances / My Doctor Fired Me / Love App-tually

TIME

IT PROMISES TO SOLVE SOME OF HUMANITY'S
MOST COMPLEX PROBLEMS. IT'S BACKED
BY JEFF BEZOS, NASA AND THE CIA.
EACH ONE COSTS \$10,000,000 AND OPERATES
AT 459° BELOW ZERO. AND NOBODY KNOWS
HOW IT ACTUALLY WORKS

THE INFINITY MACHINE

BY LEV GROSSMAN



Back to quantum mechanics...

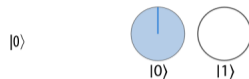
→ Any state can be written as

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix},$$

with $|\alpha|^2 + |\beta|^2 = 1$

→ Representation

Possible values of a qubit Graphical representation



Back to quantum mechanics...

→ Any state can be written as

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix},$$

$$\text{with } |\alpha|^2 + |\beta|^2 = 1$$

→ In quantum mechanics (and in quantum computing), any operation A is unitary

$$\psi \xrightarrow{A} \psi'$$

$$A |\psi\rangle = |\psi'\rangle \quad \text{with} \quad A^\dagger A = AA^\dagger = \text{Id}$$



Example of gates (I)

Pauli-X (X)

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\psi = \alpha |0\rangle + \beta |1\rangle \rightarrow \psi' = \alpha |1\rangle + \beta |0\rangle$$



Example of gates (I)

Pauli-X (X)

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\psi = \alpha |0\rangle + \beta |1\rangle \rightarrow \psi' = \alpha |1\rangle + \beta |0\rangle$$



Hadamard (H)

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$\psi = \alpha |0\rangle + \beta |1\rangle \rightarrow \psi' = \frac{\alpha+\beta}{\sqrt{2}} |0\rangle + \frac{\alpha-\beta}{\sqrt{2}} |1\rangle$$



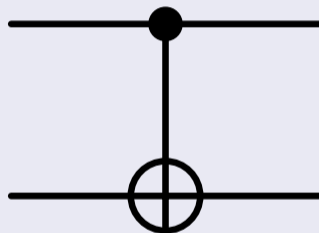
Example of gates (II)

Controlled not (CNOT, CX)

$$CX = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

$|00\rangle \rightarrow |00\rangle$; $|01\rangle \rightarrow |01\rangle$;
 $|10\rangle \rightarrow |11\rangle$; $|11\rangle \rightarrow |10\rangle$.

- If 0 \rightarrow nothing happens; If 1 \rightarrow Pauli-X (X)!
- *Control* qubit (top) and *target* qubit (bottom)



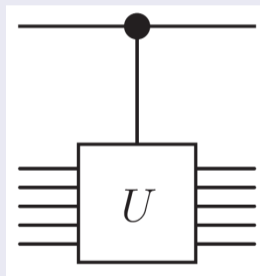
Example of gates (III)

Generalised controlled gate (CU)

$$CU = \begin{bmatrix} \text{Id}_2 & 0 \\ 0 & U \end{bmatrix}$$

→ One *control* qubit and ...
... many *target* qubits

NB Also possible:
Many *control* qubits and ...
... one/*many target* qubit



Measurement process

→ State: $|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$

→ Probabilities: $\begin{cases} \text{for } |0\rangle : \alpha^2 \\ \text{for } |1\rangle : \beta^2 \end{cases}$ with $|\alpha|^2 + |\beta|^2 = 1!$

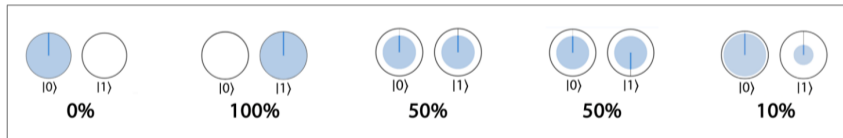
→ Measurement = squaring/collapse the amplitude!

Measurement process

→ State: $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$

→ Probabilities: $\begin{cases} \text{for } |0\rangle : \alpha^2 \\ \text{for } |1\rangle : \beta^2 \end{cases}$ with $|\alpha|^2 + |\beta|^2 = 1!$

→ Measurement = squaring/collapse the amplitude!



NB: Probability obtained after many measurements/shots! $\langle\psi|\psi\rangle!$

→ α and β encode information that can be measured/computed!

- Live example with QISKIT [open-source software development kit from IBM]

- Live example with QISKIT [open-source software development kit from IBM]

Quantum-computing applications

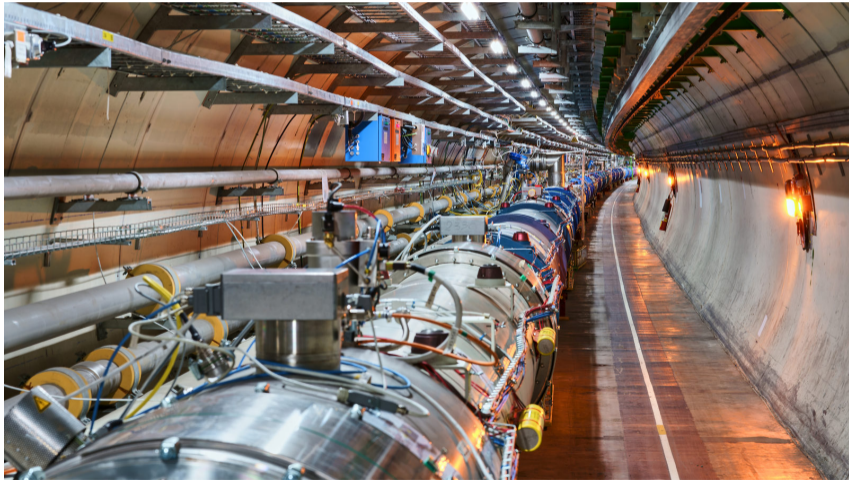
- Key parameters of quantum computers:
 - Number of qubits
 - amount of operations in parallel
 - Depth/number of gate operations
 - amount of operations in sequence (before decoherence/noise)
 - Error rates
 - Fault-tolerant era \sim 2030 (depending on the technology)
- Expressibility is a combination of number of qubits and depth
- Implementation of a problem in quantum computing
 - = implementation with low-level language with restrictions (unitary operations)

Why quantum computing and high-energy physics?

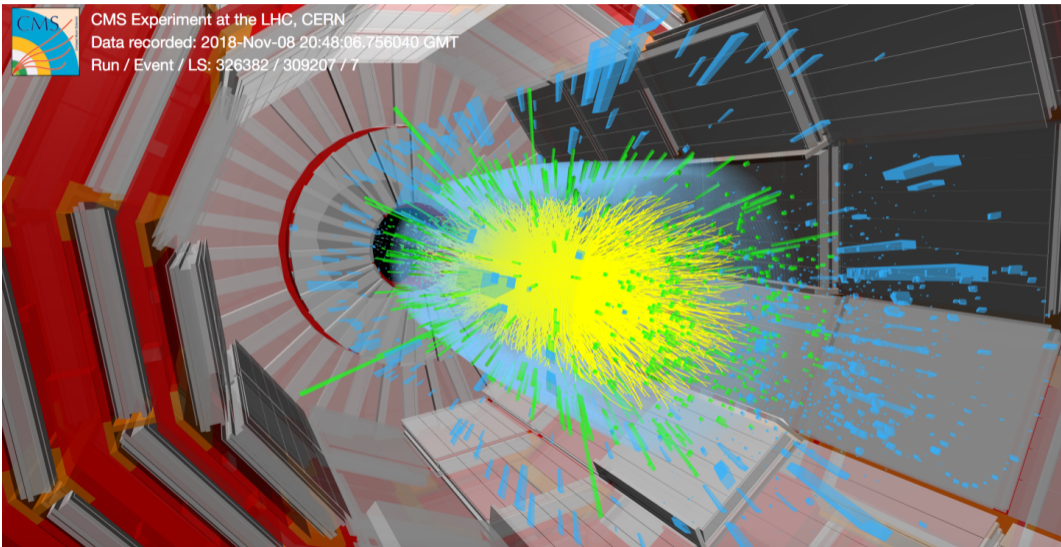
- Basics of high-energy physics and context

High-energy physics = LHC @ CERN, Geneva (Switzerland)

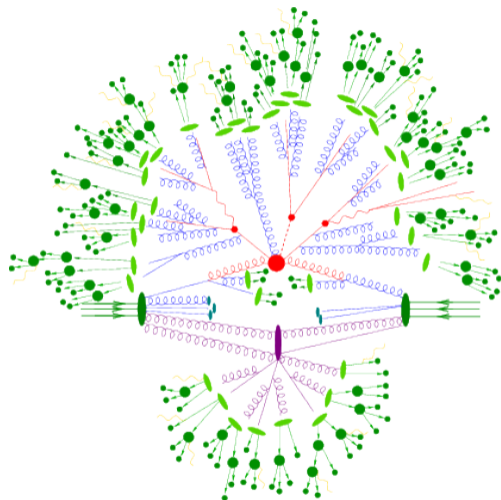
- 27km-long tunnel where protons are collided at high-energy (13.6 TeV currently)
- Allow to probe fundamental interactions



Life at the LHC (in reality)



- Factorisation principle:
 - **high-energy** (perturbative)
 - Hard-scattering, parton-shower
 - **low-energy** (non-perturbative)
 - parton distribution function, hadronisation, underlying events, ...
- **Focus of the presentation:**
high energy!
 - Calculated from first principles



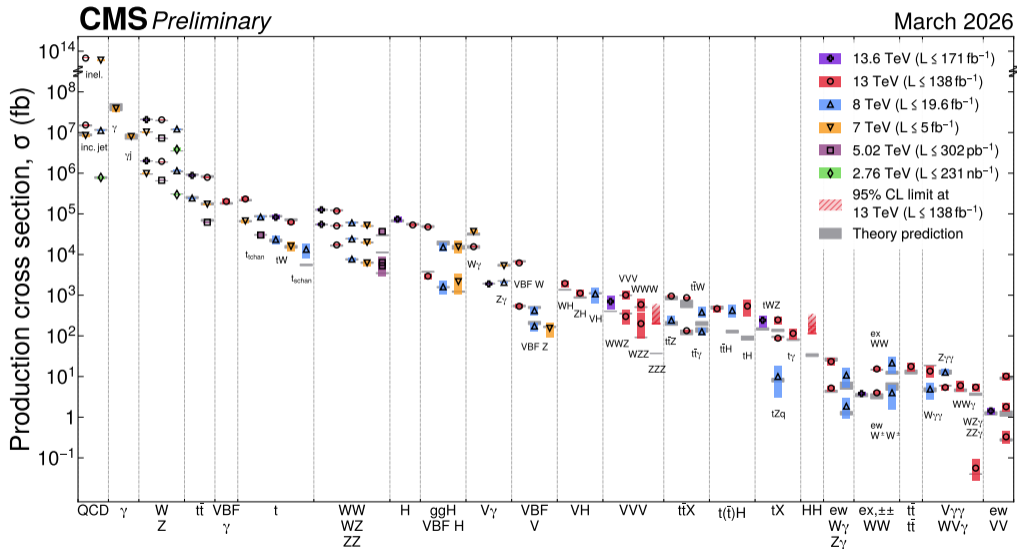
[source: Sherpa]

- Comparison between theory and experiment to learn about particle physics!



[source: bing image creator, Engel, Pellen]

- Comparison between theory and experiment to learn about particle physics!

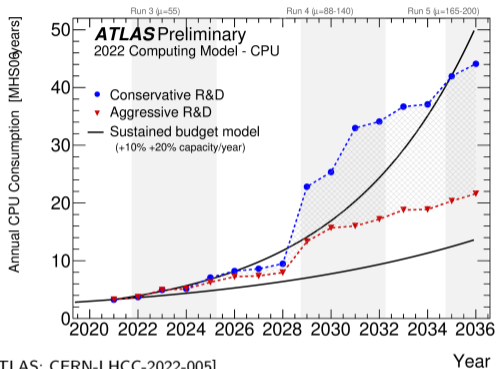


→ Event generation:

~ 15% of ~ 3 billion cpu.h.y^{-1} for ATLAS

[Buckley; 1908.00167], [Valassi et al.; 2004.13687]

Computing problem in high-energy physics



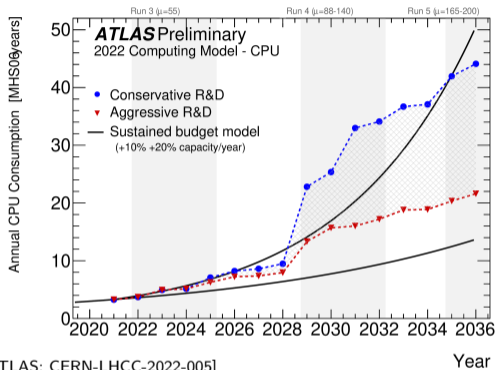
[ATLAS; CERN-LHCC-2022-005]

→ Event generation:

~ 15% of ~ 3 billion cpu.h.y^{-1} for ATLAS

[Buckley; 1908.00167], [Valassi et al.; 2004.13687]

Computing problem in high-energy physics



[ATLAS; CERN-LHCC-2022-005]

→ Event generation:

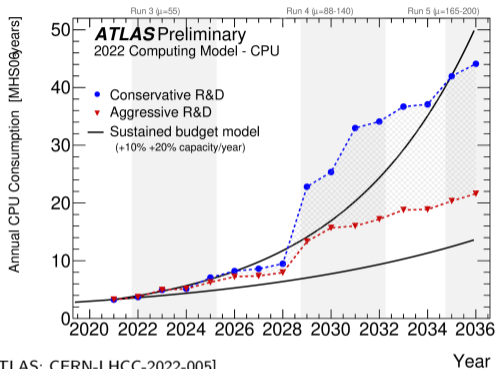
~ 15% of ~ 3 billion cpu.h.y^{-1} for ATLAS

[Buckley; 1908.00167], [Valassi et al.; 2004.13687]

• Solutions: GPU & AI

[Borowka et al.; 1811.11720], [Carrazza et al.; 2002.12921, 2009.06635, 2106.10279], [Bothmann et al.; 2106.06507], [Cruz-Martinez, De Laurentis, MP; 2502.07060], [Seymour, Sule; 2403.08692, 2511.19633], [Valassi; 2510.05392], [Bothmann et al.; 2604.03511]

Computing problem in high-energy physics



[ATLAS; CERN-LHCC-2022-005]

→ Event generation:

~ 15% of ~ 3 billion cpu.h.y^{-1} for ATLAS

[Buckley; 1908.00167], [Valassi et al.; 2004.13687]

• Solutions: GPU & AI

[Borowka et al.; 1811.11720], [Carrazza et al.; 2002.12921, 2009.06635, 2106.10279], [Bothmann et al.; 2106.06507], [Cruz-Martinez, De Laurentis, MP; 2502.07060], [Seymour, Sule; 2403.08692, 2511.19633], [Valassi; 2510.05392], [Bothmann et al.; 2604.03511]

- Can quantum computing be of any use in HEP?
 - to compute things faster/more efficiently?
 - to compute new things?
- Quantum computing provides a natural framework for QFT calculations?

- [Gray, Terashi; Gray:2022fou] (selected topics), [Delgado et al.; 2203.08805] (Snowmass), [Klco et al.; 2107.04769] (lattice), [Di Meglio et al.; 2307.03236] (CERN Quantum Initiative)

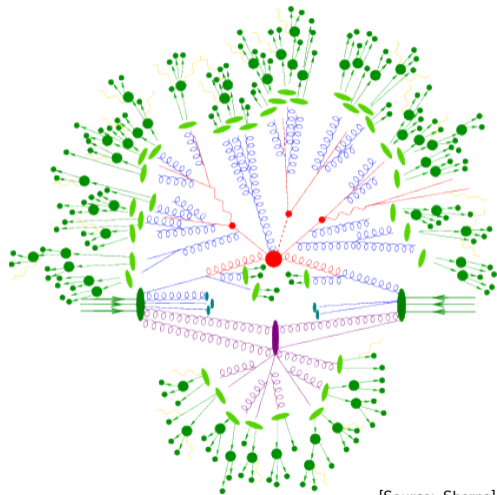
Selected references

- Amplitude/loop integrals: [Ramirez-Uribe et al.; 2105.08703], [Bepari, Malik, Spannowsky, Williams; 2010.00046], [Chawdhry, MP; 2303.04818], [Chawdhry, MP, Williams; 2507.07194], [Bashore, Moretti, Vitos; 2507.14252], [Ochoa-Oregon et al.; 2508.04019]
- Monte Carlo (integration): [Bravo-Prieto et al; 2110.06933], [Agliardi, Grossi, MP, Prati; 2201.01547], [Martínez de Lejarza et al.; 2401.03023], [Martínez de Lejarza, Cieri, Rodrigo; 2204.06496], [Cruz-Martinez, Robbiati, Carrazza; 2308.05657], [Williams, MP; 2502.14647], [Martínez de Lejarza et al.; 2409.12236], [Pyretzidis, Martínez de Lejarza, Rodrigo; 2506.19965]
- Parton shower: [Bauer, de Jong, Nachman, Provasoli; 1904.03196, 1901.08148], [Bepari, Malik, Spannowsky, Williams; 2010.00046], [Williams, Malik, Spannowsky, Bepari; 2109.13975], [Deliyannis et al.; 2203.10018], [Chigusa, Yamazaki; 2204.12500], [Gustafson, Prestel, Spannowsky, Williams; 2207.10694], [Bauer, Chigusa, Yamazaki; 2310.19881]
- Others: [Perez-Salinas, Cruz-Martinez, Alhajri, Carrazza; 2011.13934], [Bauer, Freytsis, Nachman; 2102.05044], [Martínez de Lejarza et al.; 2503.16073]

Quantum applications in high-energy physics

→ Quantum applications still in their infancy!

- Is it possible?

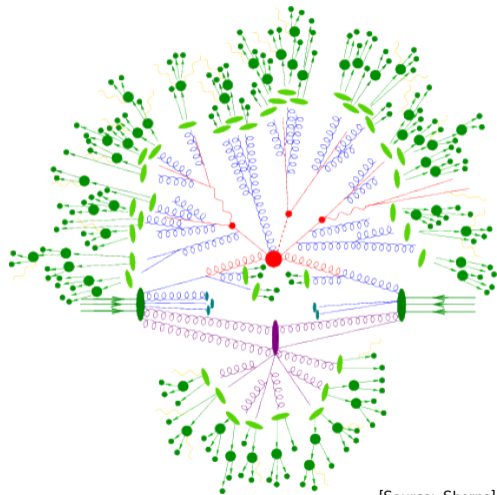


[Source: Sherpa]

Quantum applications in high-energy physics

→ Quantum applications still in their infancy!

- Is it possible?
- Is there a (theoretical) quantum advantage?

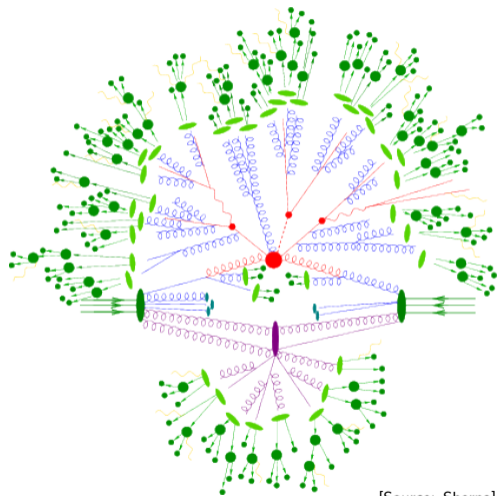


[Source: Sherpa]

Quantum applications in high-energy physics

→ Quantum applications still in their infancy!

- Is it possible?
- Is there a (theoretical) quantum advantage?
- Is it more resource efficient/faster than CPU/GPU/ML approaches?



[Source: Sherpa]

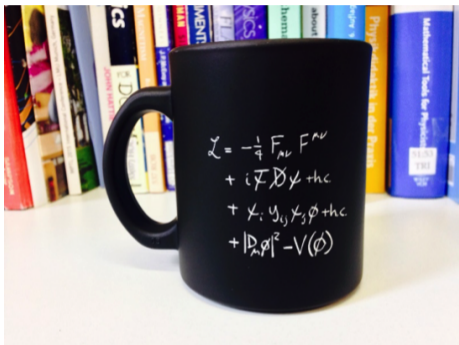
How?

- Applications

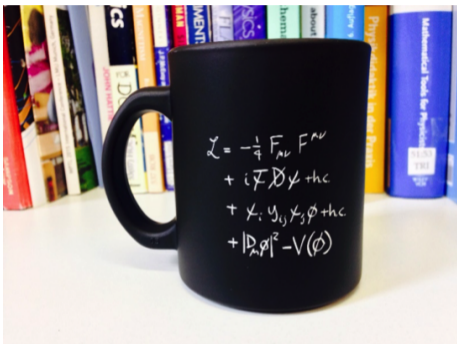
- Cross section → quantum integration
- Amplitudes in QCD → dedicated quantum algorithms
- QCD parton shower → dedicated quantum algorithms

The Standard Model of particle physics

- Particles are represented by fields
- Lagrangian displays interactions of fields
- Based on gauge symmetry
 - $SU_{\text{QCD}}(3) \times SU_{\text{L}}(2) \times U_{\text{Y}}(1)$



[source: CERN]



[source: CERN]

The Standard Model of particle physics

- Particles are represented by fields
- Lagrangian displays interactions of fields
- Based on gauge symmetry
 - $SU_{\text{QCD}}(3) \times SU_{\text{L}}(2) \times U_{\text{Y}}(1)$

- Strong force
 - Quantum chromodynamics (QCD) mediated by gluon
- Electromagnetic force
 - Quantum electrodynamics (QED) mediated by photon
- Weak force
 - Weak interaction mediated by W and Z

→ W and Z get their mass from Higgs mechanism

$$\begin{aligned}
& -\frac{1}{2}g_s g_u^a g_u^a g_u^a - g_s f^{abc} \partial_\mu g_u^b g_u^c - \frac{1}{2}g_s^2 f^{abc} f^{abc} g_u^a g_u^b g_u^c + \\
& \frac{1}{2}ig_s^2 (\bar{q}^i \gamma^\mu q^j) g_\mu^a + G^a \partial^2 G^a + g_s f^{abc} \partial_\mu G^a G^b G^c - \partial_\mu W_\nu^+ \partial_\mu W_\nu^- - \\
2. & M^2 W_\mu^+ W_\mu^- - \frac{1}{2} \partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2} \partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2} \partial_\mu H \partial_\mu H - \\
& \frac{1}{2} m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2} \partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w} M \phi^0 \phi^0 - \beta_h [\frac{2M^2}{g^2} + \\
& \frac{2M}{g} H + \frac{1}{2} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-)] + \frac{2M^4}{g^2} \alpha_h - ig_{cw} [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\mu^- W_\nu^+) - Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\nu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\nu W_\mu^+) - ig_{sw} [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\nu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\nu^+) - \frac{1}{2} g^2 W_\nu^+ W_\nu^- W_\nu^+ W_\nu^- + \\
& \frac{1}{2} g^2 W_\nu^+ W_\nu^- W_\nu^+ W_\nu^- + g^2 s_w^2 (Z_\mu^0 W_\nu^+ Z_\mu^0 W_\nu^- - Z_\mu^0 Z_\mu^0 W_\nu^+ W_\nu^-) + \\
& g^2 s_w^2 (A_\mu W_\nu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\nu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H \phi^0 \phi^0 + 2H \phi^+ \phi^-] - \\
& \frac{1}{8} g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
& g M W_\nu^+ W_\nu^- H - \frac{1}{2} g \frac{M}{c_w} Z_\mu^0 Z_\mu^0 H - \frac{1}{2} ig [W_\nu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - \\
& W_\nu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2} g [W_\nu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\nu^- (H \partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu H)] + \frac{1}{2} g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) - ig \frac{2M}{c_w} M Z_\mu^0 (W_\nu^+ \phi^- - W_\nu^- \phi^+) + \\
& ig_{sw} M A_\mu (W_\nu^+ \phi^- - W_\nu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + \\
& ig_{sw} A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4} g^2 W_\nu^+ W_\nu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
& \frac{1}{4} g^2 \frac{1}{c_w^2} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)\phi^+ \phi^-] - \frac{1}{2} g^2 \frac{2s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\nu^+ \phi^- + \\
& W_\nu^- \phi^+) - \frac{1}{2} ig^2 \frac{2s_w^2}{c_w} Z_\mu^0 H (W_\nu^+ \phi^- - W_\nu^- \phi^+) + \frac{1}{2} g^2 s_w A_\mu \phi^0 (W_\nu^+ \phi^- + \\
& W_\nu^- \phi^+) + \frac{1}{2} ig^2 s_w A_\mu H (W_\nu^+ \phi^- - W_\nu^- \phi^+) - g^2 \frac{2s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\
& g^4 s_w^2 A_\mu A_\mu \phi^+ \phi^- - e^{\lambda} (\gamma \partial + m_e^{\lambda}) e^{\lambda} - \bar{\nu}^{\lambda} \gamma \partial \nu^{\lambda} - \bar{u}_e^{\lambda} (\gamma \partial + m_u^{\lambda}) u_e^{\lambda} - \\
3. & d_1^{\lambda} (\gamma \partial + m_d^{\lambda}) d_1^{\lambda} + ig_{sw} A_\mu [-(e^{\lambda} \nu^{\lambda} e^{\lambda}) + \frac{2}{3} (\bar{u}_e^{\lambda} \gamma^{\mu} u_e^{\lambda}) - \frac{1}{3} (\bar{d}_1^{\lambda} \gamma^{\mu} d_1^{\lambda})] + \\
& \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^{\lambda} \gamma^{\mu} \nu^{\lambda}) + (e^{\lambda} \gamma^{\mu} e^{\lambda}) + (4s_w^2 - 1 - \gamma^5) e^{\lambda}] + (\bar{u}_e^{\lambda} \gamma^{\mu} (\frac{1}{3} s_w^2 - \\
& 1 - \gamma^5) u_e^{\lambda}) + (\bar{d}_1^{\lambda} \gamma^{\mu} (1 - \frac{2}{3} s_w^2 - \gamma^5) d_1^{\lambda})] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^{\lambda} \gamma^{\mu} (1 + \gamma^5) e^{\lambda}) + \\
& (\bar{u}_e^{\lambda} \gamma^{\mu} (1 + \gamma^5) C_{\lambda e} d_1^{\lambda})] + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^{\lambda} \gamma^{\mu} (1 + \gamma^5) \nu^{\lambda}) + (\bar{d}_1^{\lambda} C_{\lambda e}^{\dagger} \gamma^{\mu} (1 + \\
& \gamma^5) u_e^{\lambda})] + \frac{ig}{2\sqrt{2}} \frac{m_h^2}{M^2} [-\phi^+ (\bar{\nu}^{\lambda} (1 - \gamma^5) e^{\lambda}) + \phi^- (\bar{e}^{\lambda} (1 + \gamma^5) \nu^{\lambda})] - \\
4. & \frac{g}{2} \frac{m_h^2}{M^2} [H (\bar{e}^{\lambda} e^{\lambda}) + i\phi^0 (\bar{e}^{\lambda} \gamma^5 e^{\lambda})] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_h^2 (\bar{u}_e^{\lambda} C_{\lambda e} (1 - \gamma^5) d_1^{\lambda}) + \\
& m_h^2 (\bar{u}_e^{\lambda} C_{\lambda e} (1 + \gamma^5) d_1^{\lambda})] + \frac{ig}{2M\sqrt{2}} \phi^- [m_h^2 (\bar{d}_1^{\lambda} C_{\lambda e}^{\dagger} (1 + \gamma^5) u_e^{\lambda}) - m_h^2 (\bar{d}_1^{\lambda} C_{\lambda e}^{\dagger} (1 - \\
& \gamma^5) u_e^{\lambda})] - \frac{g}{2} \frac{m_h^2}{M^2} H (\bar{u}_e^{\lambda} u_e^{\lambda}) - \frac{g}{2} \frac{m_h^2}{M^2} H (\bar{d}_1^{\lambda} d_1^{\lambda}) + \frac{ig}{2M} \phi^0 (\bar{u}_e^{\lambda} \gamma^5 u_e^{\lambda}) - \\
& \frac{ig}{2M} \phi^0 (\bar{d}_1^{\lambda} \gamma^5 d_1^{\lambda}) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \\
5. & \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig_{cw} W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + ig_{sw} W_\mu^+ (\partial_\mu \bar{Y} X^- - \\
& \partial_\mu \bar{X}^+ Y) + ig_{cw} W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + ig_{sw} W_\mu^- (\partial_\mu \bar{X}^- Y - \\
& \partial_\mu \bar{Y} X^+) + ig_{cw} Z_\mu^0 (\partial_\mu \bar{X}^+ X^- - \partial_\mu \bar{X}^- X^+) + ig_{sw} A_\mu (\partial_\mu \bar{X}^+ X^- + \\
& \partial_\mu \bar{X}^- X^+) - \frac{1}{2} g M [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\
& \frac{1-2c_w^2}{2c_w} ig M [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} ig M [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
& ig M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2} ig M [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
\end{aligned}$$

[source: symmetrymagazine, Thomas Gutierrez]

The Standard Model of particle physics

- Particles are represented by fields
- Lagrangian displays interactions of fields
- Based on gauge symmetry
 $\rightarrow SU_{\text{QCD}}(3) \times SU_{\text{L}}(2) \times U_{\text{Y}}(1)$

- Strong force
 \rightarrow Quantum chromodynamics (QCD) mediated by gluon
- Electromagnetic force
 \rightarrow Quantum electrodynamics (QED) mediated by photon
- Weak force
 \rightarrow Weak interaction mediated by W and Z

\rightarrow W and Z get their mass from Higgs mechanism

Quark and gluon interaction

$$\mathcal{L}_\psi = \sum_q \bar{\psi}_q \left(i\not{\partial} - g_s \frac{\lambda_a}{2} \not{G}^a - m_q \right) \psi_q$$

After quantisation:

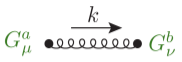
(e.g. with path integral formalism)

→ *Feynman rules* in momentum space

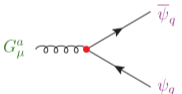
Perturbation theory and Feynman diagrams

Quark and gluon interaction

$$\mathcal{L}_\psi = \sum_q \bar{\psi}_q \left(i\not{\partial} - g_s \frac{\lambda_a}{2} \not{G}^a - m_q \right) \psi_q$$


$$G_\mu^a \xrightarrow{k} G_\nu^b \quad \frac{-i\delta^{ab}g_{\mu\nu}}{k^2 + i\epsilon}$$


$$\bar{\psi}^a \xleftarrow{k} \psi^b \quad \frac{i(\not{k} + m_q)}{k^2 - m_q^2 + i\epsilon}$$


$$G_\mu^a \xrightarrow{k} \begin{cases} \bar{\psi}_q \\ \psi_q \end{cases} \quad -i g_s \frac{\lambda^a}{2} \gamma_\mu$$

After quantisation:

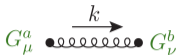
(e.g. with path integral formalism)

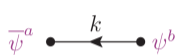
→ *Feynman rules* in momentum space

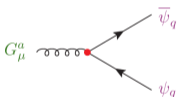
Perturbation theory and Feynman diagrams

Quark and gluon interaction

$$\mathcal{L}_\psi = \sum_q \bar{\psi}_q \left(i\not{\partial} - g_s \frac{\lambda_a}{2} \not{G}^a - m_q \right) \psi_q$$


$$\frac{-i\delta^{ab}g_{\mu\nu}}{k^2 + i\epsilon}$$


$$\frac{i(\not{k} + m_q)}{k^2 - m_q^2 + i\epsilon}$$


$$-ig_s \frac{\lambda^a}{2} \gamma_\mu$$

After quantisation:

(e.g. with path integral formalism)

→ *Feynman rules* in momentum space

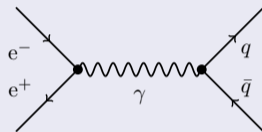
→ Building blocks for transition **amplitudes**:

$$\langle f | S | i \rangle = (2\pi)^4 \delta \left(\sum_j p_j - \sum_j k_j \right) i\mathcal{M}_{fi}, \quad |i\rangle \neq |f\rangle$$

→ **Cross section:** Probability to measure a scattering process

$$\sigma \propto \int d\Phi |\mathcal{M}|^2$$

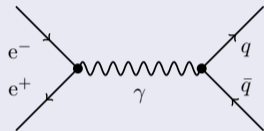
$$e^+(p_1) e^-(p_2) \rightarrow q(k_1) \bar{q}(k_2) \quad [\text{QED}]$$



→ **Cross section:** Probability to measure a scattering process

$$\sigma \propto \int d\Phi |\mathcal{M}|^2$$

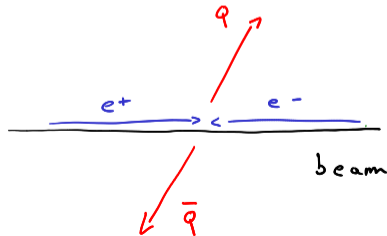
$$e^+(p_1) e^-(p_2) \rightarrow q(k_1) \bar{q}(k_2) \quad [\text{QED}]$$



- $d\Phi \propto d\phi d\cos\theta$

- $|\mathcal{M}|^2 \propto 1 + \cos^2\theta$

$$\Rightarrow \sigma_{\text{incl}} \propto \int_0^{2\pi} d\phi \int_{-1}^{+1} d\cos\theta (1 + \cos^2\theta)$$



- Experimental measurements performed in parts of the phase space
→ where detectors are placed
- Experiments also measure observables
→ transverse mom. $\left[p_T(p) = \sqrt{p_x^2 + p_y^2} \right]$, rapidity $\left[y(p) = \frac{1}{2} \log \left(\frac{|\vec{p}| + p_z}{|\vec{p}| - p_z} \right) \right]$...
- Analytical integration not an option!

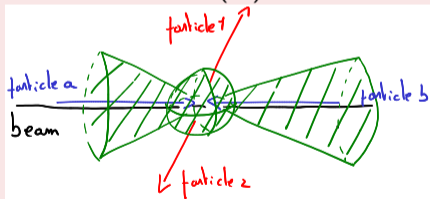
- Experimental measurements performed in parts of the phase space
→ where detectors are placed
- Experiments also measure observables
→ transverse mom. $\left[p_T(p) = \sqrt{p_x^2 + p_y^2} \right]$, rapidity $\left[y(p) = \frac{1}{2} \log \left(\frac{|\vec{p}| + p_z}{|\vec{p}| - p_z} \right) \right]$...
- Analytical integration not an option!

⇒ Solution: **Monte Carlo integration!**

→ Trade integration variables (angles, energies) vs. random numbers (x_i)

Example: $\phi = 2\pi \cdot x_1$, $\cos \theta = 2x_2 - 1$

⇒ $\sigma \propto \int_0^1 dx_1 \cdots \int_0^1 dx_n f(x_1, \dots, x_n) \Theta(g(x_1, \dots, x_n))$



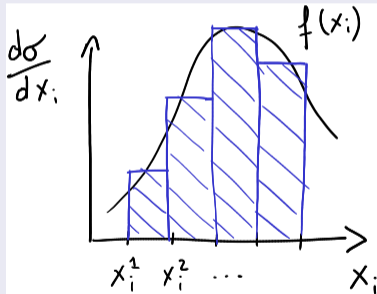
- $\Theta(t)$ encodes experimental selection
- g highly complex function with singularities

$$\sigma \propto \int_0^1 dx_1 \cdots \int_0^1 dx_n f(x_1, \dots, x_n) \Theta(g(x_1, \dots, x_n))$$

- Cross section:

$$\sigma = \sum_{i,l} c_{il} \frac{d\sigma}{dx_l^i}$$

- **Integrating = guessing the values of a function at specific points (Riemann sum)**
- More complex calculations ...
 - ... more integration variables ...
 - ... more computing resources!



Quantum Amplitude Estimate (QAE)

[Brassard, Hoyer, Mosca, Tapp; Quantum Amplitude Amplification and Estimation; quant-ph/0005055]

$$\mathcal{A}|0\rangle = \sqrt{1-a}|\Psi_0\rangle + \sqrt{a}|\Psi_1\rangle$$

QAE estimates a with high probability such that the estimation error scales as $\mathcal{O}(1/M)$ [as opposed to $\mathcal{O}(1/\sqrt{M})$], M : number of applications of \mathcal{A}

Quantum Amplitude Estimate (QAE)

[Brassard, Hoyer, Mosca, Tapp; Quantum Amplitude Amplification and Estimation; quant-ph/0005055]

$$\mathcal{A}|0\rangle = \sqrt{1-a}|\psi_0\rangle + \sqrt{a}|\psi_1\rangle$$

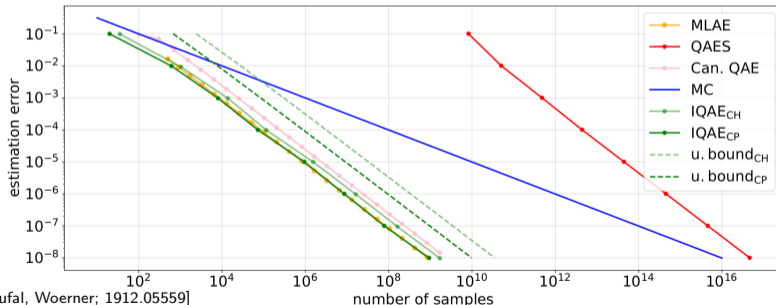
QAE estimates a with high probability such that the estimation error scales as $\mathcal{O}(1/M)$ [as opposed to $\mathcal{O}(1/\sqrt{M})$], M : number of applications of \mathcal{A}

Quantum Amplitude Estimate (QAE)

[Brassard, Hoyer, Mosca, Tapp; Quantum Amplitude Amplification and Estimation; quant-ph/0005055]

$$\mathcal{A}|0\rangle = \sqrt{1-a}|\Psi_0\rangle + \sqrt{a}|\Psi_1\rangle$$

QAE estimates a with high probability such that the estimation error scales as $\mathcal{O}(1/M)$ [as opposed to $\mathcal{O}(1/\sqrt{M})$], M : number of applications of \mathcal{A}



[Grinko, Gacon, Zoufal, Woerner; 1912.05559]

Resulting estimation error for $a = 1/2$ and 95% confidence level with respect to the required total number of oracle queries.

Quantum integration

Extension to

$$\mathcal{A}|0\rangle = \sum_i a_i |\Psi_i\rangle$$

→ Definition of a piece-wise function with $f(x_i) = a_i$.

Quantum integration

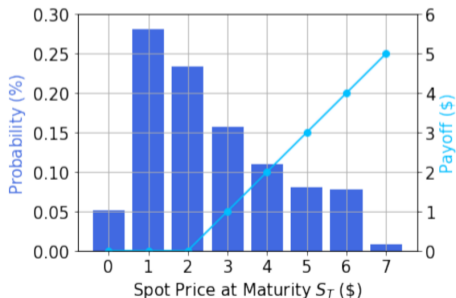
Extension to

$$\mathcal{A}|0\rangle = \sum_i a_i |\Psi_i\rangle$$

→ Definition of a piece-wise function with $f(x_i) = a_i$.

Initially used in finance for simple functions in 1D [Zoufal, Lucchi, Woerner; 1904.00043]

[Woerner and Egger; 1806.06893], [Stamatopoulos et al.; 1905.02666, 2111.12509], [Rebentrost, Gupta, Bromley; Phys.Rev.A 98 (2018) 022321]



$$I = \int dx f(x)g(x)$$

[Zoufal, Lucchi, Woerner; 1904.00043]

Quantum integration

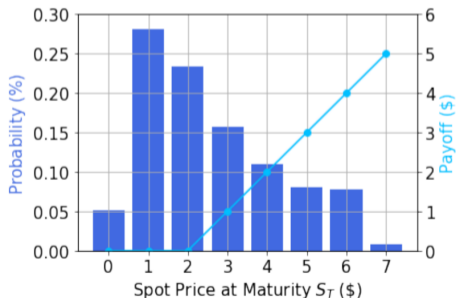
Extension to

$$\mathcal{A}|0\rangle = \sum_i a_i |\Psi_i\rangle$$

→ Definition of a piece-wise function with $f(x_i) = a_i$.

Initially used in finance for simple functions in 1D [Zoufal, Lucchi, Woerner; 1904.00043]

[Woerner and Egger; 1806.06893], [Stamatopoulos et al.; 1905.02666, 2111.12509], [Rebentrost, Gupt, Bromley; Phys.Rev.A 98 (2018) 022321]



[Zoufal, Lucchi, Woerner; 1904.00043]

$$I = \int dx f(x)g(x)$$

- In finance:
 - f : probability
 - g : payoff
- In high-energy physics:
 - f : $|\mathcal{M}|^2$
 - g : $\Theta(\Phi - \Phi_c)$

Idea to apply QAE to high-energy physics [Agliardi, Grossi, MP, Prati; 2201.01547]

→ Followed by [Martínez de Lejarza et al.; 2204.06496, 2401.03023, 2409.12236,2506.19965], [Williams, MP; 2502.14647]

Idea to apply QAE to high-energy physics [Agliardi, Grossi, MP, Prati; 2201.01547]

→ Followed by [Martínez de Lejarza et al.; 2204.06496, 2401.03023, 2409.12236,2506.19965], [Williams, MP; 2502.14647]

Two-step procedure

- 1 Encode the $a_i = f(x_i)$ of the integrand f into the state $\mathcal{A}|0\rangle = \sum_i a_i |\Psi_i\rangle$
- 2 Guess (=integrate) the a_i using QAE

Idea to apply QAE to high-energy physics [Agliardi, Grossi, MP, Prati; 2201.01547]

→ Followed by [Martínez de Lejarza et al.; 2204.06496, 2401.03023, 2409.12236,2506.19965], [Williams, MP; 2502.14647]

Two-step procedure

- 1 Encode the $a_i = f(x_i)$ of the integrand f into the state $\mathcal{A}|0\rangle = \sum_i a_i |\Psi_i\rangle$
- 2 Guess (=integrate) the a_i using QAE

Encode the a_i

- Analytical expression of the integrand (quantum arithmetics)
- Numerical implementation of the integrand (following application)

→ Both approaches used for classical Monte Carlo

[Williams, MP; 2502.14647]

Quantum arithmetics + use of QUANTINUUM Quantum Monte Carlo Integrator

[Akhalya et al.; 2308.06081], based on [Herbert; 2105.09100]

[Williams, MP; 2502.14647]

Quantum arithmetics + use of QUANTINUUM Quantum Monte Carlo Integrator

[Akhilwaya et al.; 2308.06081], based on [Herbert; 2105.09100]

- 2D integral

$$\int \prod_{i=1}^{N_I} dx_i \frac{\sum_{S_k \in I} \alpha_k \prod_{j \in S_k} x_j^{n_j}}{\prod_{p=1}^{N_P} (x_p - M_{op}^2)^2 + M_{op}^2 \Gamma_{op}^2} \rightarrow \int_0^s \int_0^{s-s_2} ds_1 ds_2 \frac{s_1^2 s - s_1 M_\tau^2 s + s_1 M_\tau^2 s_2}{(s_2 - M_W^2)^2 + (M_W \Gamma_W)^2}$$

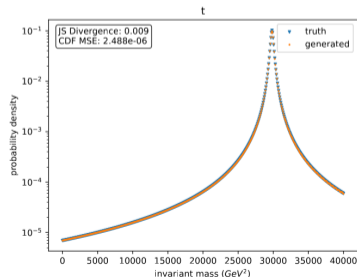
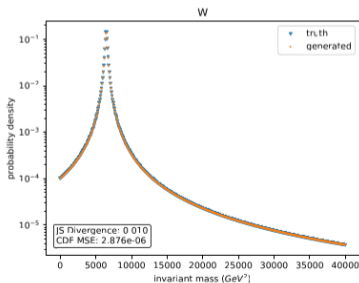
Quantum arithmetics + use of QUANTINUUM Quantum Monte Carlo Integrator

[Akhmalwaya et al.; 2308.06081], based on [Herbert; 2105.09100]

- 2D integral

$$\int \prod_{i=1}^{N_I} dx_i \frac{\sum_{S_k \in I} \alpha_k \prod_{j \in S_k} x_j^{n_j}}{\prod_{p=1}^{N_P} (x_p - M_{op}^2)^2 + M_{op}^2 \Gamma_{op}^2} \rightarrow \int_0^s \int_0^{s-s_2} ds_1 ds_2 \frac{s_1^2 s - s_1 M_T^2 s + s_1 M_T^2 s_2}{(s_2 - M_W^2)^2 + (M_W \Gamma_W)^2}$$

- Building blocks \rightarrow Propagator (using Fourier expansion) [Rosenkranz et al.; 2405.21058]



Quantum arithmetics + use of QUANTINUUM Quantum Monte Carlo Integrator

[Akhilwaya et al.; 2308.06081], based on [Herbert; 2105.09100]

- 2D integral

$$\int \prod_{i=1}^{N_I} dx_i \frac{\sum_{S_k \in I} \alpha_k \prod_{j \in S_k} x_j^{n_j}}{\prod_{p=1}^{N_P} (x_p - M_{op}^2)^2 + M_{op}^2 \Gamma_{op}^2} \rightarrow \int_0^s \int_0^{s-s_2} ds_1 ds_2 \frac{s_1^2 s - s_1 M_T^2 s + s_1 M_T^2 s_2}{(s_2 - M_W^2)^2 + (M_W \Gamma_W)^2}$$

- Building blocks → Propagator (using Fourier expansion) [Rosenkranz et al.; 2405.21058]
- Resource estimate → Fault-tolerant era

Compilation	Resource	Metric	Precision		
			10%	1%	0.1%
NISQ	Number of qubits	Largest across circuits	28	28	28
	CX gates	Total number across circuits	7.39×10^7	6.15×10^8	5.09×10^9
		Total depth across circuits	4.34×10^7	3.61×10^8	2.99×10^9
		Number in largest circuit	3.39×10^7	2.71×10^8	1.20×10^9
		Depth of largest circuit	1.99×10^7	1.59×10^8	7.03×10^8
	All gates	Total number across circuits	1.50×10^8	1.24×10^9	1.03×10^{10}
		Total depth across circuits	8.02×10^7	6.67×10^8	5.52×10^9
		Number in largest circuit	6.86×10^7	5.49×10^8	2.42×10^9
		Depth of largest circuit	3.68×10^7	2.94×10^8	1.30×10^9
	Fault tolerant	Number of qubits	Largest across circuits	41	41
T gates		Total number across circuits	3.39×10^9	4.08×10^{10}	2.72×10^{11}
		Total depth across circuits	3.29×10^9	3.95×10^{10}	2.63×10^{11}
		Number in largest circuit	1.65×10^9	2.14×10^{10}	6.97×10^{10}
		Depth of largest circuit	1.60×10^9	2.07×10^{10}	6.75×10^{10}

How?

- Applications

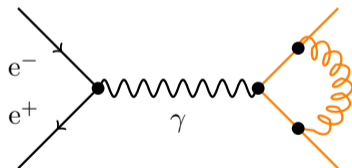
- Cross section → quantum integration
- Amplitudes in QCD → dedicated quantum algorithms
- QCD parton shower → dedicated quantum algorithms

Quantum implementation of QCD amplitude

→ Numerical implementation of integrand

$$\mathcal{M} \sim \sum \text{Tr}(T^{a_1} \dots T^{a_n}) \mathcal{K}(1, \dots, n)$$

- **Kinematic part:** made of spinors and tensors (and kinematic invariants)
- **Colour part:** made of SU(3) generators of QCD

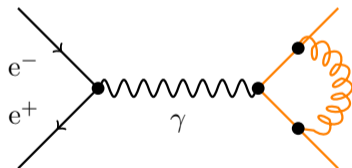


Quantum implementation of QCD amplitude

→ Numerical implementation of integrand

$$\mathcal{M} \sim \sum \text{Tr}(T^{a_1} \dots T^{a_n}) \mathcal{K}(1, \dots, n)$$

- **Kinematic part:** made of spinors and tensors (and kinematic invariants)
- **Colour part:** made of SU(3) generators of QCD



- First step towards a full quantum amplitude/Monte Carlo
- Useful for a quantum parton shower (see later)

- $[T^a, T^b] = if_{abc} T^c$.
- T^a, T^c, \dots : SU(3) generators
- Gell-Mann matrices

$$T^1 = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad T^2 = \frac{1}{2} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad T^3 = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad T^4 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix},$$
$$T^5 = \frac{1}{2} \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix}, \quad T^6 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad T^7 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad T^8 = \frac{1}{2\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}.$$

Attention! T^a are not unitary!

→ In our example, colour factor: $T_{ij}^a T_{jk}^a = C_F \delta_{ik}$

Quantum implementation of colour

→ Implementation of the (discrete) colour into qubits:

- Gluon: 8 colours → 3 qubits ($2^3 = 8$)
- Quark: 3 colours → 2 qubits ($2^2 = 4$)

Quantum implementation of colour

→ Implementation of the (discrete) colour into qubits:

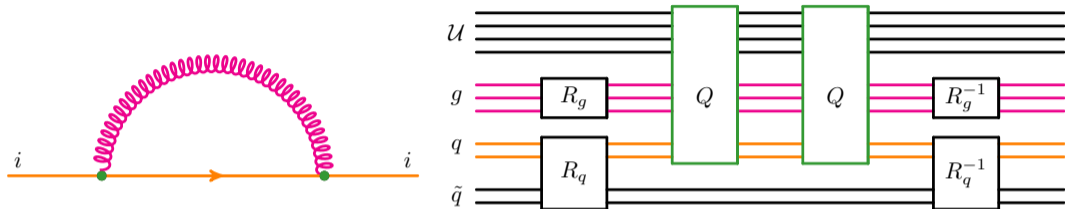
- Gluon: 8 colours → 3 qubits ($2^3 = 8$)
- Quark: 3 colours → 2 qubits ($2^2 = 4$)

Make non-unitary matrices unitary again!

→ Extend dimension and modify them

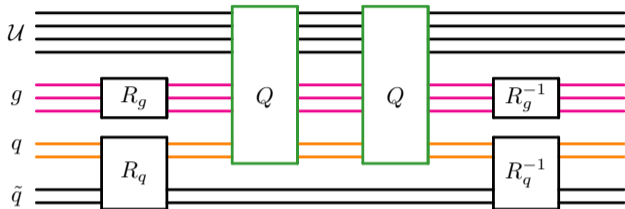
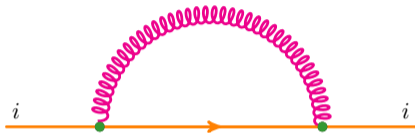
$$\begin{aligned} \overline{T^1} &= \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^2} &= \frac{1}{2} \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^3} &= \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^4} &= \frac{1}{2} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \\ \overline{T^5} &= \frac{1}{2} \begin{pmatrix} 0 & 0 & -i & 0 \\ 0 & 1 & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^6} &= \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^7} &= \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, & \overline{T^8} &= \frac{1}{2\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}. \end{aligned}$$

Example

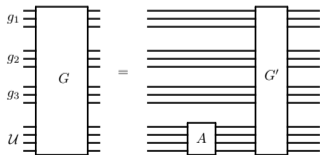


- One-to-one correspondence between Feynman diagram and circuit
- Gates for qqg (Q) and ggg (G) vertices to simulate QCD (colour) interaction

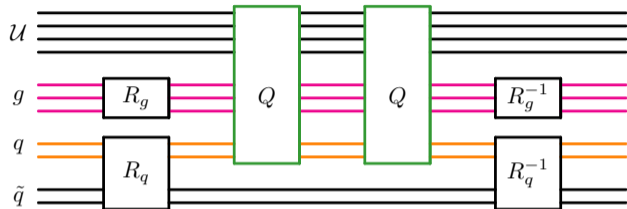
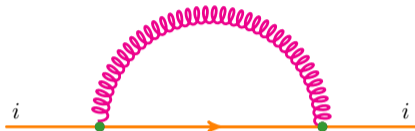
Example



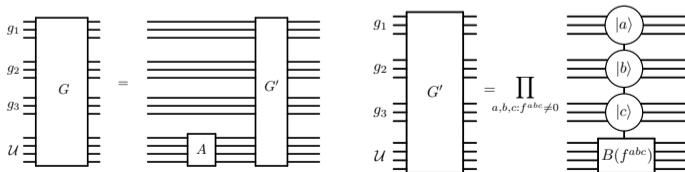
- One-to-one correspondence between Feynman diagram and circuit
- Gates for qqg (Q) and ggg (G) vertices to simulate QCD (colour) interaction



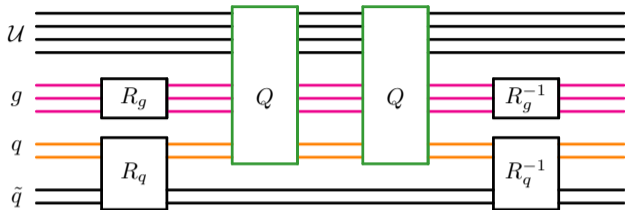
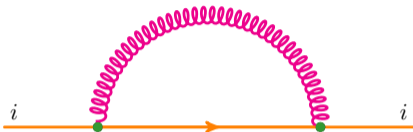
Example



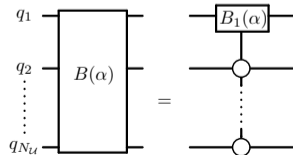
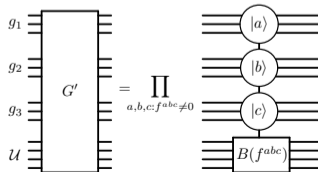
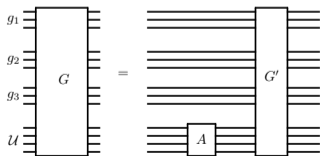
- One-to-one correspondence between Feynman diagram and circuit
- Gates for qqg (Q) and ggg (G) vertices to simulate QCD (colour) interaction



Example

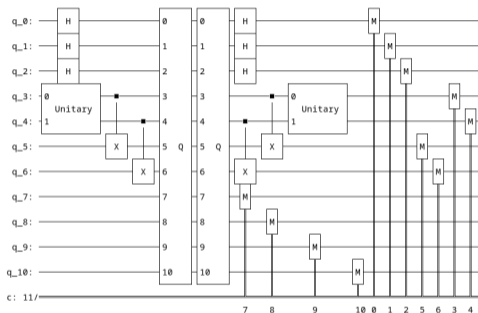


$$\langle \Omega |_{\mathcal{U}} \prod_{i=1}^{N_{ops}} \{B(\alpha_i)A\} | \Omega \rangle_{\mathcal{U}} = \prod_{i=1}^{N_{ops}} \alpha_i$$



→ Trace defined in $|00000000000\rangle = |0_{11}\rangle$ state: $|\psi\rangle = \frac{c}{N}|0_{11}\rangle + \dots$

$$\Rightarrow \frac{27415}{N_{\text{shots}}=1000000} \sim \left(\frac{c}{N} = \frac{(N_c=3)C_F}{N_c^{n_q=1}(N_c^2-1)^{n_g=1}} \right)^2 \sim \alpha^2 \text{ in intro}$$

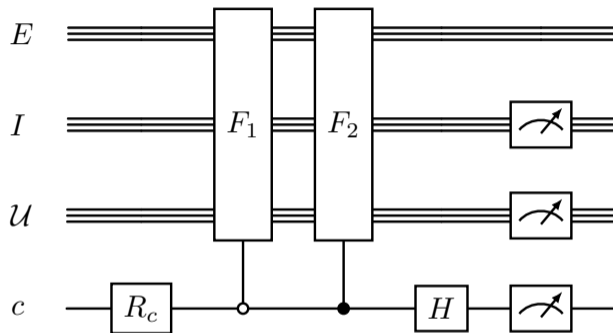


- Colour factors encoded in one single state $|0_{11}\rangle$ (as needed for QAE)
- Any colour factor can be computed

[Chawdhry, MP; 2303.04818]

```
Total counts are: ('00000010101': 226, '00010010011': 342, '00000010110': 225, '00000001101': 696, '00010010010': 362, '00000010001': 872, '00010001101': 2006, '00000001100': 643, '00010001100': 2006, '00010001110': 1051, '00000001111': 638, '00000010010': 904, '00010000010': 1057, '00100010101': 52353, '00010010100': 3342, '00100000101': 6942, '00010000111': 1046, '00000010111': 210, '00100001101': 36223, '00100010111': 51877, '00000001110': 643, '00010010000': 4838, '00000010100': 5421, '00010000100': 1035, '00100000000': 187280, '00010000011': 1075, '00100000111': 6795, '00100001011': 145043, '00100001010': 10275, '00010000000': 65548, '00000000000': 27415, '00010001111': 2031, '00100000110': 7004, '00000001011': 2551, '00100001111': 36471, '00010010111': 8866, '00010000101': 1077, '00100010110': 52220, '00010010110': 9080, '00100010100': 185856, '00100001100': 36173, '00010010101': 8860, '00100010000': 14129, '00100001110': 36925, '00100000100': 6858, '00100010011': 10182, '00100000001': 6950, '00010001110': 2018, '00100000011': 6983, '00000010011': 815, '00010001011': 7957, '00010010001': 340, '00100010001': 10163, '00010000001': 1092, '00100000010': 7010)
```

→ Interference between two amplitudes (F_1 and F_2)

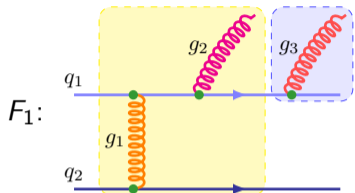


→ c is an equal superposition of F_1 and F_2

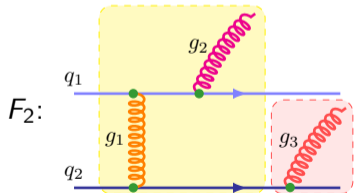
→ Making the measurement (\sim square/collapse the amplitude):

$$|\mathcal{M}|^2 = |F_1 + F_2|^2 = |F_1|^2 + |F_2|^2 + 2 \operatorname{Re} \{F_1 \cdot F_2^*\}$$

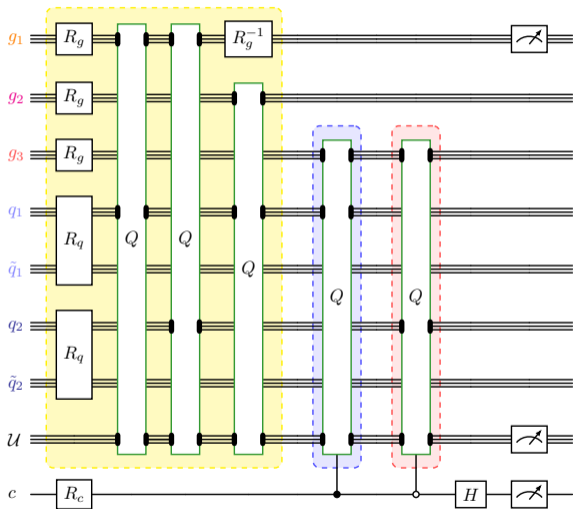
→ Common building block factored out



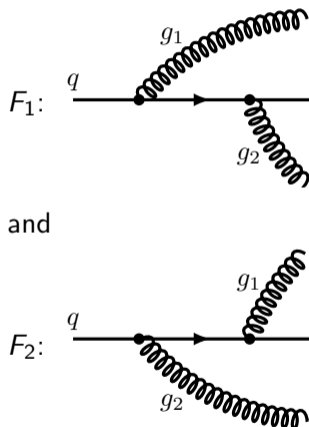
and



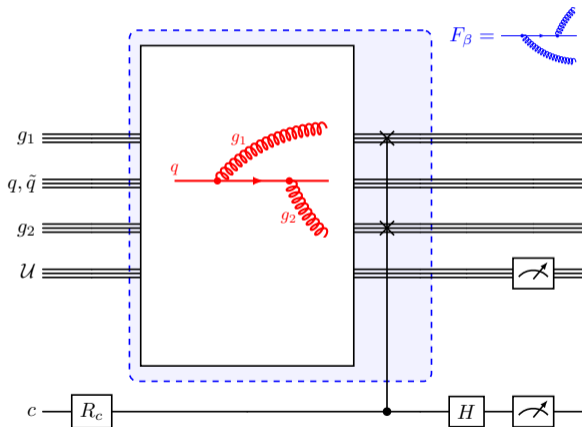
[Chawdhry, MP, Williams; 2507.07194]



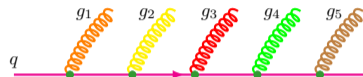
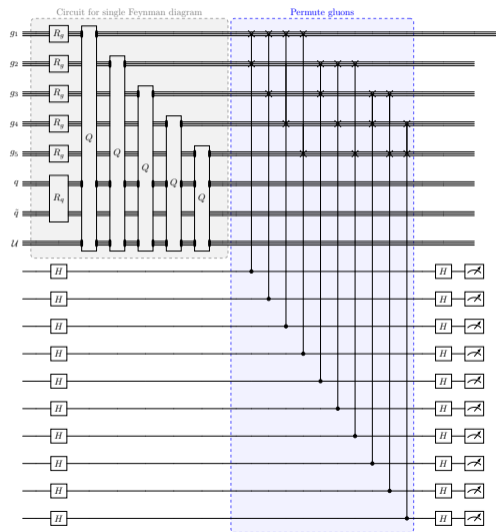
Common building block factored out + simple operation to obtain second amplitude
 → Hint for quadratic improvement with respect to classical approach



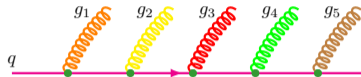
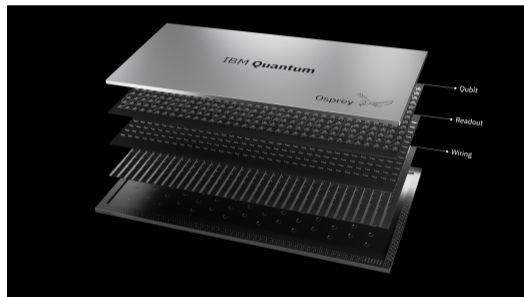
[Chawdhry, MP, Williams; 2507.07194]



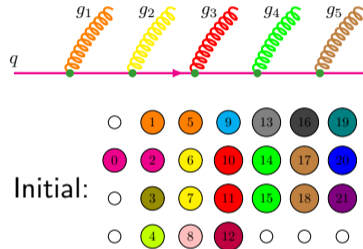
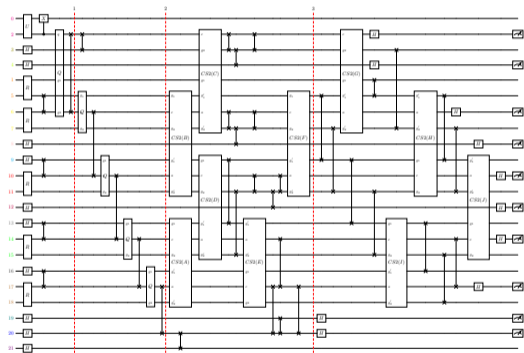
→ Application to IBM 120 qubits superconducting with squared lattice
(Nightawk generation) [Chawdhry, Pascuzzi, MP; in preparation]



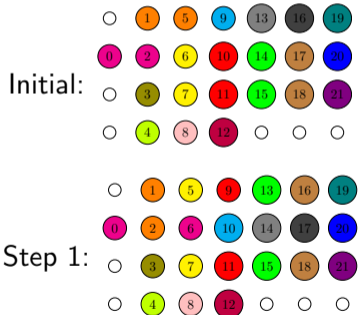
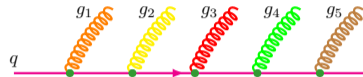
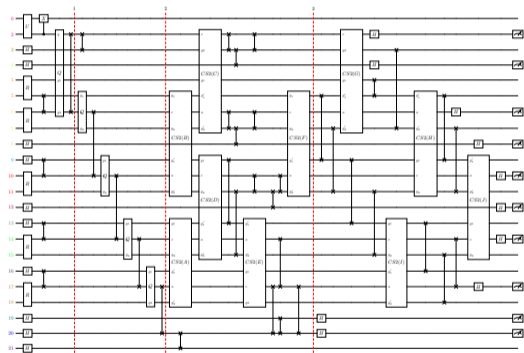
→ Application to IBM 120 qubits superconducting with squared lattice
(Nighthawk generation) [Chawdhry, Pascuzzi, MP; in preparation]



→ Application to IBM 120 qubits superconducting with squared lattice
 (Nightawk generation) [Chawdhry, Pascuzzi, MP; in preparation]



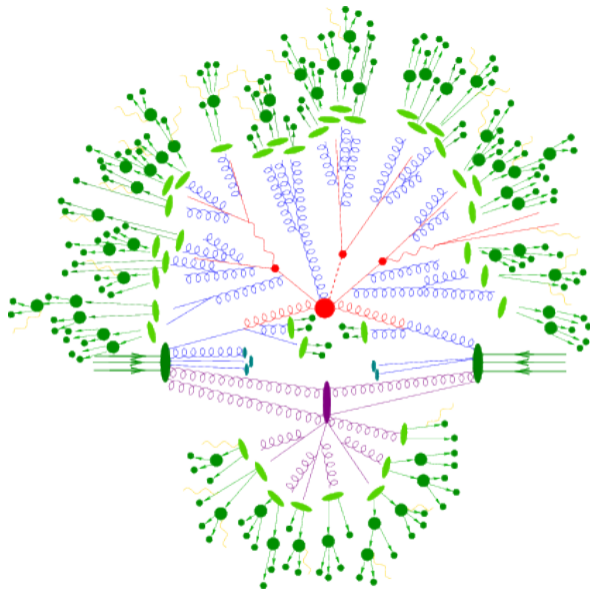
→ Application to IBM 120 qubits superconducting with squared lattice
 (Nightawk generation) [Chawdhry, Pascuzzi, MP; in preparation]



How?

- Applications

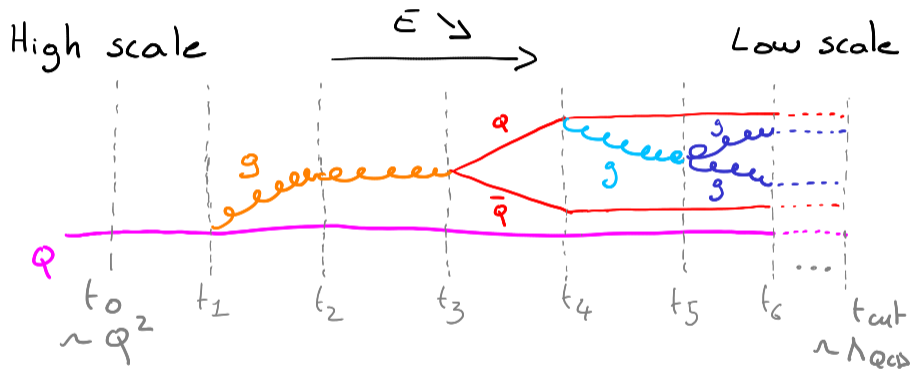
- Cross section → quantum integration
- Amplitudes in QCD → dedicated quantum algorithms
- QCD parton shower → dedicated quantum algorithms



[source: Sherpa]

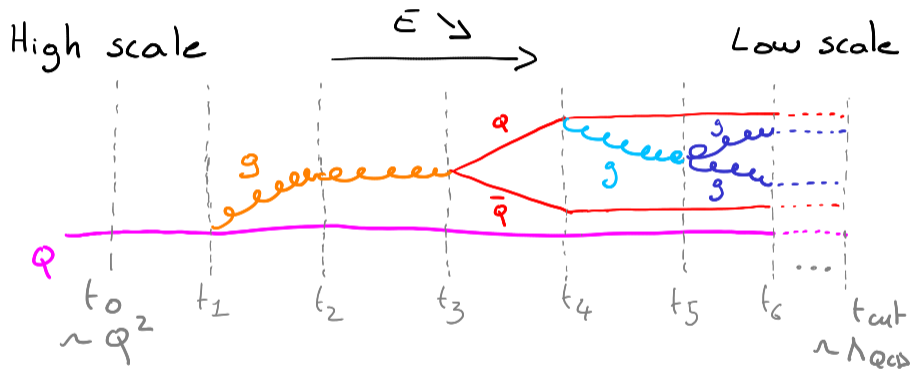
Parton shower

Algorithmic and approximate way of emitting an arbitrary number of particles



Parton shower

Algorithmic and approximate way of emitting an arbitrary number of particles



Advantages

- Universal

Drawbacks

- Approximate

→ First proposal for Quantum Parton Shower [Bauer, de Jong, Nachman, Provasoli; 1904.03196]

Toy model:
$$\mathcal{L} = \bar{f}_1(i\not{\partial} + m_1)f_1 + \bar{f}_2(i\not{\partial} + m_2)f_2 + (\partial_\mu\phi)^2$$
$$+ g_1\bar{f}_1f_1\phi + g_2\bar{f}_2f_2\phi + g_{12}[\bar{f}_1f_2 + \bar{f}_2f_1]\phi$$

- Splittings: $f_i \rightarrow f_j\phi$ and $\phi \rightarrow f_i\bar{f}_j$
- Mixing between f_1 and f_2

→ First proposal for Quantum Parton Shower [Bauer, de Jong, Nachman, Provasoli; 1904.03196]

Toy model:
$$\mathcal{L} = \bar{f}_1(i\not{\partial} + m_1)f_1 + \bar{f}_2(i\not{\partial} + m_2)f_2 + (\partial_\mu\phi)^2$$
$$+ g_1\bar{f}_1f_1\phi + g_2\bar{f}_2f_2\phi + g_{12}[\bar{f}_1f_2 + \bar{f}_2f_1]\phi$$

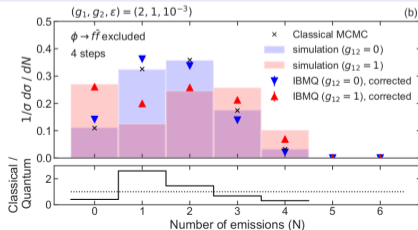
- Splittings: $f_i \rightarrow f_j\phi$ and $\phi \rightarrow f_i\bar{f}_j$
- Mixing between f_1 and f_2
- **Idea:** Evolution with (quantum) superposition of f_1 and f_2
 - allows to take into account interference between different histories
- Classically requires a reweighting strategy at each step (costly)
 - claim of exponential speed-up

→ First proposal for Quantum Parton Shower [Bauer, de Jong, Nachman, Provasoli; 1904.03196]

Toy model:
$$\mathcal{L} = \bar{f}_1(i\not{\partial} + m_1)f_1 + \bar{f}_2(i\not{\partial} + m_2)f_2 + (\partial_\mu\phi)^2$$

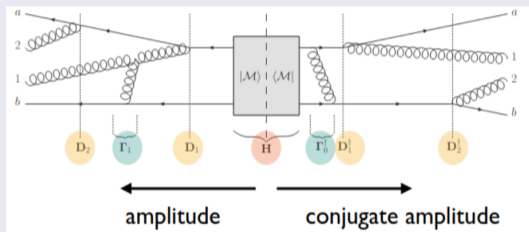
$$+ g_1\bar{f}_1f_1\phi + g_2\bar{f}_2f_2\phi + g_{12} [\bar{f}_1f_2 + \bar{f}_2f_1] \phi$$

- Splittings: $f_i \rightarrow f_j\phi$ and $\phi \rightarrow f_i\bar{f}_j$
- Mixing between f_1 and f_2
- **Idea:** Evolution with (quantum) superposition of f_1 and f_2
 - allows to take into account interference between different histories
- Classically requires a reweighting strategy at each step (costly)
 - claim of exponential speed-up



Outlook

- In SM, mixing for $\gamma/Z \rightarrow$ EW PS
- Account for different PS history (and interference) \rightarrow Amplitude PS
[Nagy, Soper; 0706.0017,1902.02105], [Plätzer et al.; 1201.0260,1802.08531,1808.00332], [Isaacson, Prestel; 1806.10102]



- Perfect framework for implementation of previous quantum interaction gates

- Is it possible?
→ **Yes.**

Quantum simulation in high-energy physics

- Is it possible?

→ **Yes.**

- Is there a quantum advantage?

→ **In principle, yes. In practice, not yet.**

Quantum simulation in high-energy physics

- Is it possible?

→ **Yes.**

- Is there a quantum advantage?

→ **In principle, yes. In practice, not yet.**

- Is it more resource efficient/faster than CPU/GPU/AI approaches?

→ **At the moment, not known.**

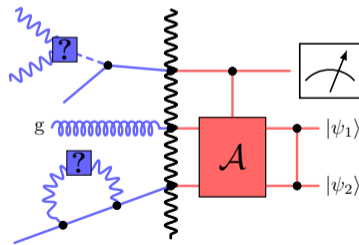


- Are we witnessing a quantum revolution?
- When can we do *state-of-the-art* quantum computations?

Quantum simulation of High-Energy physics

→ High-energy physics: computationally-intensive field

- **GPU & AI:** Technical, general-purpose approaches
- **Quantum computing:** Tailored to quantum systems (at very high energies)

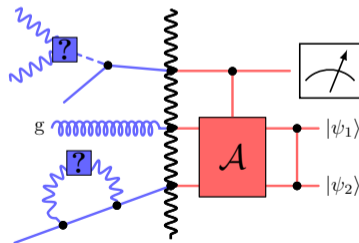


Quantum simulation of High-Energy physics

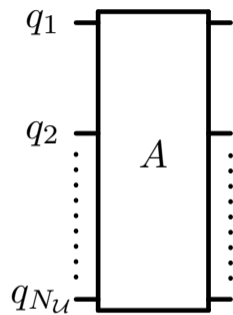
→ High-energy physics: computationally-intensive field

- **GPU & AI:** Technical, general-purpose approaches
- **Quantum computing:** Tailored to quantum systems (at very high energies)

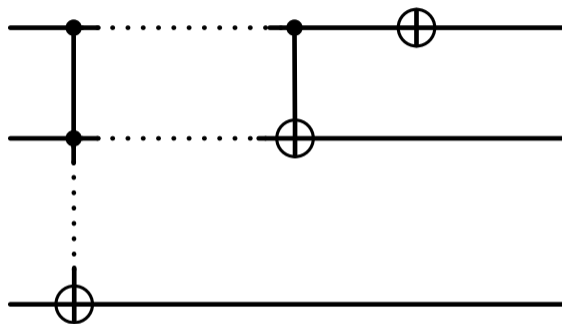
- Several interesting applications already
- Powerful quantum computers becoming available
→ entering era with non-trivial applications
- Exciting emerging field with many new ideas!



BACK-UP



=



$$B_1(\alpha) = \begin{pmatrix} \sqrt{1 - |\alpha|^2} & \alpha \\ -\alpha & \sqrt{1 - |\alpha|^2} \end{pmatrix} \quad (1)$$

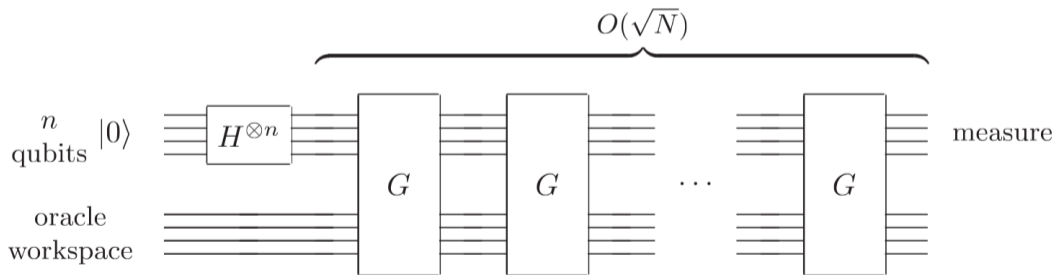
$$B(\alpha)A|k\rangle = \begin{cases} \alpha|0\rangle + \sqrt{1 - |\alpha|^2}|1\rangle & \text{if } k = 0 \\ |k + 1\rangle & \text{if } 0 < k < 2^{N_{\mathcal{U}}} - 1 \\ \sqrt{1 - |\alpha|^2}|0\rangle - \alpha|1\rangle & \text{if } k = 2^{N_{\mathcal{U}}} - 1 \end{cases} \quad (2)$$

$$\langle \Omega|_{\mathcal{U}} B(\alpha)A|\Omega\rangle_{\mathcal{U}} = \alpha \quad (3)$$

$$\langle \Omega|_{\mathcal{U}} \prod_{i=1}^{N_{ops}} \{B(\alpha_i)A\} |\Omega\rangle_{\mathcal{U}} = \prod_{i=1}^{N_{ops}} \alpha_i \quad (4)$$

Grover algorithm/iteration

- Very general quantum algorithm
- Quadratic speed up
→ $\mathcal{O}(\sqrt{N})$ operations instead of $\mathcal{O}(N)$
- Most famous example: unstructured database search



Quantum Amplitude Estimate (QAE)

[Brassard, Hoyer, Mosca, Tapp; Quantum Amplitude Amplification and Estimation; quant-ph/0005055]

$$\mathcal{A}|0\rangle = \sqrt{1-a}|\Psi_0\rangle + \sqrt{a}|\Psi_1\rangle$$

QAE estimates a with high probability such that the estimation error scales as $\mathcal{O}(1/M)$
[as opposed to $\mathcal{O}(1/\sqrt{M})$]

M : number of applications of \mathcal{A}

Quantum Amplitude Estimate (QAE)

[Brassard, Hoyer, Mosca, Tapp; Quantum Amplitude Amplification and Estimation; quant-ph/0005055]

$$\mathcal{A}|0\rangle = \sqrt{1-a}|\Psi_0\rangle + \sqrt{a}|\Psi_1\rangle$$

QAE estimates a with high probability such that the estimation error scales as $\mathcal{O}(1/M)$ [as opposed to $\mathcal{O}(1/\sqrt{M})$]

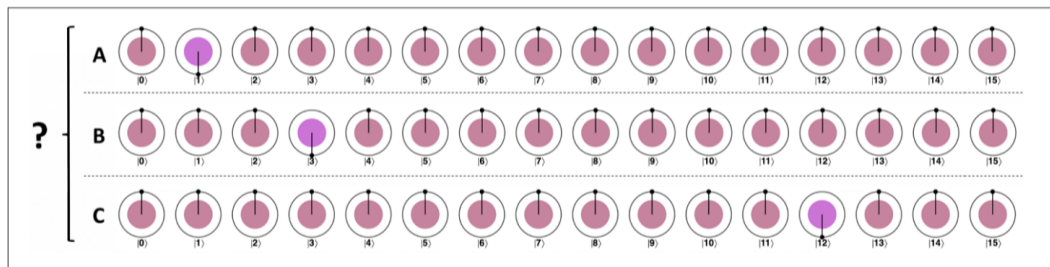
M : number of applications of \mathcal{A}

→ What the (original) algorithm provides:

- An estimate: $\tilde{a} = \sin^2(\tilde{\theta}_a)$
with $\tilde{\theta}_a = y\pi/M$, $y \in \{0, \dots, M-1\}$, and $M = 2^n$
- A success probability (that can be increased by repeating the algorithm)
- A bound: $|a - \tilde{a}| \leq \mathcal{O}(1/M)$

Grover algorithm/iteration

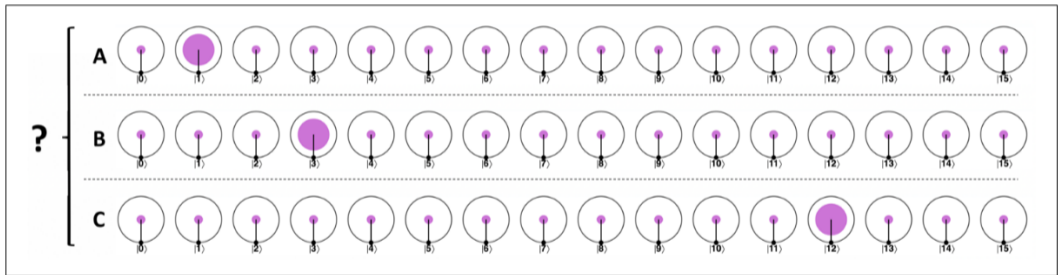
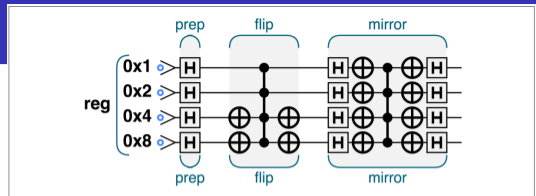
- Example (from [Johnston, Harrigan, Gimeno-Segovia; Programming Quantum Computers])



→ What solution is contained in our quantum register?

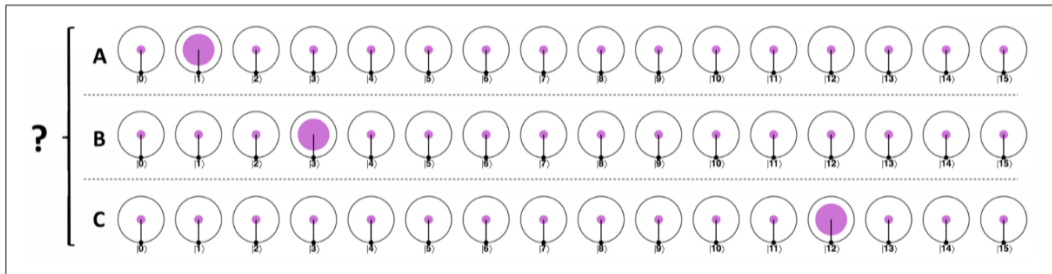
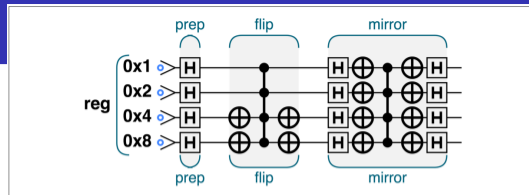
Grover algorithm/iteration

→ Applying a Grover iteration

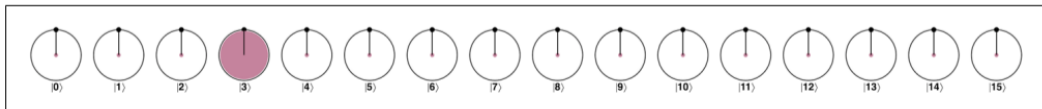


Grover algorithm/iteration

→ Applying a Grover iteration



→ Applying it twice (e.g. solution B)



- $e^+e^- \rightarrow q\bar{q}$ (in QED) \rightarrow 1D problem

$$\sigma \sim \int_{-1}^1 \int_0^{2\pi} d\cos\theta d\phi (1 + \cos^2\theta)$$

- $e^+e^- \rightarrow q\bar{q}'W$ \rightarrow 2D problem

$$\begin{aligned}\sigma &\sim \int_{M_W^2}^s \int_0^{s_1^{\text{Max}}} \int_{-1}^1 \int_0^{2\pi} \int_0^{2\pi} d\Phi_3 |\mathcal{M}_{e^+e^- \rightarrow q\bar{q}'W}|^2 \\ &\sim \int_{M_W^2}^s \int_0^{s_1^{\text{Max}}} d\tilde{\Phi}_3 |\mathcal{M}'|^2\end{aligned}$$

with $\mathcal{M}' = \mathcal{M}_{e^+e^- \rightarrow q\bar{q}'W}(\cos\theta_1 = 0, \phi_1 = \pi/2, \phi_2 = \pi/2)$.

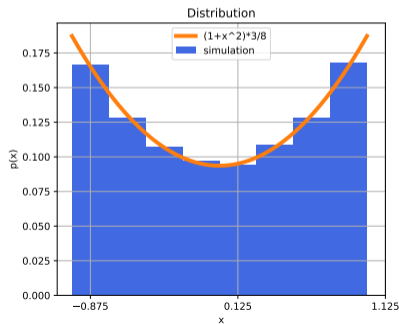
Encoding the distribution to be integrated into qubits (coefficients of states)

- Exact loading [Shende, Bullock, Markov, quant-ph/0406176] (resource intensive)
- Using quantum machine learning (qGAN) [Zoufal, Lucchi, Woerner; 1904.00043] (not exact)

Loading of distribution / encoding into qubits

Encoding the distribution to be integrated into qubits (coefficients of states)

- Exact loading [Shende, Bullock, Markov, quant-ph/0406176] (resource intensive)
- Using quantum machine learning (qGAN) [Zoufal, Lucchi, Woerner; 1904.00043] (not exact)
- Example: Exact loading - $1 + x^2$



→ 3 qubits: $2^3 = 8$ bins

- Matching boundary of integration (3 qubits $\Rightarrow 2^3$ bins)

Domain	low stat.		high stat.		very high stat.		exact	
	σ	$\delta[\%]$	σ	$\delta[\%]$	σ	$\delta[\%]$	σ	$\delta[\%]$
$[-0.75; 0]$	0.345	-3.31	0.332	0.706	0.334	0.0331	0.334	-8.31×10^{-3}
$[-0.5; 0]$	0.215	-5.86	0.201	1.15	0.203	0.0986	0.203	-0.0161
$[-0.25; 0]$	0.112	-17.1	0.0939	1.87	0.0960	-0.284	0.0957	-0.0389

- Matching boundary of integration (3 qubits $\Rightarrow 2^3$ bins)

Domain	low stat.		high stat.		very high stat.		exact	
	σ	δ [%]	σ	δ [%]	σ	δ [%]	σ	δ [%]
$[-0.75; 0]$	0.345	-3.31	0.332	0.706	0.334	0.0331	0.334	-8.31×10^{-3}
$[-0.5; 0]$	0.215	-5.86	0.201	1.15	0.203	0.0986	0.203	-0.0161
$[-0.25; 0]$	0.112	-17.1	0.0939	1.87	0.0960	-0.284	0.0957	-0.0389

- Non-matching boundary of integration

Qubits number	$[-0.7; 0.6]$				$[-0.625; 0.375]$			
	high stat.		exact		high stat.		exact	
	σ	δ [%]	σ	δ [%]	σ	δ [%]	σ	δ [%]
3	0.402	-28.0	0.406	-27.1	0.296	-28.1	0.299	-27.5
4	0.463	-17.0	0.468	-16.0	0.408	-1.07	0.412	5.96×10^{-3}
5	0.527	-5.46	0.532	-4.62	0.408	-1.07	0.412	5.96×10^{-3}
6	0.542	-2.76	0.547	-1.81	0.408	-1.07	0.412	5.96×10^{-3}

Remarks

- Use Qiskit (IBM python software) subroutines and noiseless quantum simulation (perfect quantum computer)
- For present application, too many qubits for test on real hardware
 - 4 qubits for representation → 9 total qubits
 - 6 qubits for representation → 13 total qubits
- Largest quantum computer on IBM quantum experience:
7 qubits previously (127 qubits now)
 - Simulators can go up to 5000 qubits

Remarks

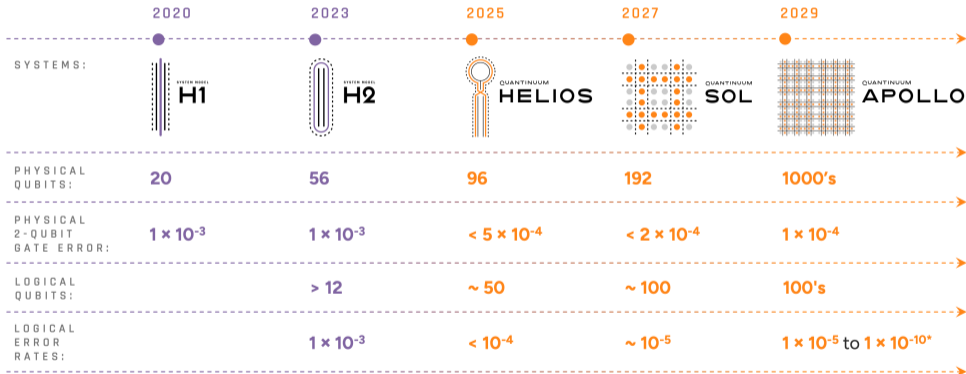
- Use Qiskit (IBM python software) subroutines and noiseless quantum simulation (perfect quantum computer)
- For present application, too many qubits for test on real hardware
 - 4 qubits for representation → 9 total qubits
 - 6 qubits for representation → 13 total qubits
- Largest quantum computer on IBM quantum experience:
7 qubits previously (127 qubits now)
 - Simulators can go up to 5000 qubits

Summary

- First application of quantum integration in HEP
- Theoretical quadratic speed-up
- Main challenge: error estimate

[Williams, MP; 2502.14647]

Development roadmap



© 2024 Quantinuum. All Rights Reserved.

*analysis based on recent literature in new, novel error correcting codes predict that error could be as low as $1E-10$ in Apollo (ref: arXiv:2403.16054, arXiv:2308.07915)